

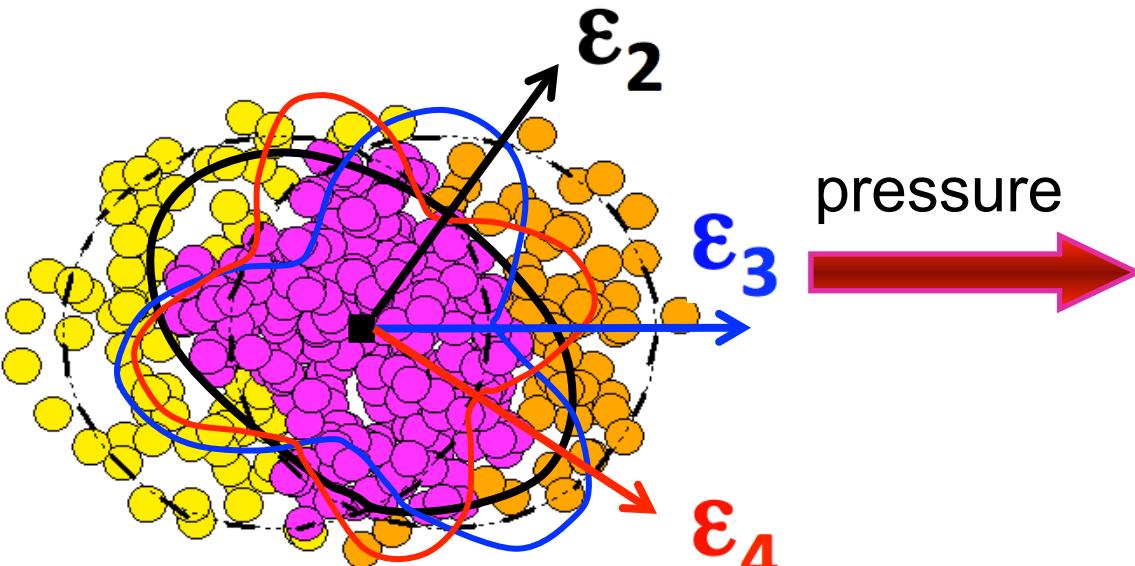
Event-by-event flow and initial geometry from LHC

Jiangyong Jia



Twelfth Workshop on Non-Perturbative Quantum Chromodynamics

Initial geometry & momentum anisotropy



$$\varepsilon_n = \sqrt{\frac{\langle r^n \cos(n\phi) \rangle + \langle r^n \sin(n\phi) \rangle}{r^n}} \quad \tan(n\Phi_n^*) = \frac{\langle r^n \sin(n\phi) \rangle}{\langle r^n \cos(n\phi) \rangle}$$

Single particle distribution

$$\frac{dN}{d\phi} \propto 1 + 2 \sum_n v_n \cos n(\phi - \Phi_n)$$

Pair distribution

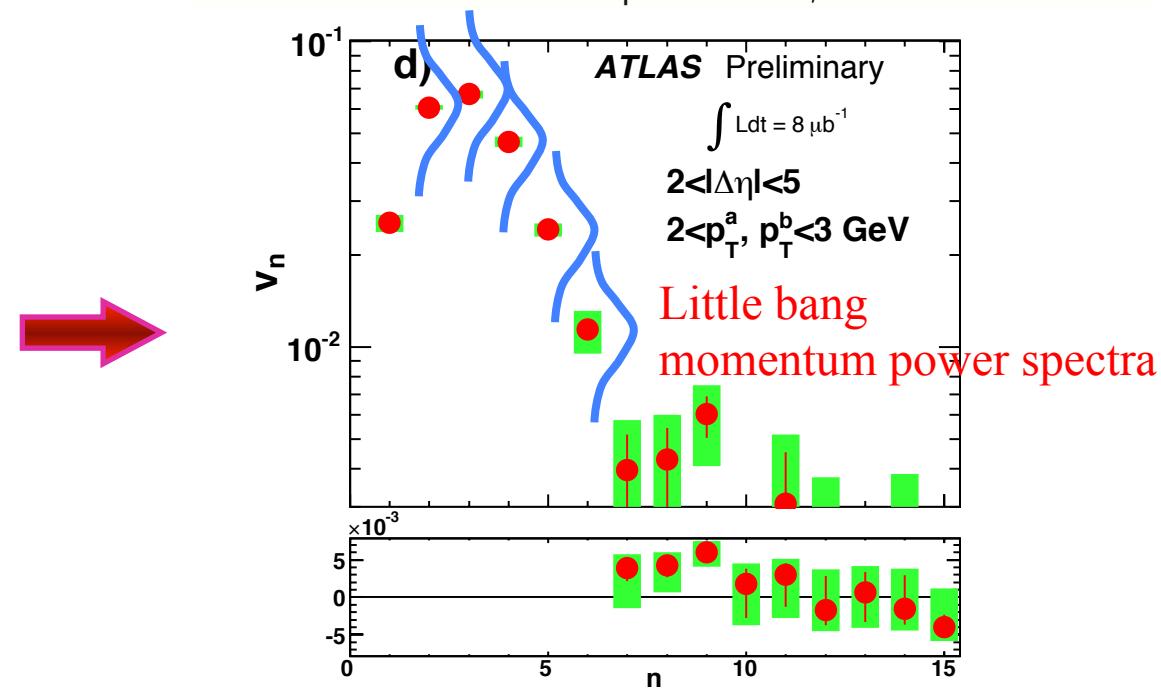
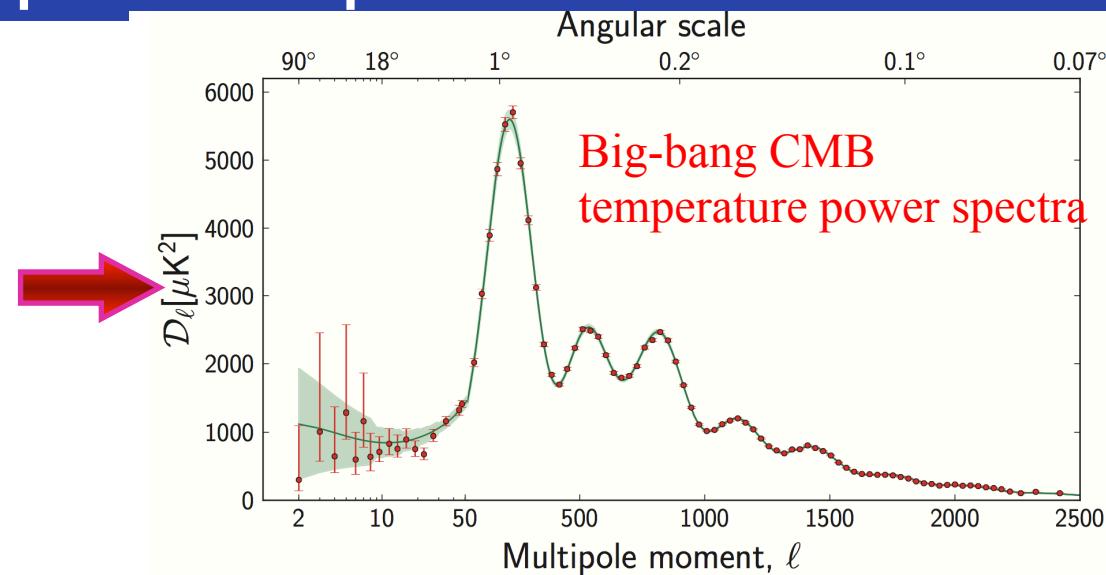
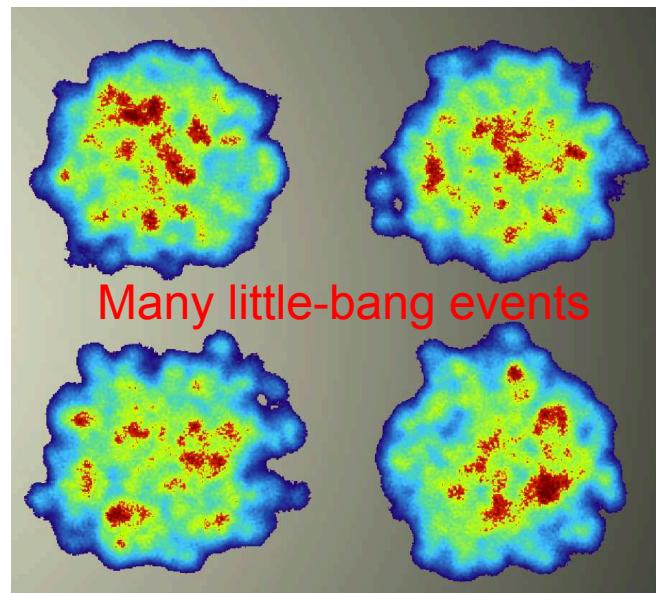
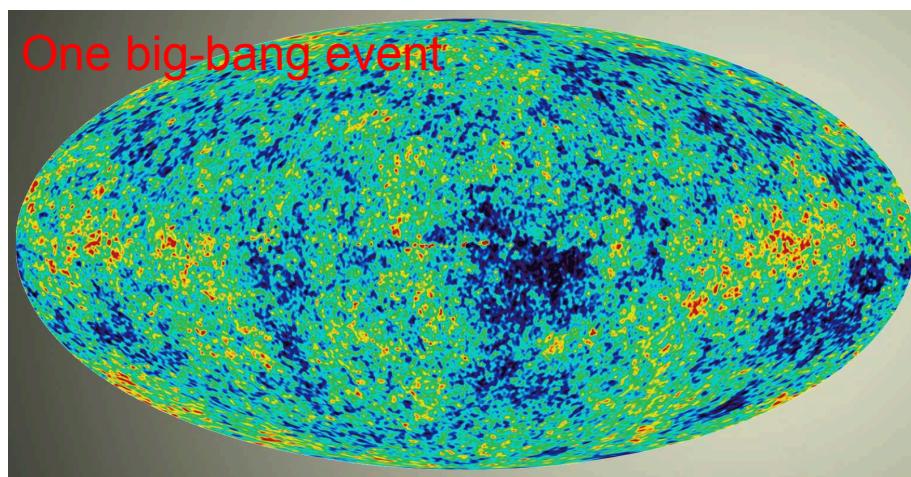
$$\frac{dN}{d\Delta\phi} = \left[\frac{dN}{d\phi_a} * \frac{dN}{d\phi_b} \right] \propto 1 + \sum_n 2 v_n^a v_n^b \cos(n\Delta\phi)$$

