

Figure 3. Lattice constants  $c$  in  $\text{ZrO}_2$  (■),  $\text{Al}_2\text{O}_3$  (+),  $\text{TiO}_2$  (□),  $\text{B}_2\text{O}_3$  (△) coated, and bare  $\text{LiCoO}_2$  (●) as a function of  $x$  in  $\text{Li}_{1-x}\text{CoO}_2$  during the first charge ( $c$  = lattice constant).

nonuniform strain during delithiation produces a new-class of cathode materials, which are extremely tolerant to electrochemical cycling. Note that the lattice constant  $c$  of the  $\text{ZrO}_2$ -coated  $\text{LiCoO}_2$  exhibits negligible shift in the range  $0 < x < 0.7$ , which indicates that the original hexagonal symmetry is unaffected by delithiation. A sol-gel coating of  $\text{ZrO}_2$ , which has the highest fracture toughness, leads to the formation of a fracture-toughened thin-film solid solution (tens of nm thick) near the particle surface. This film significantly improves the structural stability of the cathode material by suppressing  $c$ -axis expansion (or phase transition), thereby preventing capacity fading during electrochemical cycling.

In conclusion, the encapsulation of  $\text{LiCoO}_2$  powders by thin-film coating of high-fracture-toughness oxides effectively suppresses the lattice-constant changes during electrochemical cycling and thereby suppresses phase transitions. The order of capacity retention correlates well with the fracture toughness of the coated thin-film oxides. This innovative technology can be applied to any cathode material that has an accompanying lattice strain (or phase transition) during cycling. The commercial potential of a rechargeable Li-ion battery is enormous.

### Experimental Section

Metal (or boron) ethylhexanate-diisopropoxide ( $\text{M}(\text{OOC}_8\text{H}_{15})(\text{OC}_3\text{H}_7)_2$ ) was dissolved in 2-propanol, and stirred continuously for 20 h at 21 °C. The  $\text{LiCoO}_2$  powder was then mixed with the coating solution, and the resulting solution was further aged at 50 °C for 24 h to give strong bonds between the metal-oxide (or boron-oxide) gel and the  $\text{LiCoO}_2$  particle surface. After drying, each coated  $\text{LiCoO}_2$  powder was fired at 400 °C for 10 h.

To test the cycle-life performance of each cathode material, a slurry of the desired viscosity was prepared by mixing the cathode materials, super P carbon black, and a poly(vinylidene fluoride) (PVDF) binder in a weight ratio of 92:4:4 in *N*-methyl-2-pyrrolidone (NMP; water content was below 0.01 wt %). The coin-type half cell (2016-size) contained a test cathode, a lithium metal counter-and-reference electrode, a 15  $\mu\text{m}$  thick microporous polyethylene separator, and an electrolyte solution of 1M  $\text{LiPF}_6$  in EC/DMC (1:1 vol % EC = ethylene carbonate, DMC = dimethyl carbonate). Each cell was aged for 24 h at room temperature prior to commencing the electrochemical tests at 21 °C. The coin-type half cells in  $\text{Li}/\text{LiCoO}_2$  were charged with 0.1 C rate to the pre-determined voltages, then potentiostated for 10 h until the current density was equivalent to 0.01 C.

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## Aromatic Mercury Clusters in Ancient Amalgams\*\*

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Mercury and its alloys, named amalgams, were known to the Ancients. Amalgams were used for metal extraction from about 500 BC and one of the most interesting influences of

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amalgams (many of which have a golden appearance) on human history was the idea of transmutation (the conversion of base metal into gold). Efforts of transmutation flourished during the time of alchemy, when many gullible rulers of that time found themselves in embarrassing circumstances thanks to the golden appearance of the amalgams. A wide range of alkali metal amalgams is now known with different stoichiometries and structures.<sup>[1]</sup> Low mercury-content amalgams contain isolated Hg atoms or clusters solvated by the alkali atoms, whereas high mercury-content amalgams tend to have network structures with the alkali atoms occupying interstitial positions. A substantial amount of charge transfer is believed to occur from the alkali metal atoms to mercury in the amalgams. One may think that amalgams, having such a rich and long history, should be well understood today. However, as we will show herein, that is not the case, and the chemistry, structure, and bonding of the amalgams are still not completely elucidated.

