Anisotropic flow in ALICE at the LHC Ante Bilandzic (for the ALICE collaboration)

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Introduction



Anisotropies in momentum space S. Voloshin and Y. Zhang (1996) =>

$$E\frac{d^3N}{d^3\vec{p}} = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} 2v_n \cos\left(n\left(\phi - \Psi_{\rm RP}\right)\right) \right)$$
$$v_n = \left\langle \cos(n(\phi - \Psi_{\rm RP})) \right\rangle$$

- Coordinate => momentum space
- In case of fluctuations $\Psi_{\text{RP}} \Rightarrow \Psi_n$
 - Reaction plane is a plane spanned by the impact parameter and beam axis
- Harmonics v_n quantify anisotropic flow
 - v_1 is directed flow, v_2 is elliptic flow, v_3 is triangular flow, etc.



Why flow?



- Goal: expect to see/study the QGP (deconfined quark-gluon matter)
- Test: whether hydrodynamic (liquid) description can describe the data
 - If yes, can we extract the transport properties of this matter (e.g. viscosity)?
 - Other items of interest: Equation of state (EoS), energy loss...
- For detailed description: need realistic time evolution => initial conditions (fluctuations), deconfined phase (hydro), hadronization
- Anisotropic flow is sensitive to the system evolution (in particular it probes the properties of the created matter such as shear viscosity)
 - Perfect liquid
 shear viscosity zero





 Shear viscosity characterizes quantitatively the resistance of the liquid to displacement of its layers



Analysis outline



- Data:
 - 2010 + 2011 => Pb-Pb events at 2.76 TeV
 - Acceptance => $|\eta| < 0.8$ (Time Projection Chamber)
 - $|\eta| < 5.1$ (Forward Multiplicity Detector)
- Charged particle tracking:
 - Time Projection Chamber
 - Inner Tracking System
- Centrality determination
 - VZERO detectors
 - VZERO-A => 2.8 < η < 5.1
 - VZERO-C => -3.7 < η < -1.7
- Systematic uncertainties:
 - Nonflow
 - Centrality determination
 - Inefficiencies in detectors azimuthal acceptance
 - Variation of track quality cuts
 - Secondaries in the material











- Elliptic flow is geometrical quantity => need to classify all events in terms of initial geometry
- Another geometrical quantity available: Multiplicity
 - In central collisions more nucleons within nuclei interact than in the peripheral collisions => more particles are produced in the central collisions than in the peripheral
- Glauber model: Quantitative description of multiplicity distribution, centrality classes of events
 FALICE PD-PD at VSw = 2.76 TeV
 - Most central => Centrality class 0-5%
 - Peripheral => Centrality class 70-80%







How do we measure anisotropic flow?





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Azimuthal correlations

• Definition of flow harmonics :

$$v_n = \langle \cos(n(\varphi - \Psi_{\rm RP})) \rangle$$

• Two particle azimuthal correlations (double brackets indicate an event sample and particles in each event averages):

$$\left\langle \left\langle e^{in(\phi_1 - \phi_2)} \right\rangle \right\rangle = \left\langle \left\langle e^{in(\phi_1 - \Psi_{\rm RP} - (\phi_2 - \Psi_{\rm RP}))} \right\rangle \right\rangle \\ = \left\langle \left\langle e^{in(\phi_1 - \Psi_{\rm RP})} \right\rangle \left\langle e^{-in(\phi_2 - \Psi_{\rm RP})} \right\rangle \right\rangle = \left\langle v_n^2 \right\rangle$$

- Systematic biases:
 - Few-particle correlations unrelated to the initial geometry (nonflow)
 - Multiplicity fluctuations
- Suppress nonflow with multi-particle cumulants

N. Borghini, P. M. Dinh and J.-Y. Ollitrault, "Flow analysis from multiparticle azimuthal correlations," PRC 64 (2001) 054901 Ryogo Kubo, "Generalized Cumulant Expansion Method"



2-particle cumulant



• The most general decomposition of 2-particle correlation

 $\langle X_1 X_2 \rangle = \langle X_1 \rangle \langle X_2 \rangle + \langle X_1 X_2 \rangle_c$