Momentum anisotropy probes:

initial geometry and transport properties of the QGP

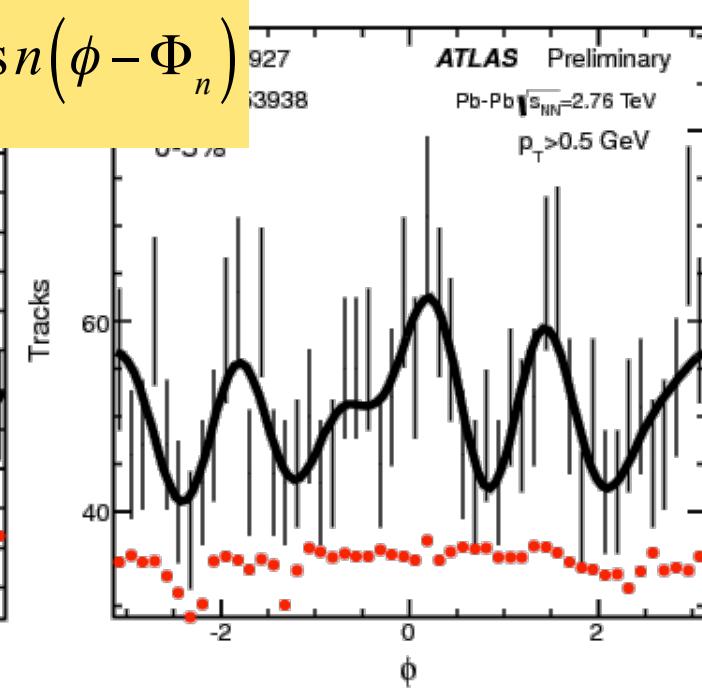
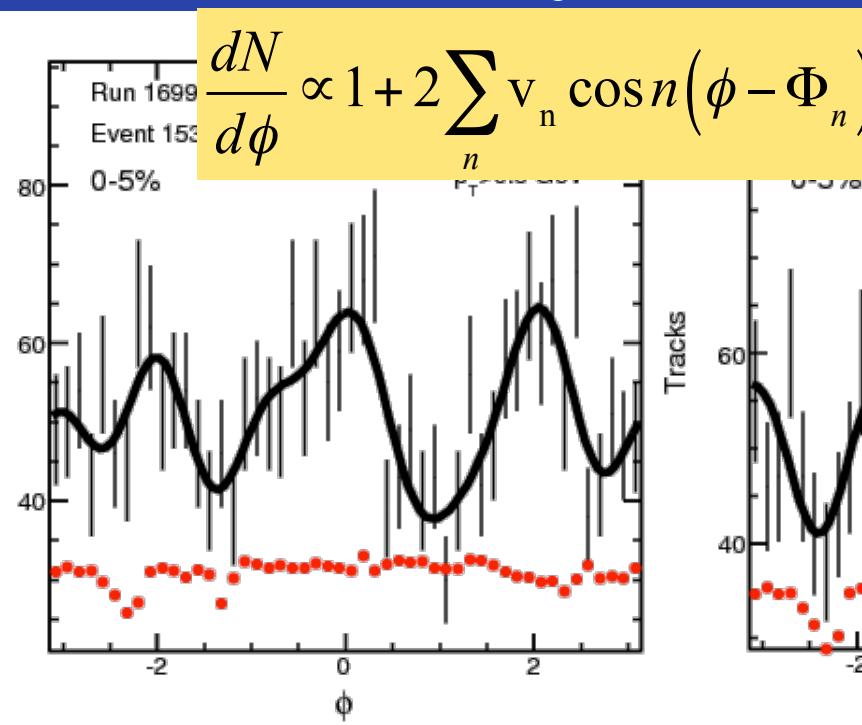
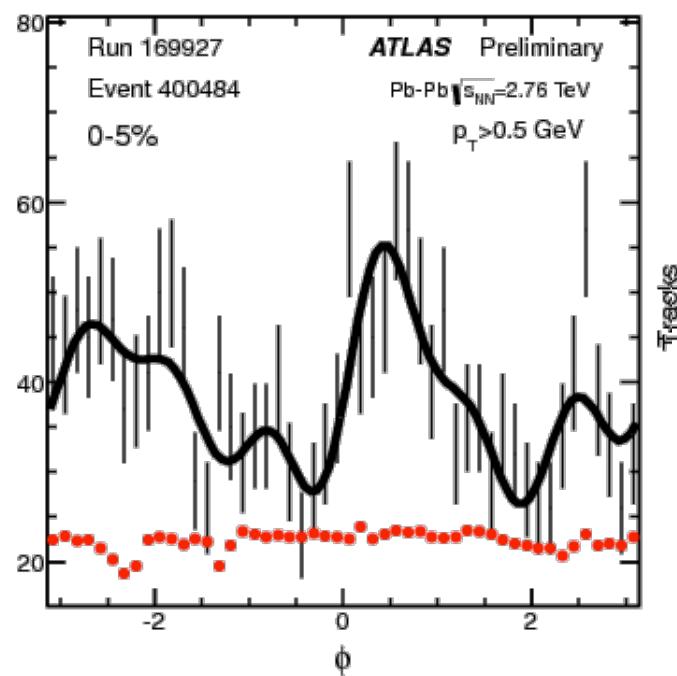
by MADAI.us