What we want to address here is a class of amalgams that contain  $\text{Hg}_4^{6-}$  square units as their building blocks, typified by that found in solid  $\text{Na}_3\text{Hg}_2$  (Figure 1).<sup>[1,2]</sup> Why  $\text{Hg}_4^{6-}$  is a

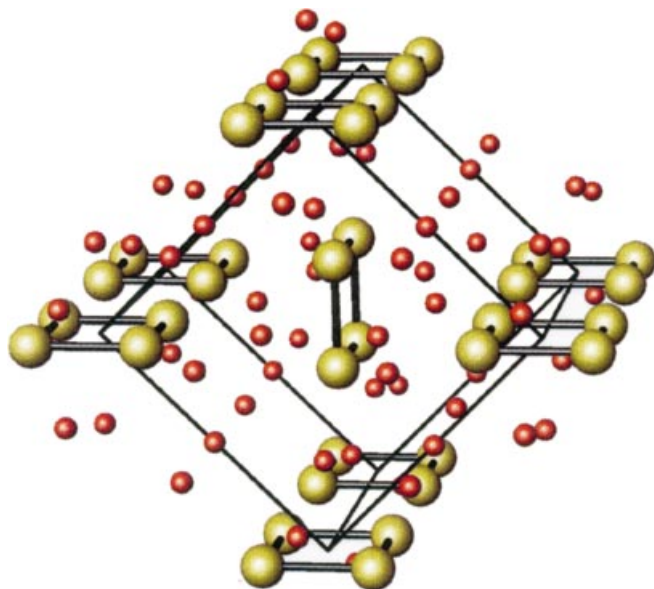


Figure 1. Crystal structure of  $\text{Na}_3\text{Hg}_2$ , showing the  $\text{Hg}_4$  square units (yellow) surrounded by Na (red).

particularly stable building block for the amalgams and why it assumes a perfect square structure are questions that have not been understood. To the best of our knowledge, Corbett was the first to examine these questions using molecular orbital (MO) calculations.<sup>[3]</sup> More recently the electronic structure of mercury amalgams has been studied by Liao using relativistic density functional calculations.<sup>[4]</sup> However, the aromatic character of the  $\text{Hg}_4^{6-}$  square units were not discussed. We have obtained a key insight into these questions through the recent discovery of aromaticity in purely metallic clusters.<sup>[5-7]</sup> We concluded that the square-planar  $\text{Hg}_4^{6-}$  ion is directly related to this finding and its structure and stability can in fact be attributed to aromaticity.

We recently reported experimental and theoretical evidence of aromaticity in an all-metal system, the  $\text{Al}_4^{2-}$  ion in a series of bimetallic and ionic clusters,  $\text{MAl}_4^-$  ( $\text{M} = \text{Li}, \text{Na}, \text{or Cu}$ ).<sup>[5]</sup> The  $\text{Al}_4^{2-}$  ion, both as an isolated species and in the bimetallic clusters, was found to be square-planar and to possess two delocalized  $\pi$  electrons, thus conforming to the structural criterion and the  $(4n+2)$  electron-counting rule for aromaticity. A simple electron counting for the  $\text{Hg}_4^{6-}$  ion immediately revealed that it is isosteric and isoelectronic with the  $\text{Al}_4^{2-}$  ion if we ignore the d electrons, which can be considered as part of the core. The analogy between  $\text{Hg}_4^{6-}$  and  $\text{Al}_4^{2-}$  suggests that the former may be aromatic as well, thus explaining its unique stability and structure.