 X_1 and X_2 are two observables (e.g. particle density)

• 2-particle cumulant (non-factorisable 2-particle correlation):

 $\langle X_1 X_2 \rangle_c = \langle X_1 X_2 \rangle - \langle X_1 \rangle \langle X_2 \rangle$



3-particle cumulant



• The most general decomposition of 3-particle correlation:

$$\begin{array}{lcl} X_{1}X_{2}X_{3}\rangle &=& \langle X_{1}\rangle \langle X_{2}\rangle \langle X_{3}\rangle \\ &+& \langle X_{1}X_{2}\rangle_{c} \langle X_{3}\rangle + \langle X_{1}X_{3}\rangle_{c} \langle X_{2}\rangle + \langle X_{2}X_{3}\rangle_{c} \langle X_{1}\rangle \\ &+& \langle X_{1}X_{2}X_{3}\rangle_{c} \end{array}$$

• Using expression for 2-particle cumulants obtain:

$$\begin{array}{rcl} X_1 X_2 X_3 \rangle_c &=& \langle X_1 X_2 X_3 \rangle \\ && - & \langle X_1 X_2 \rangle \langle X_3 \rangle - \langle X_1 X_3 \rangle \langle X_2 \rangle - \langle X_2 X_3 \rangle \langle X_1 \rangle \\ && + & 2 \langle X_1 \rangle \langle X_2 \rangle \langle X_3 \rangle \end{array}$$

=> Recursively one can obtain cumulants for any number of observables



Cumulants in flow analysis



• Observables in the context of anisotropic flow analysis (Ollitrault *et al*)

$$egin{array}{ll} X_1 \equiv e^{in\phi_1}, & X_2 \equiv e^{in\phi_2} \ X_3 \equiv e^{-in\phi_3}, & X_4 \equiv e^{-in\phi_4} \end{array}$$

• An event average of single particle observable vanish:

$$\langle \langle X_i \rangle \rangle = \langle \langle e^{\pm i n \varphi} \rangle \rangle = 0$$
 for all n

Correlate only distinct particles (exclude self correlation):

$$\langle 2 \rangle \equiv \langle \cos n(\phi_1 - \phi_2) \rangle , \qquad \phi_1 \neq \phi_2 \langle 4 \rangle \equiv \langle \cos n(\phi_1 + \phi_2 - \phi_3 - \phi_4) \rangle , \qquad \phi_1 \neq \phi_2 \neq \phi_3 \neq \phi_4$$



Flow vector and cumulants



• Q_n-vector (or flow vector) evaluated for harmonic *n*, and for event with multiplicity *M*:

$$Q_n = \sum_{i=1}^M e^{in\phi_i}$$

 Analytical expressions for multi-particle azimuthal correlations in terms of *Q*-vectors

$$\langle 2 \rangle = \frac{|Q_n|^2 - M}{M(M-1)}$$

$$\langle 4 \rangle = \frac{|Q_n|^4 + |Q_{2n}|^2 - 2 \cdot \operatorname{Re} \left[Q_{2n} Q_n^* Q_n^*\right] - 4(M-2) \cdot |Q_n|^2}{M(M-1)(M-2)(M-3)}$$

$$+ \frac{2}{(M-1)(M-2)}$$



Flow cumulants



 Cumulants expressed in terms of azimuthal correlations, defined in terms of *Q*-vectors:

$$c_{n}\{2\} = \langle \langle 2 \rangle \rangle$$

$$c_{n}\{4\} = \langle \langle 4 \rangle \rangle - 2 \langle \langle 2 \rangle \rangle^{2}$$

$$c_{n}\{6\} = \langle \langle 6 \rangle \rangle - 9 \langle \langle 2 \rangle \rangle \langle \langle 4 \rangle \rangle + 12 \langle \langle 2 \rangle \rangle^{3}$$

$$c_{n}\{8\} = \langle \langle 8 \rangle \rangle - 16 \langle \langle 6 \rangle \rangle \langle \langle 2 \rangle \rangle - 18 \langle \langle 4 \rangle \rangle^{2}$$

$$+ 144 \langle \langle 4 \rangle \rangle \langle \langle 2 \rangle \rangle^{2} - 144 \langle \langle 2 \rangle \rangle^{4}$$