Anisotropy power spectra

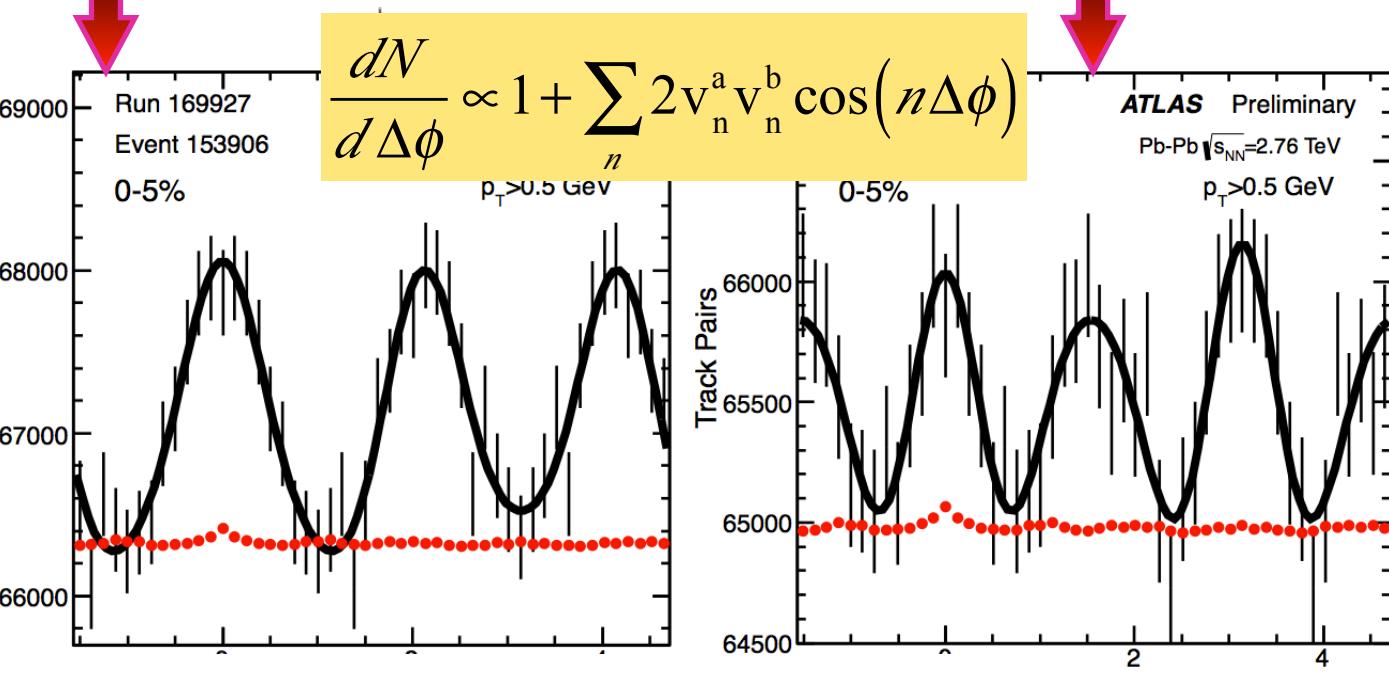
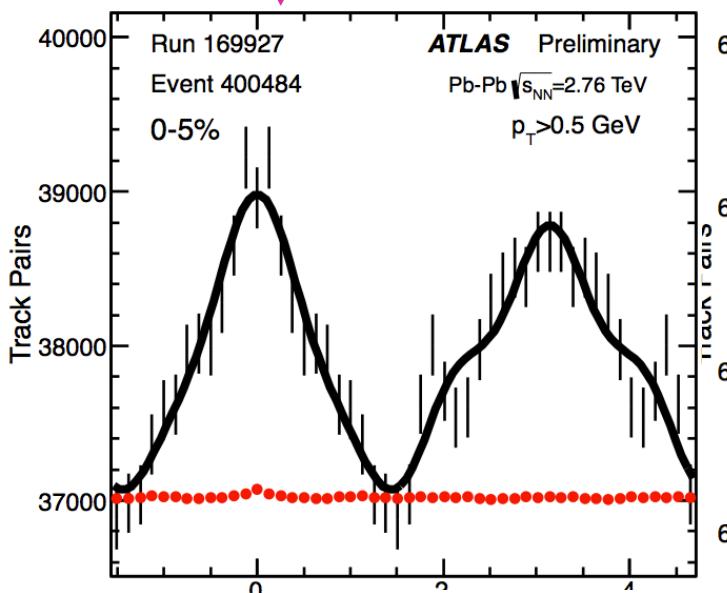
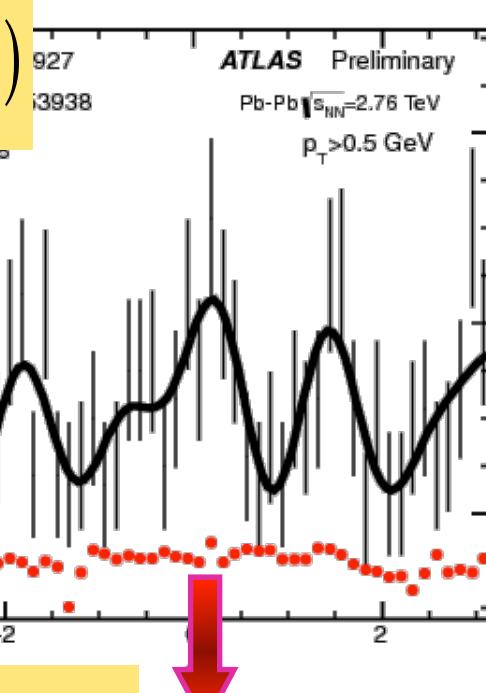
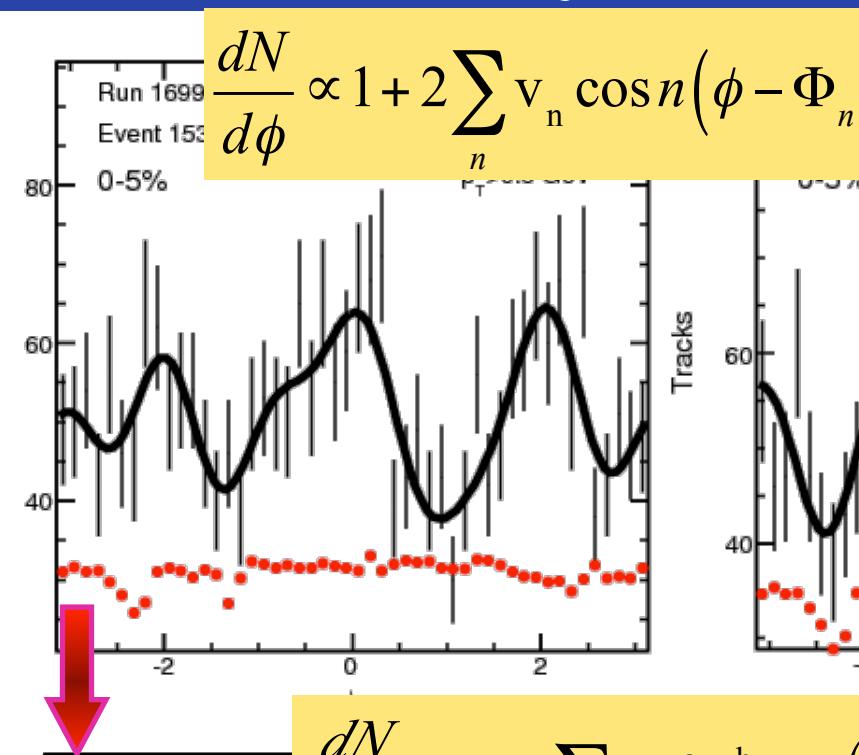
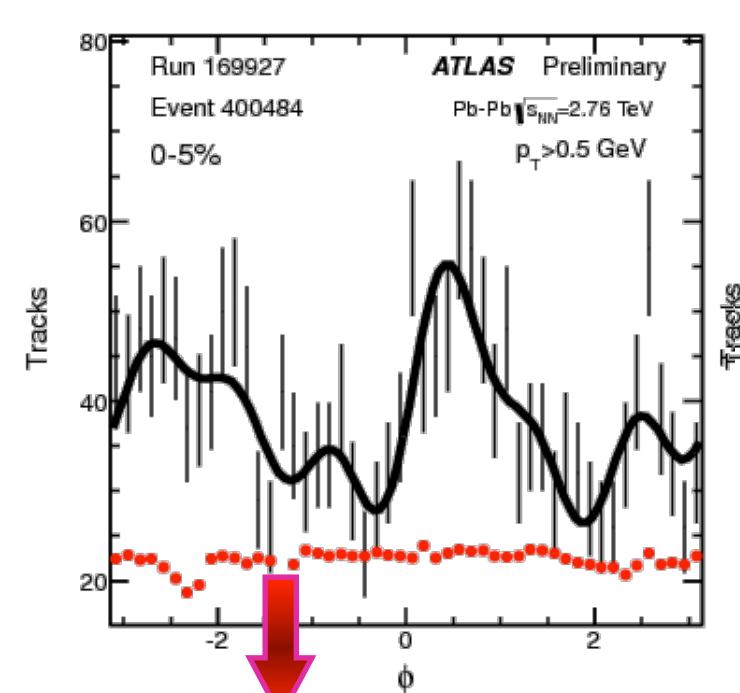


Many little-bang events → probability distributions: $p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots)$

Fluctuation event by event



Fluctuation event by event



Rich event-by-event patterns for v_n and Φ_n !