To confirm this conjecture we performed ab initio calculations on an isolated  $\text{Hg}_4^{6-}$  ion, as well as on a model neutral  $\text{Na}_6\text{Hg}_4$  cluster and a known aromatic hydrocarbon analogue  $\text{C}_4\text{H}_4^{2+}$ .<sup>[8]</sup> While the isolated  $\text{Hg}_4^{6-}$  ion is not expected to be electronically stable as a gaseous species, it is a common practice in chemistry to divide a crystal into building blocks including highly charged anions.<sup>[9]</sup> For example, *nido*- $\text{Ti}_4^{8-}$  is accepted as a building block in  $\text{Na}_2\text{Ti}$ .<sup>[10]</sup> These charges are only formal charges and do not represent the real electron transfers in the crystals. But they are very important for the proper electron counting in chemistry and for understanding the structure and bonding of the solid building blocks. To test the effects of the counterions on the MOs and bonding on the  $\text{Hg}_4^{6-}$  building block in the  $\text{Na}_3\text{Hg}_2$  crystal, while keeping the same electron count, we further performed calculations on a model  $\text{Na}_6\text{Hg}_4$  cluster. In our calculations for  $\text{Na}_6\text{Hg}_4$  we used three theoretical methods: a hybrid of Hartree–Fock and density functional methods (B3LYP), a second-order Møller–Plesset perturbation method (MP2), and a complete active space multiconfigurational self-consistent field (CASSCF) calculation, including 103 740 electronic configurations.<sup>[11]</sup> In all calculations we used the triple-zeta valence basis sets and Stevens–Basch–Krauss effective core potentials (CEP-121G).

Since the purpose of the  $\text{Na}_6\text{Hg}_4$  calculations was not to determine its global minimum structure, but rather to examine how the presence of  $\text{Na}^+$  ions would influence the MOs and bonding in  $\text{Hg}_4^{6-}$ , it was important to ensure that the  $\text{Hg}_4$  unit in the  $\text{Na}_6\text{Hg}_4$  model has the same symmetry as in the naked  $\text{Hg}_4^{6-}$ . Our geometry optimization with the  $D_{4h}$  symmetry restriction led to the square  $\text{Hg}_4$  unit with  $R(\text{Hg}–\text{Hg}) = 2.9–3.1 \text{ \AA}$ , depending on the methods used. This calculated Hg–Hg bond length in the model  $\text{Na}_6\text{Hg}_4$  cluster is actually very close to the value of  $3.0 \text{ \AA}$  for the square  $\text{Hg}_4$  cluster observed in the  $\text{Na}_3\text{Hg}_2$  crystal.<sup>[2]</sup> This bond length is also very close to the optimized Hg–Hg bond ( $2.9–3.1$ ) in a linear Na–Hg–Hg–Na molecule. The covalent Hg–Hg bonding in the  $\text{Na}_2\text{Hg}_2$  molecule is undoubted and therefore the current result lends credence to the Hg–Hg covalent bonding in the model  $\text{Na}_6\text{Hg}_4$  cluster and the  $\text{Na}_3\text{Hg}_2$  crystal.

To examine the occupation of MOs we further performed CASSCF calculations on both  $\text{Hg}_4^{6-}$  and  $\text{Na}_6\text{Hg}_4$ . These calculations show that the Hartree–Fock wave function is predominate in the CASSCF expansion ( $C_{\text{HF}} = 0.892$ ), and the occupation numbers for the highest three Hartree–Fock orbitals are close to two,  $1.84 e (2a_{1g})$ ,  $1.87 e (1b_{2g})$ , and  $1.85 e$

