

# MnB<sub>6</sub><sup>−</sup>: An Open-Shell Metallaboron Analog of 3d Metallabenzenes

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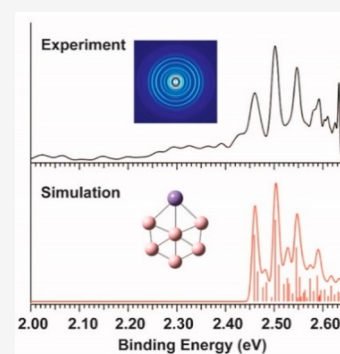
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**ABSTRACT:** We report a study of the structure and bonding of a transition-metal-doped boron cluster, MnB<sub>6</sub><sup>−</sup>, using high-resolution photoelectron imaging and quantum chemical calculations. Vibrationally resolved photoelectron spectra indicate a significant geometry change between the anionic and neutral ground states of MnB<sub>6</sub>. The electron affinity of MnB<sub>6</sub> is measured to be 2.4591(5) eV, and vibrational frequencies for five of its vibrational modes were determined. The experimental data are combined with theoretical calculations to determine the structure and bonding of MnB<sub>6</sub><sup>−</sup>, which is found to be planar with a B-centered hexagonal structure (C<sub>2v</sub> symmetry) and a quintet spin state (<sup>5</sup>A<sub>2</sub>). Nuclear-independent chemical shift calculations indicate that MnB<sub>6</sub><sup>−</sup> is aromatic. Molecular orbital analyses reveal that MnB<sub>6</sub><sup>−</sup> contains three  $\pi$  orbitals, one of which is singly occupied. Hence, MnB<sub>6</sub><sup>−</sup> can be considered as an open-shell metallaboron analog of 3d metallabenzenes.



## INTRODUCTION

Boron is known to form three-dimensional (3D) polyhedral structures with delocalized bonds in many borane compounds because of its electron deficiency.<sup>1</sup> The structures and bonding of bare boron clusters have been extensively investigated by joint experimental and computational studies in recent years.<sup>2–8</sup> The structures of small boron clusters have been found to be predominantly planar with delocalized bonding, giving rise to both  $\sigma$  and  $\pi$  aromaticity.<sup>9–15</sup> It was further found that the planarity and aromaticity of small boron clusters obey the Hückel rule. Transition-metal-doped boron clusters have also been studied by joint photoelectron spectroscopy (PES) and computational chemistry.<sup>16–26</sup> In addition, there has been a variety of computational studies on transition-metal-doped boron clusters.<sup>27–35</sup> Notably, a class of aromatic borometallic molecular wheels (M@B<sub>n</sub><sup>−</sup>) ( $n = 8–10$ ) have been experimentally discovered.<sup>16,17</sup> While Re@B<sub>8</sub><sup>−</sup> and Re@B<sub>9</sub><sup>−</sup> have been found to belong to the class of borometallic molecular wheels,<sup>36</sup> a recent study showed that ReB<sub>6</sub><sup>−</sup> is planar with the Re atom being on the periphery of the cluster.<sup>37</sup> Chemical bonding analyses revealed that ReB<sub>6</sub><sup>−</sup> is aromatic with six  $\pi$  electrons, whereas AlB<sub>6</sub><sup>−</sup> with four  $\pi$  electrons is antiaromatic with an out-of-plane distortion. The bonding in ReB<sub>6</sub><sup>−</sup> was further shown to be similar to that in rhenabenzene. Thus, ReB<sub>6</sub><sup>−</sup> is the first aromatic metallaboron analog of metallabenzenes.<sup>37</sup>

Metallabenzenes are a class of organometallic compounds in which one CH group in benzene is substituted by a metal atom.<sup>38–41</sup> There have been increasing studies on metallabenzenes due to their interesting chemical and structural

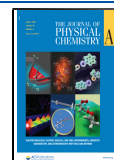
properties. While there have been numerous reports on the syntheses of metallabenzenes with various 5d transition metals,<sup>39–47</sup> as well as the 4d ruthenium,<sup>48,49</sup> it has been challenging to synthesize metallabenzenes with the first-row and other second-row transition metals.<sup>40,50</sup> Very recently, a tetralithio metalla-aromatic compound has been synthesized with two five-membered manganacycles spiro-fused by a high-spin Mn atom.<sup>51</sup> However, the synthesis of monocyclic metallabenzenes with any 3d transition metals has remained elusive, although there has been a previous theoretical study.<sup>52</sup>