Outline

$$p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots)$$

- Event-plane correlations $\rightarrow p(\Phi_n, \Phi_m, \dots)$

ATLAS-CONF-2012-49

- Event-by-event v_n distributions $\rightarrow p(v_n)$

1305.2942

$$p(\Phi_n, \Phi_m, \dots)$$

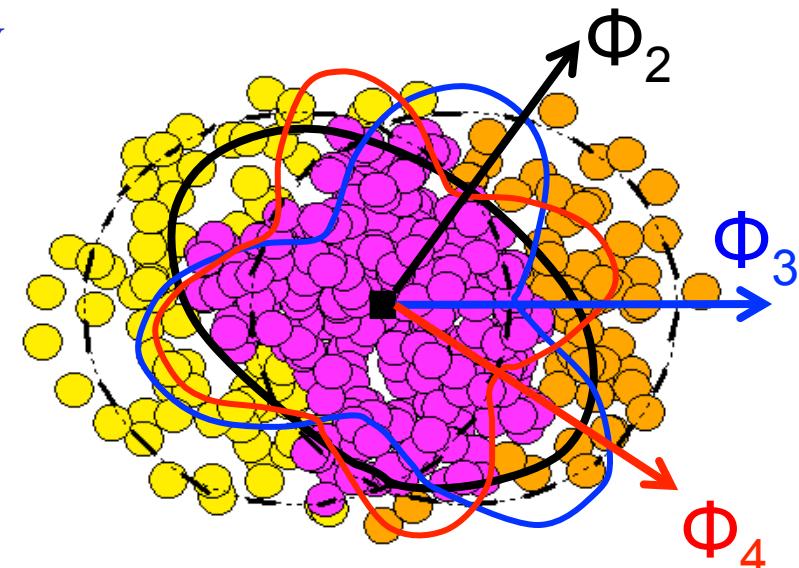
- Correlation can exist in the initial geometry and also generated during hydro evolution
- The correlation quantified via correlators