( $1a_{2u}$ ). These results indicate that our one-electron MO picture can be used to understand the bonding of the  $\text{Hg}_4^{6-}$  cluster in the  $\text{Na}_3\text{Hg}_2$  amalgams. There is considerable charge transfer from Na to Hg and the model  $\text{Na}_6\text{Hg}_4$  cluster can be viewed formally as a  $\text{Hg}_4^{6-}$  cluster solvated by six  $\text{Na}^+$  ions. Figure 2a displays the seven valence MOs of the square-planar  $\text{Hg}_4^{6-}$  and Figure 2b displays those of the  $\text{Na}_6\text{Hg}_4$  model system. Comparing Figure 2b to 2a, we see clearly that the same set of occupied MOs is present in both cases. The counterions appear to have little effect on the occupied MOs of the  $\text{Hg}_4^{6-}$  building block; they alter slightly the MO ordering. The MOs shown in Figure 2a are identical to those of  $\text{Al}_4^{2-}$ , except for the slight difference in ordering among the top three MOs, which are bonding combinations of the valence p orbitals. The bottom four MOs are combinations of the valence s orbitals, which are essentially lone pairs in both  $\text{Hg}_4^{6-}$  and  $\text{Al}_4^{2-}$  with little s–p mixing. Based on the presence of the same two  $\pi$  electrons and its square-planar structure, we concluded that  $\text{Hg}_4^{6-}$  is indeed an aromatic all-metal cluster. To further confirm the aromatic nature of  $\text{Hg}_4^{6-}$ , we compare its MOs and bonding to those of an aromatic hydrocarbon  $\text{C}_4\text{H}_4^{2+}$ , as shown in Figure 2c.  $\text{C}_4\text{H}_4^{2+}$  has four more valence electrons than  $\text{Hg}_4^{6-}$ . We observed that except for the degenerate HOMO-1 in  $\text{C}_4\text{H}_4^{2+}$ , the remaining seven valence MOs of  $\text{C}_4\text{H}_4^{2+}$  are nearly identical to those of  $\text{Hg}_4^{6-}$ , with its characteristic aromatic  $\pi$  HOMO.

Mercury has a closed shell electron configuration ( $6s^2$ ) and neutral  $\text{Hg}_4$  is a van der Waals cluster (the bottom four MOs of Figure 2a or 2b provide no net bonding). The top three MOs of  $\text{Hg}_4^{6-}$  are completely bonding orbitals formed by perpendicular in-plane  $p_\sigma$  orbitals ( $1b_{2g}$ ), out-of-plane  $p_\pi$  orbitals ( $1a_{2u}$ ), and radial  $p_\sigma$  orbitals ( $2a_{1g}$ ), which are filled owing to electron donation from the alkali metals. On the other hand, the eight  $\sigma$  MOs in  $\text{C}_4\text{H}_4^{2+}$  actually describe four C–C bonds and four C–H bonds, consistent with the classical bonding picture for this saturated hydrocarbon molecule. Compared to  $\text{C}_4\text{H}_4^{2+}$ , the aromatic  $\text{Hg}_4^{6-}$  and  $\text{Al}_4^{2-}$  species are electron-deficient systems because we do not have four  $\sigma$  MOs to describe four two-electron two-center metal–metal bonds, as the square structure would suggest. Instead, in  $\text{Hg}_4^{6-}$  and  $\text{Al}_4^{2-}$ , we only have two four-center  $\sigma$  MOs describing the metal–metal bonding, in addition to the  $\pi$  MO. These multicenter  $\sigma$  bonds have been called  $\sigma$  aromaticity.<sup>[12]</sup> Hence, the structural and electronic stability of the square-planar  $\text{Hg}_4^{6-}$  and  $\text{Al}_4^{2-}$  should be attributed not only to  $\pi$  aromaticity due to the presence of the two  $\pi$  electrons in the