 If all correlations are expressed analytically in terms of *Q*-vectors => *Q*-cumulants (QC)

R. Snellings, S. Voloshin, AB: *"Flow analysis with cumulants: Direct calculations",* PRC 83, 044913 (2011)



Flow contribution to cumulants



 In the absence of flow fluctuations, flow harmonics enter in a power of the cumulant order:

$$c_{n}\{2\} = v_{n}^{2}$$

$$c_{n}\{4\} = -v_{n}^{4}$$

$$c_{n}\{6\} = 4v_{n}^{6}$$

$$c_{n}\{8\} = -33v_{n}^{8}$$

These relations hold for any harmonic

- Each of the equations above gives an independent observable for v_n => v_n{2}, v_n{4}, v_n{6}, etc.
 - Important for flow fluctuations studies (discussed later)





Results





A central heavy-ion collision as seen by ALICE







First results on the elliptic flow at the LHC





=> Elliptic flow increases by ~ 30% when compared to RHIC energies

Phys. Rev. Lett. 105, 252302 (2010)

Cited by now ~ 300 times! (most cited LHC heavy-ion paper according to Spires)





Two and multi-particle cumulants in ALICE





=> The difference between v_2 {2} with and without eta gap is driven by the contribution from nonflow

=> The difference between 2- and multi-particle estimates is due to fluctuations in the initial geometry



Flow fluctuations



 Event-by-event fluctuations in the positions of participating nucleons => non-zero odd harmonics:

$$\Psi_{\text{RP}} \Rightarrow \Psi_n$$

 $v_n = \langle \cos(n(\varphi - \Psi_n)) \rangle$

- Each harmonic v_n has its own symmetry plane Ψ_n
- Experimental consequences of e-b-e flow fluctuations:
 \$\langle v_n^k \rangle\$ is not the same as \$\langle v_n \rangle^k\$







Charged particle v₃



- Phys.Rev.Lett. 107 (2011) 032301
- Nonzero v₃ develops along its own symmetry plane
- Symmetry plane of v_2 shows no correlation with the plane of v_3



Discovery



v_n vs. transverse momentum



Phys.Lett.B 719 (2013) 18-28



* Central collisions: all harmonics are similar

- * v_3 is almost independent of centrality, dominant harmonic for more central events very soon i.e. at moderate p_T values.
- * $v_4{\Psi_4} > v_4{\Psi_2} =>$ flow fluctuations

* Non-vanishing v_2 at high- p_T ($p_T > 8$ GeV/c) => path length dependence of energy loss of high- p_T partons traversing the non-isotropic medium



Estimate of v₂ fluctuations



Phys.Lett.B 719 (2013) 18-28



=> Fluctuations are similar up to $p_T \sim 6$ GeV, with the exception of most central events





=> Fluctuations do not change significantly with rapidity (A. Hansen QM12, <u>arXiv:1210.7095</u>)



p.d.f. of flow fluctuations



- If for the 1st ($\langle v \rangle$) and 2nd (σ_v) moments ($\sigma_v / \langle v \rangle$)² «1 is satisfied, then all multi-particle cumulants for any p.d.f. are the same
- Odd harmonics originate from fluctuations => $(\sigma_v / \langle v \rangle)^2 \ll 1$ is never satisfied
- Bessel-Gaussian p.d.f: All higher moments degenerated $v_n{4} = v_n{6} = v_n{8} = \dots$

$$f(v) = \frac{v}{b^2} \exp\left(-\frac{v^2 + a^2}{2b^2}\right) I_0\left(\frac{va}{b^2}\right)$$
$$v\{2\} = \sqrt{a^2 + 2b^2}$$
$$4, 6, \ldots\} = a$$

Voloshin et al: PLB 659, 537 (2008)

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v₁, v₂, and v₃ from multi-particle cumulants

• Established experimentally that $v_n{4} \sim v_n{6} \Rightarrow p.d.f.$ of e-b-e flow fluctuations must have non-negligible 3rd/higher moments (when compared to the 1st/2nd moment)