$$\frac{dN_{events}}{d(jk(\Phi_n - \Phi_m))} = 1 + 2 \sum_{j=1}^{\infty} V_{n,m}^j \cos(jk(\Phi_n - \Phi_m))$$

$$V_{n,m}^j = \langle \cos(jk(\Phi_n - \Phi_m)) \rangle$$

arXiv:1205.3585

arXiv:1203.5095



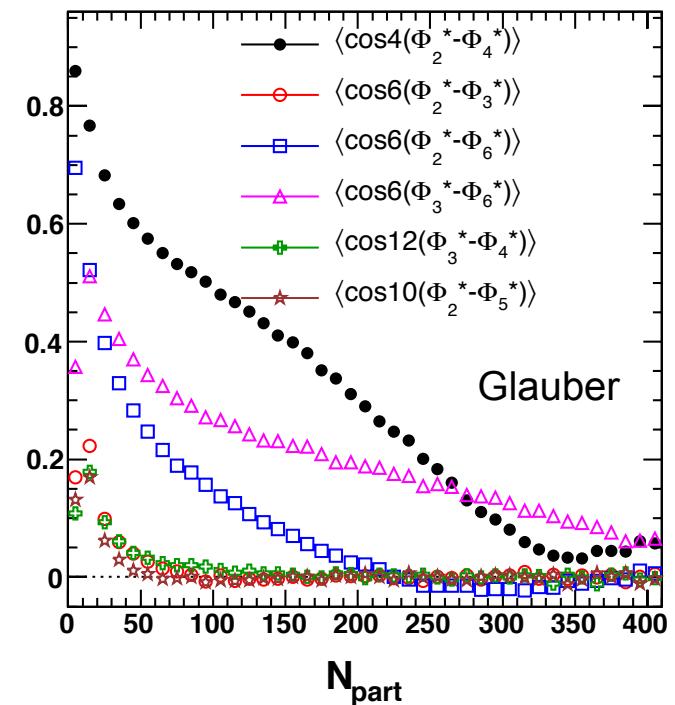
- Corrected by resolution

$$\langle \cos(jk(\Phi_n - \Phi_m)) \rangle = \frac{\langle \cos(jk(\Psi_n - \Psi_m)) \rangle}{\text{Res}(jk\Psi_n)\text{Res}(jk\Psi_m)}$$

$\Phi_n = True$, $\Psi_n = Measured$

- Generalize to multi-plane correlations

$$\langle \cos(c_1\Phi_1 + 2c_2\Phi_2 + \dots + c_l\Phi_l) \rangle \quad c_1 + 2c_2 + \dots + c_l = 0$$



A list of measured correlators

- List of two-plane correlators

$$\begin{aligned} & \langle \cos 4(\Phi_2 - \Phi_4) \rangle \\ & \langle \cos 8(\Phi_2 - \Phi_4) \rangle \\ & \langle \cos 12(\Phi_2 - \Phi_4) \rangle \\ & \langle \cos 6(\Phi_2 - \Phi_3) \rangle \\ & \langle \cos 6(\Phi_2 - \Phi_6) \rangle \\ & \langle \cos 6(\Phi_3 - \Phi_6) \rangle \\ & \langle \cos 12(\Phi_3 - \Phi_4) \rangle \\ & \langle \cos 10(\Phi_2 - \Phi_5) \rangle \end{aligned}$$

- List of three-plane correlators

“2-3-5”

$$\begin{aligned} & \langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle \\ & \langle \cos(-8\Phi_2 + 3\Phi_3 + 5\Phi_5) \rangle \end{aligned}$$

“2-4-6”

$$\begin{aligned} & \langle \cos(2\Phi_2 + 4\Phi_4 - 6\Phi_6) \rangle \\ & \langle \cos(-10\Phi_2 + 4\Phi_4 + 6\Phi_6) \rangle \end{aligned}$$

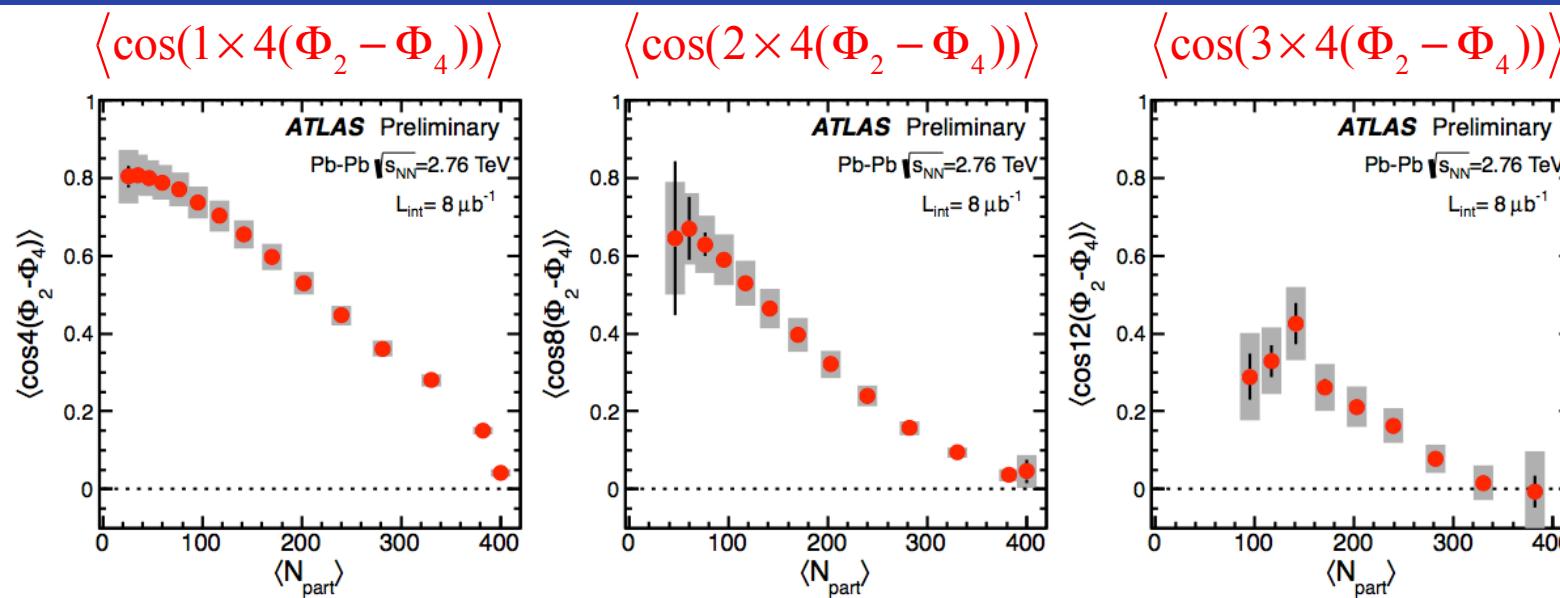
“2-3-4”

$$\begin{aligned} & \langle \cos(2\Phi_2 - 6\Phi_3 + 4\Phi_4) \rangle \\ & \langle \cos(-10\Phi_2 + 6\Phi_3 + 4\Phi_4) \rangle \end{aligned}$$