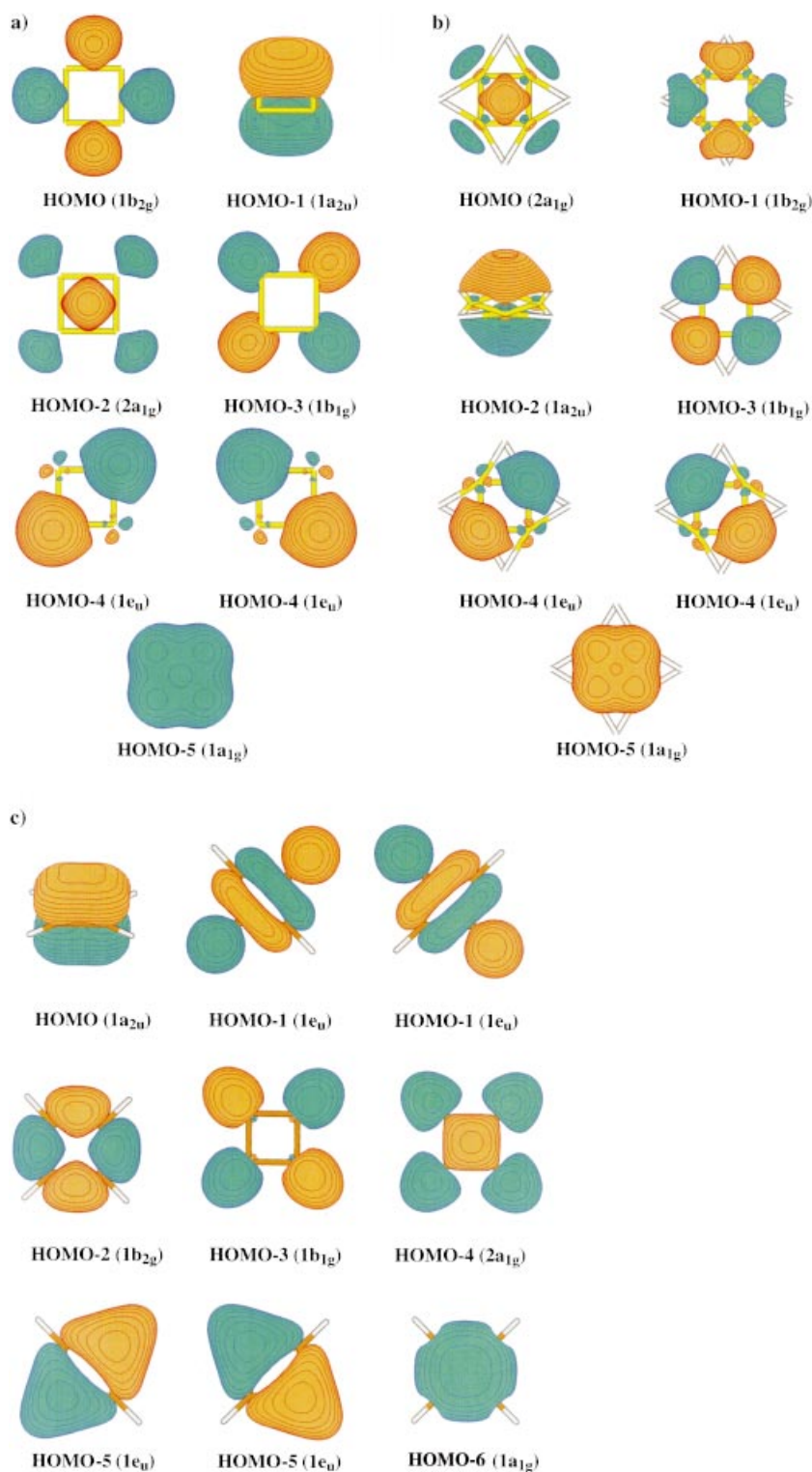


Figure 2. a) Molecular orbital pictures of square-planar  $\text{Hg}_4^{6-}$ , showing the highest occupied molecular orbital (HOMO,  $1b_{2g}$ ) and the aromatic  $\pi$  orbital (HOMO-1,  $1a_{2u}$ ) down to the fifth valence molecular orbital from the HOMO (HOMO-5,  $1a_{1g}$ ). Note that (HOMO-4) consists of a degenerate pair ( $1e_u$ ). b) Molecular orbital pictures of a model system,  $\text{Na}_6\text{Hg}_4$ , containing a square-planar  $\text{Hg}_4^{6-}$ . Note the similarity between the MOs in the two systems. c) Occupied molecular orbitals of the aromatic hydrocarbon,  $\text{C}_4\text{H}_4^{2+}$ .