=> Bessel-Gaussian function is an example of p.d.f. with $v_n{4} = v_n{6}$



Symmetry plane correlation



 Observable to determine correlation between different symmetry planes:

$$\langle \cos(n_1\phi_1 + \cdots + n_k\phi_k) \rangle = v_{n_1} \cdots v_{n_k} \cos(n_1\Psi_1 + \cdots + n_k\Psi_k)$$

R. S. Bhalerao, M. Luzum and J.-Y. Ollitrault, 'Determining initial-state fluctuations from flow measurements in heavy-ion collisions,' PRC **84** 034910 (2011)

• Example:

 $\langle \cos(3\varphi_1 + 3\varphi_2 - 2\varphi_3 - 2\varphi_4 - 2\varphi_5) \rangle = v_3^2 v_2^3 \cos[6(\Psi_3 - \Psi_2)]$



What is the relation between symmetry planes Ψ_n ?





=> Observe non-zero genuine 5-particle correlation
 => Correlation strength is related to three-plane correlations

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Three plane correlation



 $\langle \cos(\varphi_a - 3\varphi_b + 2\varphi_c) \rangle = \langle \cos(\varphi_a - 3\varphi_b + 2\Psi_2) \rangle \times v_2$



- * Observe non-zero 3-plane correlation
- * The shape at low p_T is similar to that expected from the hydrodynamic model calculations, but differs at higher p_T

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Summary



- Elliptic flow at LHC energies is 30% larger than at RHIC, compatible with expectations from hydrodynamic model calculations
- Flow fluctuations
 - Observed significant triangular flow
 - Do not change significanly neither with p_T nor with eta
 - Underlying p.d.f. is consistent with Bessel-Gaussian function
- Studied the correlation between symmetry planes along which flow develops
 - Symmetry plane of v₂ is not the same as the symmetry plane of v₃
- Observed non-zero mixed harmonic 3-particle correlation => indication of the three plane correlation
- Wealth of experimental results which demands a theoretical interpretation





Thanks!







Backup slides



Anisotropic flow



- Why we do not care about 'sinus terms'?
 - It is equally probable for a particle to be produced in directions phi and phi:

$$\sin(n\varphi) + \sin[n(-\varphi)] = \sin(n\varphi) - \sin(n\varphi) = 0$$

Can 'odd cosine terms' be non-zero for ideal geometry?
It is equally probable for a particle to be produced in directions phi and phi + Pi:

$$\cos(n\phi) + \cos[n(\phi + \pi)] = \cos(n\phi) + \cos(n\phi)\cos(n\pi) - \sin(n\phi)\sin(n\pi)$$

=
$$\cos(n\phi) + \cos(n\phi)(-1)^n - \sin(n\phi) \cdot 0$$

=
$$\cos(n\phi) \cdot (1 + (-1)^n) = 0 \text{ for odd } n$$







 Proof of the principle => Using Therminator events (realistic Monte Carlo generator of heavy-ion events, has both anisotropic flow and all resonances in the standard model)



In this regime multi-particle QCs are precision method

ALICE

Non-uniform acceptance (1/2)

- If a detector has non-uniform acceptance in azimuthal angle, than in each event we have trivial anisotropies in momentum distributions of detected particles => this has nothing to do with anisotropic flow!
 - Can we disentangle one anisotropy from another?





 Generalized Q-cumulants can correct for non-uniform acceptance very well



Grey band => v_2 {MC}; open markers => v_2 {4} from isotropic *Q*-cumulants; filled markers => v_2 {4} from generalized *Q*-cumulants



Is it really that trivial?



- Anisotropic flow measurement => 'recipe':
 - Step 1: Measure/estimate reaction (symmetry) plane in an event
 - Step 2: Take azimuthal angles of all reconstructed particles in an event
 - Step 3: Evaluate anisotropic flow harmonics via the average

$$v_n = \langle \cos(n(\phi - \Psi_{\rm RP})) \rangle$$

- In experimental practice the above prescription will not work
 - We cannot neither measure directly nor estimate reaction (symmetry) plane reliably event-by-event
- Can we estimate anisotropic flow harmonics v_n without requiring the reaction (symmetry) plane?