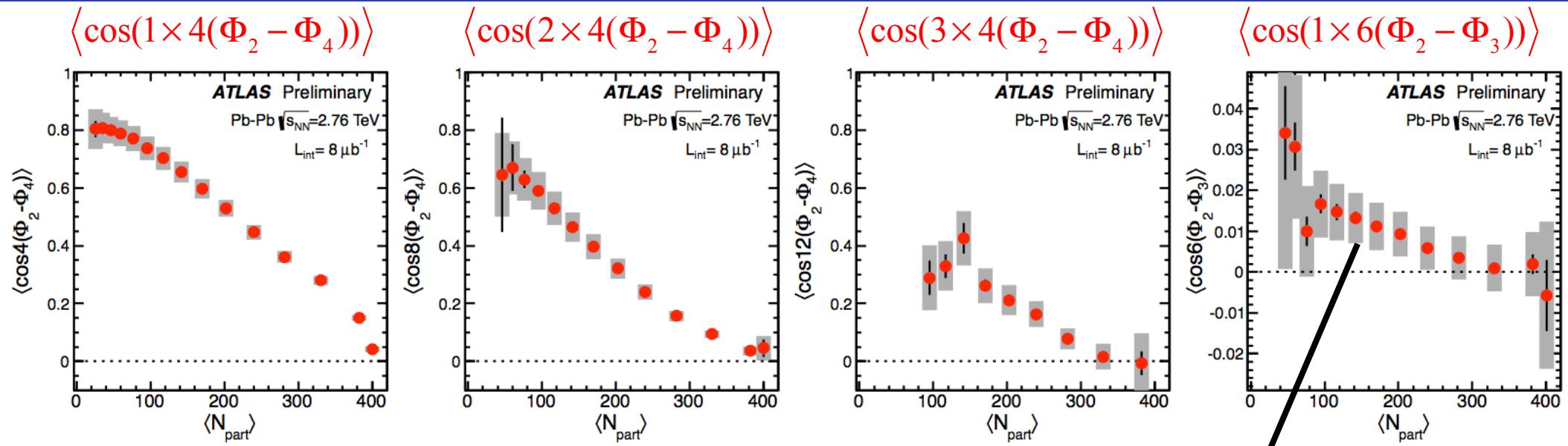
$$\begin{aligned} 2\Phi_2 + 4\Phi_4 - 6\Phi_6 &= 4(\Phi_4 - \Phi_2) - 6(\Phi_6 - \Phi_2) \\ -10\Phi_2 + 4\Phi_4 + 6\Phi_6 &= 4(\Phi_4 - \Phi_2) + 6(\Phi_6 - \Phi_2) \end{aligned}$$

Reflects correlation of two Φ_n relative to the third

Two-plane correlations



Two-plane correlations

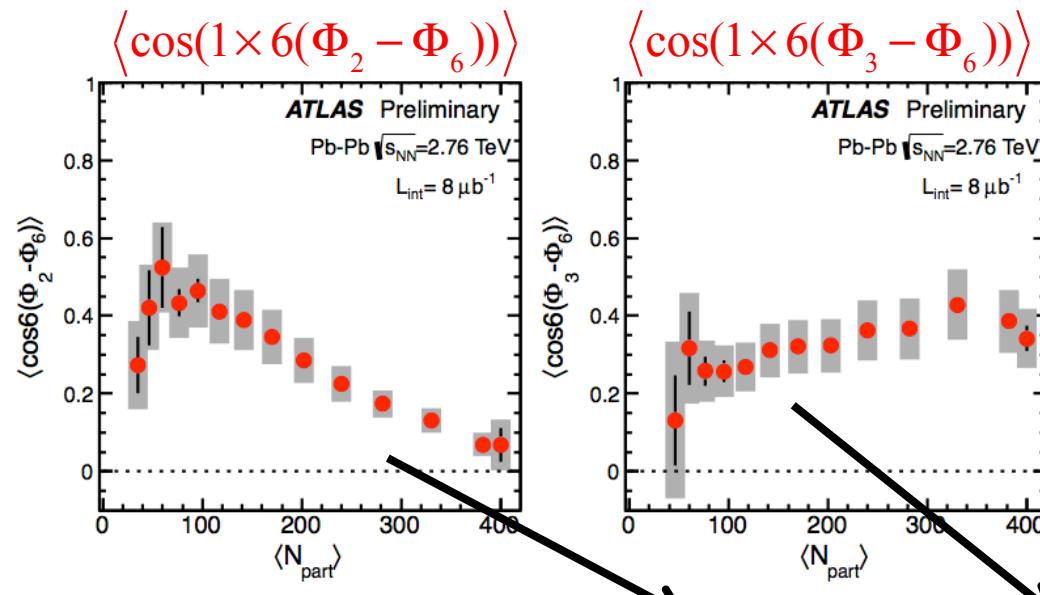
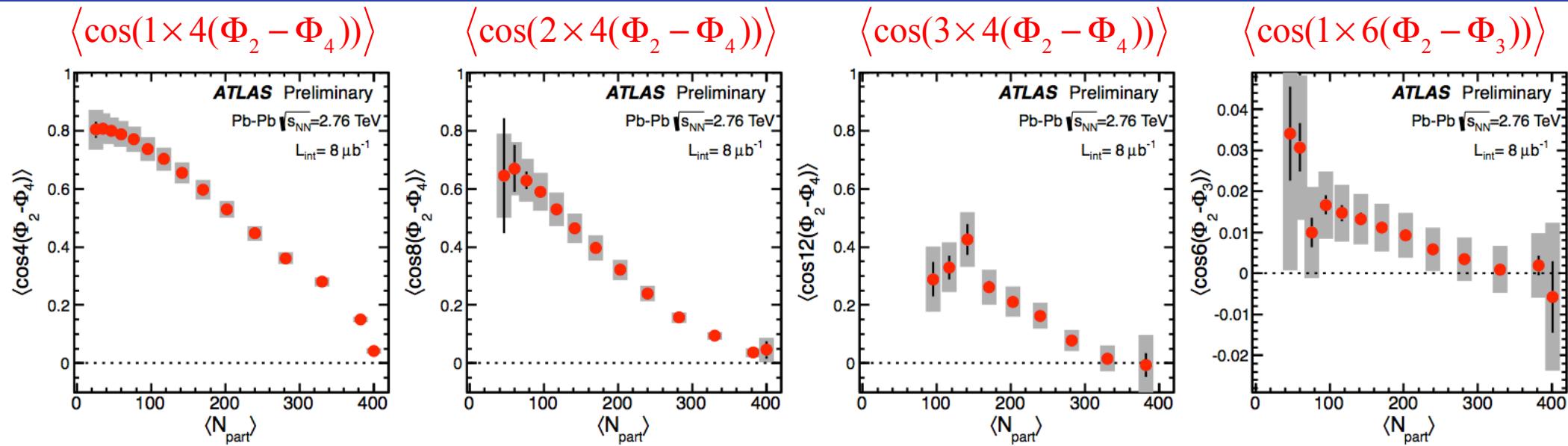