In this article, we address the possibility of producing a 3d metallaboron analog of metallabenzenes. We report an investigation of a 3d transition-metal metallaboron cluster, MnB<sub>6</sub><sup>−</sup>, produced by laser vaporization of a mixed Mn/B target. The selection of Mn was inspired by the recent synthesis of manganacycles<sup>51</sup> and the ReB<sub>6</sub><sup>−</sup> metallaboron analog of 5d metallabenzenes.<sup>37</sup> High-resolution photoelectron imaging (PEI) in conjunction with computational chemistry has shown that MnB<sub>6</sub><sup>−</sup> has a planar structure with the Mn atom on the periphery of the cluster, similar to that of ReB<sub>6</sub><sup>−</sup>. However, the MnB<sub>6</sub><sup>−</sup> cluster has a high spin ground state (<sup>5</sup>A<sub>2</sub>) and displays relatively weak Mn–B interactions. Nuclear shielding calculations using the nuclear-independent

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chemical shift (NICS) reveal aromatic characters in  $\text{MnB}_6^-$ , which can be considered as a metallaboron analog of metallabenzenes with a 3d transition metal atom.

## EXPERIMENTAL AND THEORETICAL METHODS

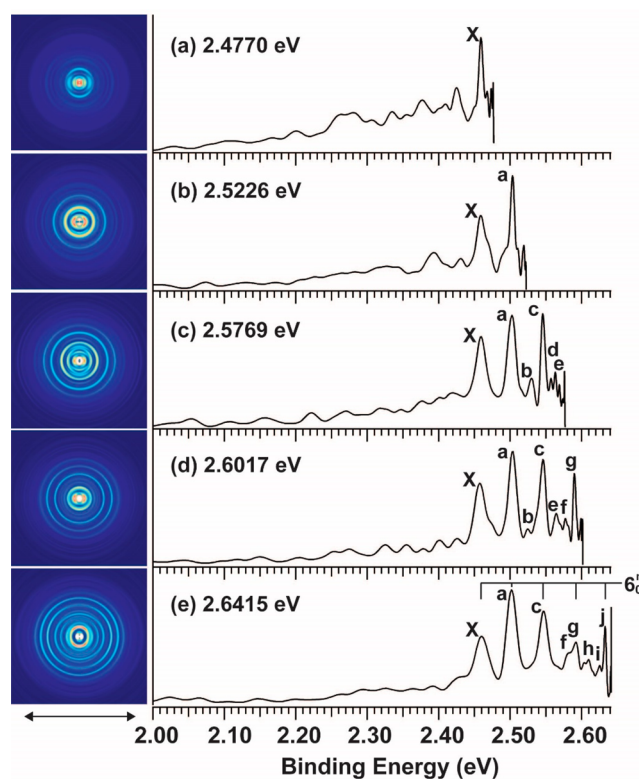
**High-Resolution Photoelectron Imaging.** The experiment was carried out using a high-resolution PEI apparatus with a laser vaporization cluster source, which has been described in detail previously.<sup>53</sup> Briefly, we produced  $\text{MnB}_6^-$  by focusing the second harmonic of a Nd:YAG laser onto a disk target, which was made of a mixture of isotopically enriched  $^{10}\text{B}$ , Mn, and Ag powders (Ag served as a binder). The laser-induced plasma was quenched by a helium carrier gas seeded with 10% argon, initiating nucleation. The nascent clusters were entrained in the carrier gas and underwent a supersonic expansion, producing a cold cluster beam. Anionic clusters were extracted perpendicularly into a time-of-flight mass spectrometer. The  $\text{MnB}_6^-$  cluster of current interest was mass-selected into the interaction zone of the velocity-map imaging (VMI) system.

A second Nd:YAG laser or a Deyang Tech dye laser pumped by a Nd:YAG laser was used to photodetach the size-selected  $\text{MnB}_6^-$  clusters. The photoelectrons were extracted from the interaction zone and focused onto a set of microchannel plates coupled with a phosphor screen and a charge-coupled device camera. Each experiment at a given photon energy required about 100 000 to 200 000 laser shots to achieve reasonable signal-to-noise ratios. The VMI lens was calibrated using the photoelectron images of  $\text{Au}^-$  and  $\text{Bi}^-$  at various photon energies. The photoelectron images were analyzed using the maximum entropy method (MEVIR and MEVELER).<sup>54</sup> The typical energy resolution of the VMI system was  $\sim 0.6\%$  for high kinetic energy electrons and could be as good as  $1.2\text{ cm}^{-1}$  for low kinetic energy electrons.<sup>53</sup>