$1a_{2u}$  orbital, but also to  $\sigma$  aromaticity due to the occupation of the two four-center  $\sigma$ -bonding orbitals,  $1b_{2g}$  and  $2a_{1g}$ .

Although an analogous aromatic  $\text{Ga}_4^{2-}$  structural unit has been recently suggested by us to exist in an organometallic

compound,<sup>[7]</sup> bulk solid materials containing the aromatic  $Al_4^{2-}$  ion have not been synthesized. Nevertheless, the current finding of aromaticity in the isoelectronic  $Hg_4^{6-}$  ion establishes a solid bridge between our gas-phase studies of aromatic clusters and bulk materials containing such species. It is a pleasant surprise that new insight can still be obtained into such an ancient and well-known material as the amalgam. It is possible that further development in the concept of aromaticity in metallic systems, through gas-phase clusters and model investigations, will lead to discovery of new classes of materials or new insights into the electronic structures and chemical bonding of complex materials.<sup>[13]</sup>

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## Allosteric Fluoride Anion Recognition by a Doubly Strapped Porphyrin\*\*


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The design of anion receptors with high selectivity is challenging.<sup>[1-6]</sup> Most systems developed use complimentary binding sites with hydrogen-bonding groups, quaternary ammonium centers, Lewis acids, and cationic metal ions.<sup>[1-6]</sup> It is often difficult to control selectivity and sensitivity among anions because of their wide range of geometries, low charge to radii ratios, sensitivities to pH, and high solvation energies.<sup>[1a]</sup> Receptors for the smallest anion, fluoride, are of special importance for monitoring the fluoride metabolism in nature, the analysis of drinking water,<sup>[7]</sup> and the detection of chemical warfare agents.<sup>[8]</sup> We report herein a highly selective allosteric fluoride recognition system<sup>[9]</sup> by a doubly strapped porphyrin<sup>[10]</sup> **1** that contains two small hydrogen-bonding cavities that are not able to bind larger anions. It is also demonstrated that a conducting polymer based upon **1** displays both electrochemical and conductivity responses to fluoride ions and no response to chloride ions.

Compound **1** was prepared by condensation of strap moieties **2** with  $\alpha\beta\alpha\beta$ -tetrakis(2-aminophenyl)porphyrin **3** under high-dilution conditions (Scheme 1).<sup>[11, 12]</sup> By comparison with related strapped porphyrins containing four linkers, we estimate that the cofacial distance between the phenyls in the straps and the porphyrin is between 3–4 Å.<sup>[12]</sup> This distance creates a small pocket that is suitable for fluoride-ion binding (the radius of a fluoride anion in an octahedral environment is 1.19 Å<sup>[1a, 13]</sup>), but should exclude larger anions. The proximate location of the strap moieties to the porphyrin in **1** is confirmed by NMR spectroscopy in  $CDCl_3$ , which revealed significant upfield shifts of specific signals. The phenyl protons appear at  $\delta = 4.24$  which is approximately 2.96 ppm upfield of the dicarboxylate strap precursor. The amide N–H protons resonate at  $\delta = 7.68$ , which is comparable to closely related phenyl-capped porphyrins which have an additional methylene between the amide carbonyl and the phenyl ether of the strap.<sup>[14]</sup> These values are very different from unstrapped porphyrins where the same resonances generally appear between  $\delta = 5$  and 6.5<sup>[10, 15]</sup> and suggest that the amide N–H protons are constrained against the porphyrin core. The Soret band of **1** (5.00  $\mu m$ ) in dichloromethane is split into two bands of equal intensity at 404 and 423 nm, which indicate that the strap presents a significant perturbation. This

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