$$\nu_2 e^{-i2\Psi_2} = \text{geometry} + \nu_1 \nu_1 e^{-i2\Psi_1} + \dots$$

$$\nu_3 e^{-i3\Psi_3} = \text{geometry} + \nu_1 \nu_2 e^{-i(\Psi_1 + 2\Psi_2)} + \dots$$

Teaney & Yan

Two-plane correlations



$$\nu_6 e^{-i6\Psi_6} = \text{geometry} + \nu_2 \nu_2 \nu_2 e^{-i6\Psi_2} + \nu_3 \nu_3 e^{-i6\Psi_3} + \dots$$

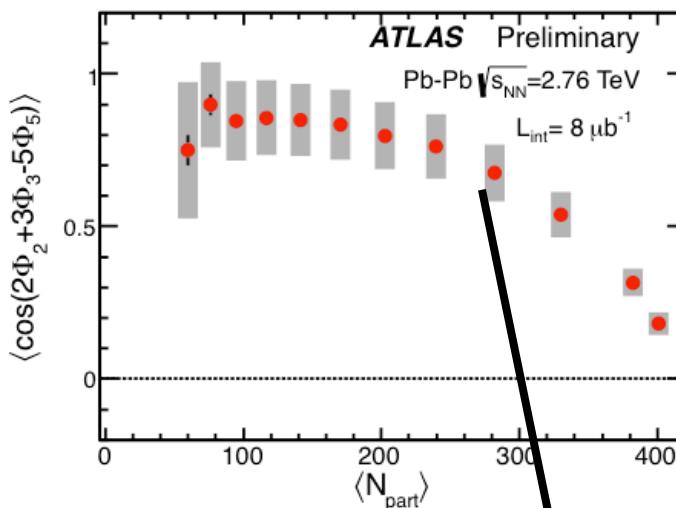
Teaney & Yan

Rich patterns for the centrality dependence

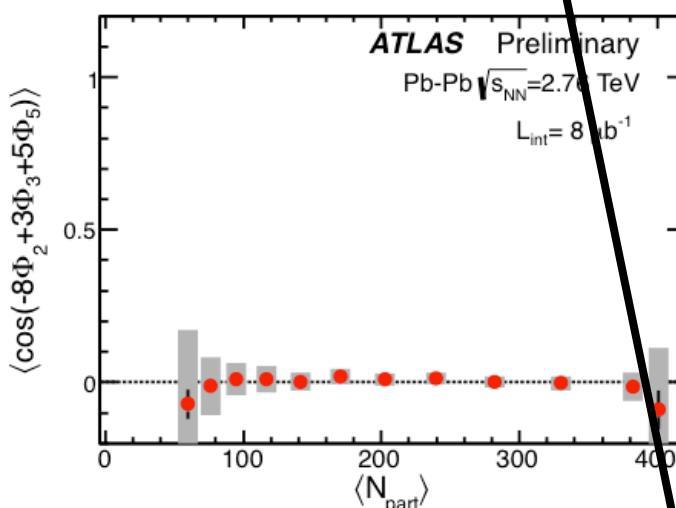
Three-plane correlations

"2-3-5" correlation

$$\langle \cos(2\Phi_2 + 3\Phi_3 - 5\Phi_5) \rangle$$



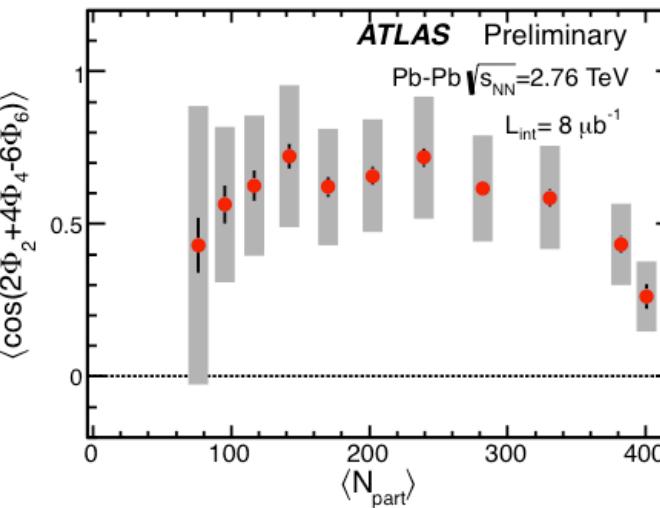
$$\langle \cos(-8\Phi_2 + 3\Phi_3 + 5\Phi_5) \rangle$$



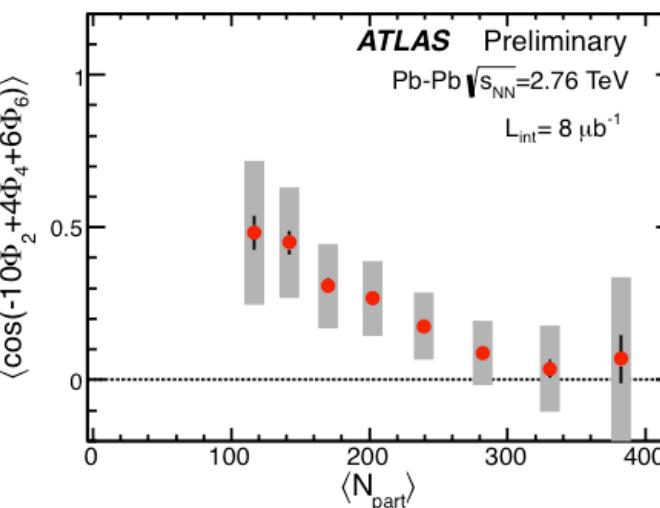
$$\nu_5 e^{-i5\Psi_5} = \text{geometry} + \nu_2 \nu_3 e^{-i(2\Psi_2 + 3\Psi_3)} + \dots$$