**Theoretical Methods.** Structural searches for  $\text{MnB}_6^-$  were performed by direct Monte Carlo sampling with local optimization. Around 2000 different structures with various spin multiplicities (singlet, triplet, quintet, and septet) were randomly generated and optimized at the PBE/6-31G\* level of theory.<sup>55–57</sup> Low-lying isomers within 1 eV of the global minima were further optimized at the B3LYP/6-311+G\* level of theory.<sup>58</sup> The adiabatic detachment energy (ADE) for the ground state transition was calculated as the energy difference between the optimized anion and its corresponding neutral. Franck–Condon simulation was performed using ezSpec-trum.<sup>59</sup> NICS calculations were carried out at the B3LYP/6-311+G\* level of theory.<sup>60</sup> All calculations were done using Gaussian 09.<sup>61</sup>

## RESULTS

**Experimental Results.** Figure 1 shows the PE images and spectra of  $\text{MnB}_6^-$  at five different photon energies. The first sharp peak labeled as X denotes the 0–0 transition, corresponding to the transition from the anion ground state to that of the neutral. Peak X defines the electron affinity (EA) of  $\text{MnB}_6^-$  to be  $2.4591 \pm 0.0005\text{ eV}$ . Extensive vibrational structures were resolved, indicating a significant geometry change between the ground states of the anion and the neutral system. The resolved vibrational peaks are labeled as a–j. The binding energies of the observed peaks and their assignments are summarized in Table 1. The weak signals on the lower binding energy side of peak X were likely due to vibrational hot



**Figure 1.** Photoelectron images and spectra of  $\text{MnB}_6^-$  at five different photon energies. The double arrows below the images denote the laser polarization.

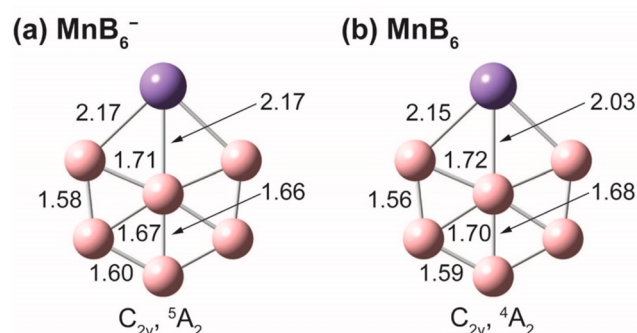
**Table 1.** Measured Binding Energies (BE), Energy Shifts Relative to the 0–0 Transition, and Assignments of the Observed Vibrational Peaks in the PE Spectra of  $\text{MnB}_6^-$  and Comparison with the Theoretical Frequencies Computed at the B3LYP/6-311+G\* Level of Theory

peak	experimental			theoretical <sup>a</sup>
	BE (eV)	assignment	energy shift ( $\text{cm}^{-1}$ )	frequency ( $\text{cm}^{-1}$ )
X	2.4591(5)	$0_0^0$		
a	2.5035(9)	$6_0^1$	358(9)	348
b	2.5292(23)	$5_0^1$	565(19)	578
c	2.5461(10)	$6_0^2$	702(9)	
d	2.5569(16)	$4_0^1$	789(14)	733
e	2.5633(13)	$3_0^1$	840(11)	818
f	2.5777(18)	$5_0^1 6_0^1$	957(15)	
g	2.5901(13)	$6_0^3$	1057(11)	
h	2.6099(18)	$1_0^1$	1216(15)	1225
i	2.6259(35)	$4_0^1 5_0^1$	1345(29)	
j	2.6330(12)	$6_0^4$	1403(10)	

<sup>a</sup>The calculated ADE for  $\text{MnB}_6^-$  is 2.43 eV at the B3LYP/6-311+G\* level of theory.

bands or metastable electronic excited states of the anion (*vide infra*). The apparent enhancement of these weak signals in Figure 1a was due to the fact that the detachment cross section for peak X was reduced at the lower photon energy.