Autocorrelations



 We have to correlate only distinct particles, because autocorrelations are dominant and useless (really!) contribution in all averages. So:

$$\langle 2 \rangle \equiv \langle \cos n(\phi_1 - \phi_2) \rangle , \qquad \phi_1 \neq \phi_2$$

 $\langle 4 \rangle \equiv \langle \cos n(\phi_1 + \phi_2 - \phi_3 - \phi_4) \rangle , \qquad \phi_1 \neq \phi_2 \neq \phi_3 \neq \phi_4$

• How to enforce above constrains in practice?

- Brute force evaluation via nested loops? => not feasible
 - Formalism of generating functions? => only approximate

$$G_n(z) \equiv \prod_{j=1}^M \left(1 + \frac{z^* e^{in\phi_j} + z e^{-in\phi_j}}{M} \right)$$

$$G_n(z) = \sum_{k=0}^{M/2} \frac{|z|^{2k}}{M^{2k}} \binom{M}{k} \binom{M-k}{k} \left\langle e^{in(\phi_1 + \dots + \phi_k - \phi_{k+1} - \dots - \phi_{2k})} \right\rangle$$

N. Borghini, P. M. Dinh and J.-Y. Ollitrault, "Flow analysis from multiparticle azimuthal correlations," PRC 64 (2001) 054901





Elliptic flow in pp?



- Both 2- and 4-particle correlations decrease with multiplicity: Typical for non-collective behavior
- Pythia and Phojet are overestimating the strength of the correlations (and these two generators are dominated by jets and resonances)
- 4-p cumulant comes with a "wrong sign" => its dominant contribution is not coming from flow
- Current status We do not see elliptic flow in pp



Flow at first sight!





(reference multiplicity)





Early results





Phys. Rev. Lett. 105, 252302 (2010)

 p_T dependence of elliptic flow at LHC close to the one at RHIC!

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Comparison to models



arXiv:1105.3865 ~ ~ Centrality 30-40% Model: Schenke et al, hydro, Glauber init, conditions full: $|\Delta \eta| > 0.2$ open: Δη > 1.0 0.3 $v_2 (\eta/s = 0.0)$ $v_2 (\eta/s = 0.08)$ $v_3 (\eta/s = 0.0)$ 0.2 $-v_{2}(\eta/s = 0.08)$ 0.1 2 3 Ω p, (GeV/*c*)

Within this model overall magnitude of v_2 and v_3 seems to be fine, but the details of p_t dependence are not well described

- More quantitative statement: The magnitude of $v_2(p_t)$ is described better with eta/s = 0, while for $v_3(pt)$ eta/s = 0.08 provides a better description
- This model fails to describe well v₂ and v₃ simultaneously
- Produced matter in Pb-Pb collisions at LHC continues to behave as a nearly perfect liquid



ALICE fluctuations



http://arxiv.org/abs/arXiv:1106.6284



Similar values for relative flow fluctuations at LHC and at RHIC



QC{5}



 Isolating the corresponding cumulants and quantifying the theoretical contributions:

$$QC\{5\} = \langle \cos(2\phi_1 + 2\phi_2 - 2\phi_3 - \phi_4 - \phi_5) \rangle$$
$$-2 \cdot \langle \cos(2\phi_1 - \phi_2 - \phi_3) \rangle \langle \cos(2\phi_1 - 2\phi_2) \rangle$$
in theory
$$= -v_2^3 v_1^2 \cos[2(\Psi_2 - \Psi_1)]$$
and:

 $QC{5}$

in theory

$$\langle \cos(3\phi_1 + 2\phi_2 - 2\phi_3 - 2\phi_4 - \phi_5) \rangle - 2 \cdot \langle \cos(3\phi_1 - 2\phi_2 - \phi_3) \rangle \langle \cos(2\phi_1 - 2\phi_2) \rangle - v_3 v_2^3 v_1 \cos[3\Psi_3 - 2\Psi_2 - \Psi_1)]$$



v_n {4} vs transverse momentum

Discovery



• Strong centrality dependence of $v_2{4} =>$ contribution from v_2 wrt. Ψ_{RP}

• Weak centrality dependence of v_3 {4} typical for pure flow fluctuations 44