**Computational Results.** The lowest energy structure of  $\text{MnB}_6^-$  was found to be perfectly planar with  $C_{2v}$  symmetry (Figure 2). The structures and relative energies of other low-lying isomers within 1 eV of the global minimum are given in Figure S1. The ground state of  $\text{MnB}_6^-$  consists of a B-centered six-membered ring with Mn being on the periphery, similar to



**Figure 2.** Optimized structures of (a)  $\text{MnB}_6^-$  and (b)  $\text{MnB}_6$ . The point group symmetries and electronic states are also given. The bond lengths are in Å.

the structure of  $\text{ReB}_6^-$ .<sup>37</sup> While  $\text{ReB}_6^-$  is closed shell,  $\text{MnB}_6^-$  is found to have a high spin state ( $^5\text{A}_2$ ) with an electron configuration of  $1a_1^2 1b_2^2 2a_1^2 2b_2^2 1b_1^2 3a_1^2 1a_2^2 3b_2^2 4a_1^2 4b_2^2 - 5a_1^2 6a_1^2 7a_1^2 2b_1^2 5b_2^1$ . The occupied valence molecular orbitals (MO) are shown in Figure 3, where the  $\beta$  orbitals are also displayed for the doubly occupied  $\pi$  MOs since they show some differences from the corresponding  $\alpha$  orbitals. The neutral ground state  $^4\text{A}_2$  is obtained by detaching an  $\alpha$  electron from the highest occupied MO (HOMO). As shown in Figure 2, there is a geometry change between the anion and the neutral ground state: the distance between the Mn atom and the central B atom is significantly shorter in the neutral. This is consistent with the removal of an electron from the HOMO, which exhibits an antibonding interaction between Mn and the central B atom.

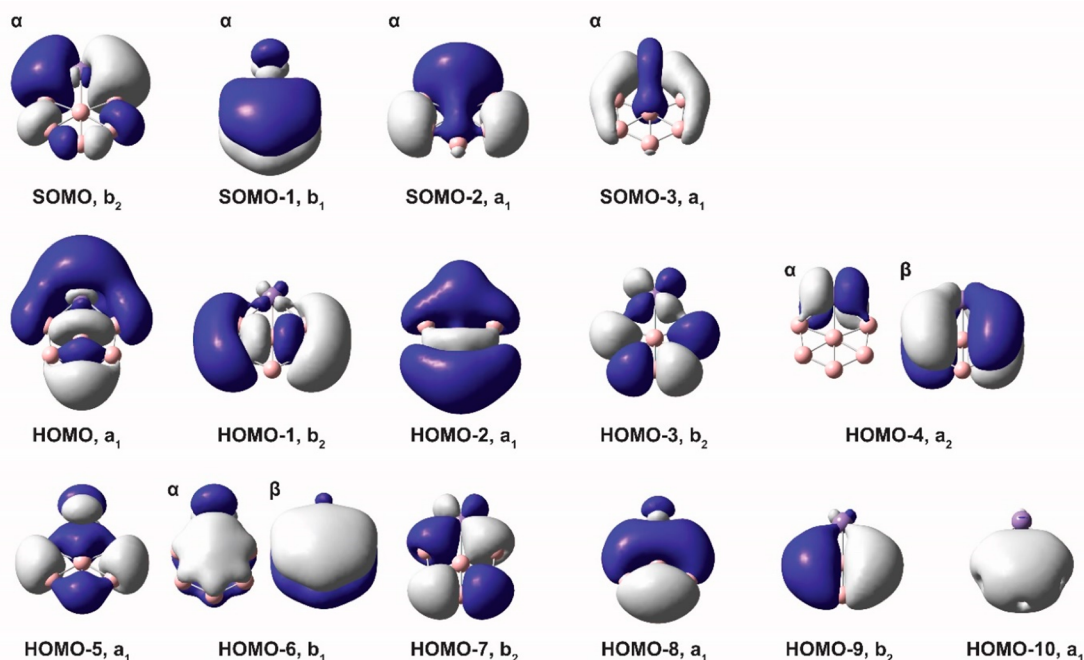
## DISCUSSION

**Comparison of the Photoelectron Spectra of  $\text{MnB}_6^-$  with the Theoretical Results.** The ADE of  $\text{MnB}_6^-$  is

calculated to be 2.43 eV at the B3LYP/6-311+G\* level of theory, which agrees well with the experimental value of 2.4591 eV. The computed vibrational frequencies also show excellent agreement with the experimental measurements (Table 1). Five totally symmetric vibrational modes,  $\nu_6$ ,  $\nu_5$ ,  $\nu_4$ ,  $\nu_3$ , and  $\nu_1$ , were observed and assigned to peaks *a*, *b*, *d*, *e*, and *h*, respectively. Their displacement vectors are given in Figure S2. As shown in the experimental spectra, the most Franck–Condon active mode is  $\nu_6$ , which corresponds to the stretching motion of the Mn atom. This observation is consistent with the fact that electron detached is from the HOMO of  $\text{MnB}_6^-$  (Figure 3). A Franck–Condon simulation was done,<sup>59</sup> as compared with the 2.6415 eV spectrum in Figure 4, which reproduced all the major vibrational features of the experimental spectra. The overall excellent agreement between the experimental and theoretical results confirms unequivocally the geometric and electronic structure of  $\text{MnB}_6^-$ .

The Franck–Condon simulation suggested that there were negligible vibrational hot band contributions to the weak signals below the 0–0 transition. These weak signals seemed to span a wide energy range down to  $\sim 2.2$  eV. A very likely source of these weak signals was due to the population of low-lying electronic excited states of the anion. As shown in Figure S1, the  $^7\text{A}_2$  excited state was computed to be  $\sim 0.22$  eV above the  $^5\text{A}_2$  ground state. Despite the fact that  $\text{MnB}_6^-$  was relatively cold vibrationally, the high-spin  $^7\text{A}_2$  state, which was expected to be metastable, might not be in equilibrium with the vibrational temperature and could be populated in small quantity in the cluster beam.

**Chemical Bonding and Aromaticity in  $\text{MnB}_6^-$ .** Although  $\text{MnB}_6^-$  has a similar geometric structure as the recently reported  $\text{ReB}_6^-$ , their electronic structures are different. In  $\text{MnB}_6^-$ , due to its large spin multiplicity, the shapes of the  $\alpha$  and  $\beta$  orbitals are not necessarily the same. Figure 5 displays the occupied  $\pi$ -MOs of  $\text{MnB}_6^-$ . The  $\beta$  orbitals are found to be more diffuse relative to their  $\alpha$



**Figure 3.** Occupied valence molecular orbitals of  $\text{MnB}_6^-$ . The  $\alpha$  orbitals are displayed. The  $\beta$  orbitals for the doubly occupied  $\pi$  MOs are also shown.



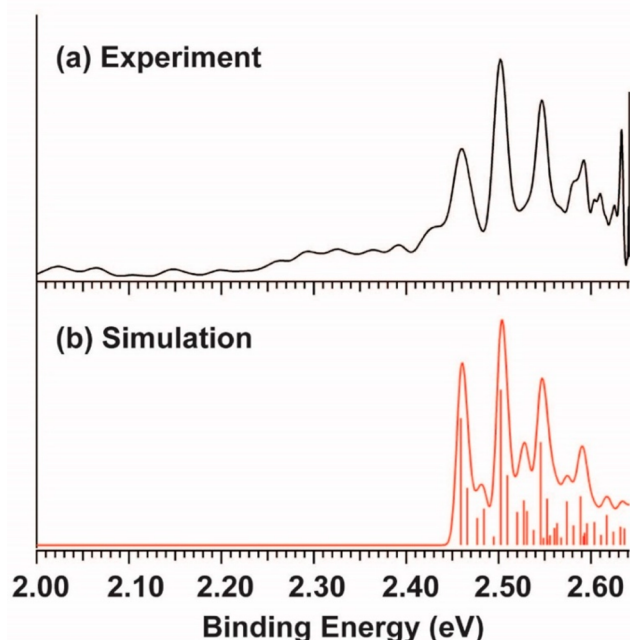


Figure 4. Franck–Condon simulation of  $\text{MnB}_6^-$  at 100 K.

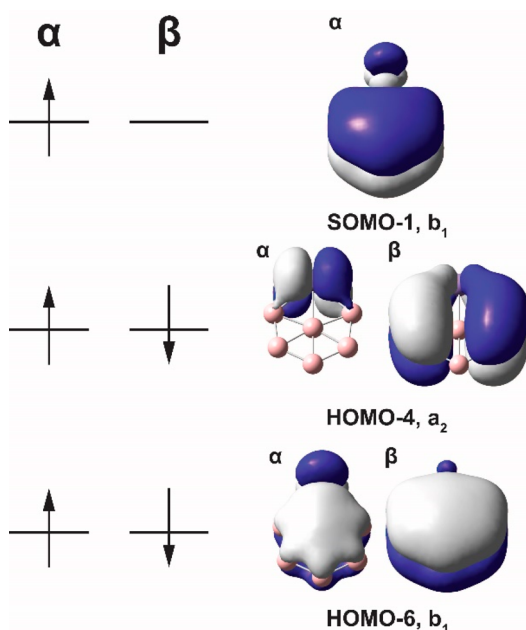


Figure 5. Occupied  $\pi$ -MOs of  $\text{MnB}_6^-$  at isovalue = 0.02 e/bohr.<sup>3</sup>

counterparts. In particular, the  $\alpha$  electron in HOMO-4 is essentially localized in the  $3d_{xz}$  orbital. On the other hand, the  $\beta$  orbital of HOMO-4 is mainly delocalized on the B atoms with a small contribution from the  $3d_{xz}$  orbital. The  $\pi$  MOs of SOMO-1 and HOMO-6 are also mainly delocalized on the  $\text{B}_6$  motif with relatively small contributions from the Mn 3d orbitals. In comparison to the 5d element Re, the 3d orbitals in Mn are spatially more contracted and are expected to have smaller overlap with the 2p orbitals of boron. The weaker bonding capacity of Mn with B can be gleaned from the significantly larger Mn–B bond distances (2.17 Å) in  $\text{MnB}_6^-$  (Figure 2a), compared with the Re–B bond distances (1.98–2.05 Å) in  $\text{ReB}_6^-$ .<sup>37</sup> There have been two previous studies on

lanthanide hexaboride clusters ( $\text{SmB}_6^-$  and  $\text{CeB}_6^-$ ),<sup>62,63</sup> which have similar planar structures as  $\text{MnB}_6^-$ . The Ln–B bond lengths are even longer, indicating even weaker interactions between the Ln atoms and boron in these clusters. In fact, the  $\text{SmB}_6^-$  cluster was characterized to consist of a  $\text{B}_6^{2-}$  unit with double antiaromaticity.<sup>62</sup>

Despite the fact that  $\text{MnB}_6^-$  possesses only five  $\pi$  electrons with one of the  $\pi$  MOs singly occupied, its structural similarity with  $\text{ReB}_6^-$  suggested that it might be aromatic, which was evaluated using the NICS index.<sup>60,64</sup> It should be noted that the open-shell nature of  $\text{MnB}_6^-$  makes it difficult to apply the Hückel rule. The NICS(0)<sub>zz</sub> and NICS(1)<sub>zz</sub> values of  $\text{MnB}_6^-$  were calculated at different sites (Figure S3) and are compared with those of benzene in Table 2. All of the the NICS(0)<sub>zz</sub> and

Table 2. NICS(0)<sub>zz</sub> and NICS(1)<sub>zz</sub> Values (in ppm) of Benzene and  $\text{MnB}_6^-$  at the B3LYP/6-311+G\* Level of Theory<sup>a</sup>

NICS(0) <sub>zz</sub>	site X	site A	site B	site C
$\text{C}_6\text{H}_6$	−14.07	—	—	—
$\text{MnB}_6^-$	−140.79	−31.40	−35.07	−32.34
NICS(1) <sub>zz</sub>	site X	site A	site B	site C
$\text{C}_6\text{H}_6$	−29.06	—	—	—
$\text{MnB}_6^-$	−8.87	−4.24	−2.72	−5.40

<sup>a</sup>The site positions are shown in Figure S3.

NICS(1)<sub>zz</sub> values are negative, indicating that  $\text{MnB}_6^-$  indeed has aromatic characters. Because the three  $\pi$  MOs are not completely filled,  $\text{MnB}_6^-$  can be considered as an open-shell weak aromatic system or a 3d metallaboron analog of metallabenzenes. The weak bonding between Mn and the  $\text{B}_6$  motif in  $\text{MnB}_6^-$  suggests that the 3d transition metals should also have weaker interactions with carbon in metallabenzenes, consistent with the challenges to synthesize such compounds.

## CONCLUSIONS

In conclusion, we have produced and studied a 3d metallaboron cluster ( $\text{MnB}_6^-$ ) using high-resolution photoelectron imaging and quantum chemical calculations. Vibrationally resolved photoelectron spectra were obtained and used to determine the structure of  $\text{MnB}_6^-$  by comparing them with the theoretical results.  $\text{MnB}_6^-$  was found to have a planar B-centered hexagonal structure with the Mn atom on the periphery and a high-spin ground state. Molecular orbital analyses and NICS calculations suggest that  $\text{MnB}_6^-$  is aromatic and can be viewed as an open-shell metallaboron analog of 3d metallabenzenes.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jpca.0c00949>.

The relative energies of low-lying isomers, the computed vibrational frequencies, and the site positions for the NICS calculations of  $\text{MnB}_6^-$  (PDF)

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## Notes

The authors declare no competing financial interest.

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