"2-4-6" correlation

$$\langle \cos(2\Phi_2 + 4\Phi_4 - 6\Phi_6) \rangle$$

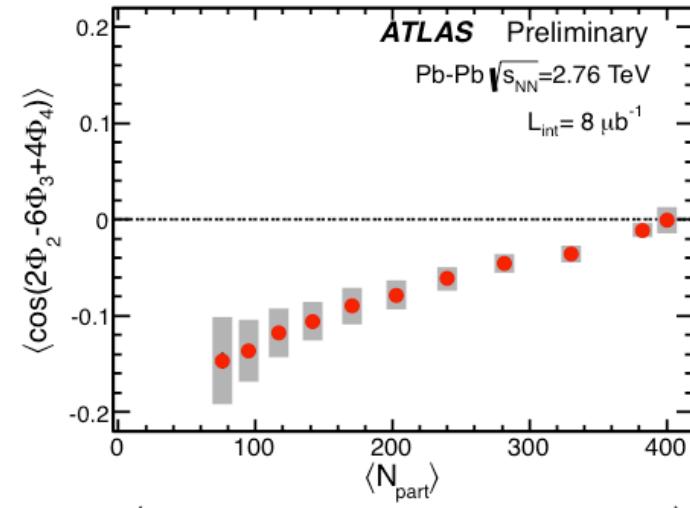


$$\langle \cos(-10\Phi_2 + 4\Phi_4 + 6\Phi_6) \rangle$$

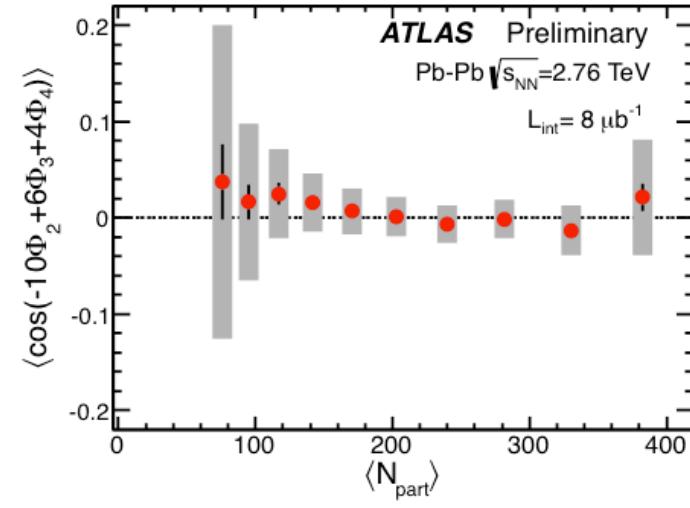


"2-3-4" correlation

$$\langle \cos(2\Phi_2 - 6\Phi_3 + 4\Phi_4) \rangle$$



$$\langle \cos(-10\Phi_2 + 6\Phi_3 + 4\Phi_4) \rangle$$



Rich patterns for the centrality dependence

Compare with EbE hydro calculation: 2-plane

Initial geometry
+ hydrodynamic

ATLAS data

MC-Glb., $\eta/s = 0.08$

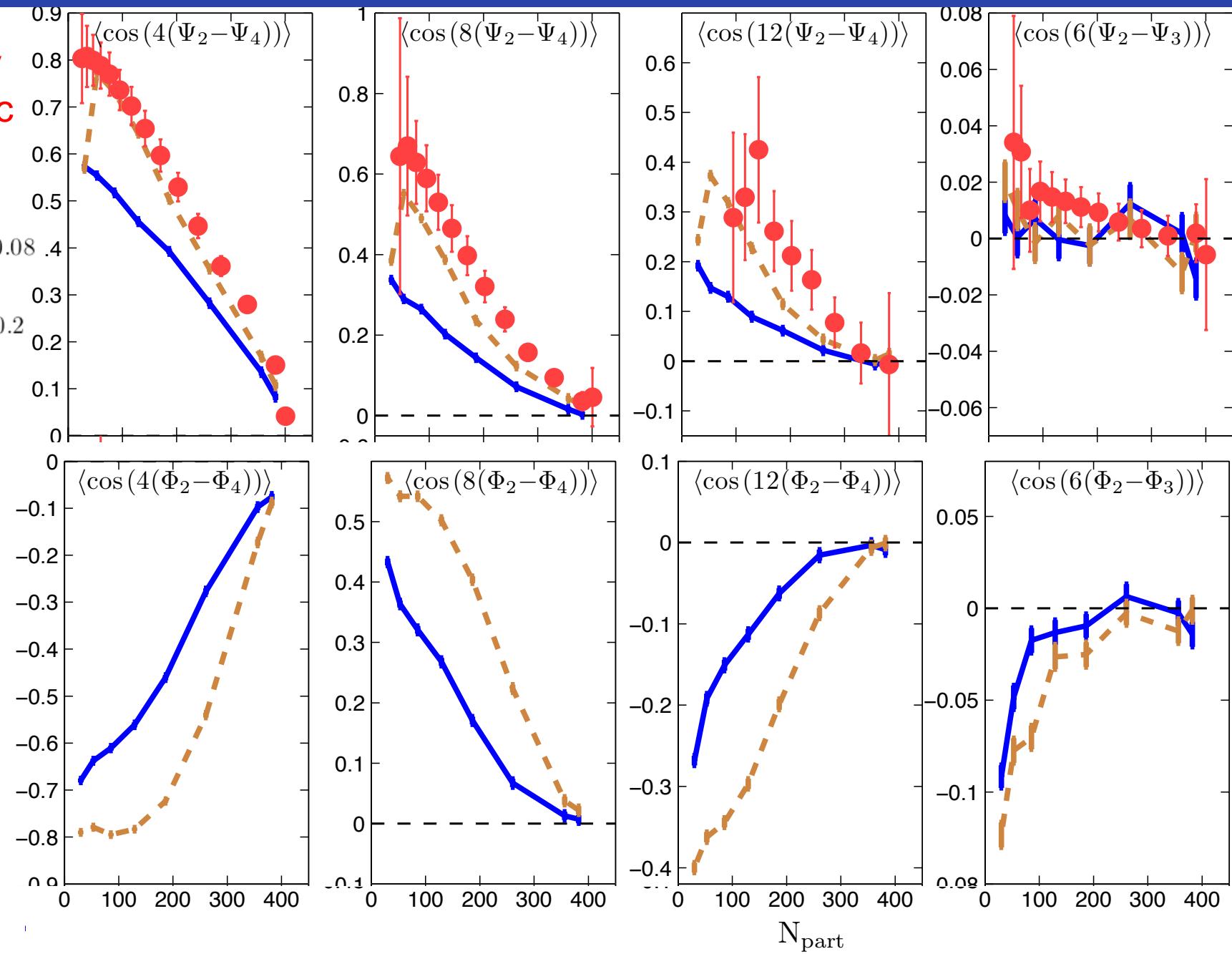
MC-KLN, $\eta/s = 0.2$

Zhe & Heinz

geometry only

MC-Glb.

MC-KLN,



EbyE hydro qualitatively reproduce features in the data

Compare with EbE hydro calculation: 3-plane

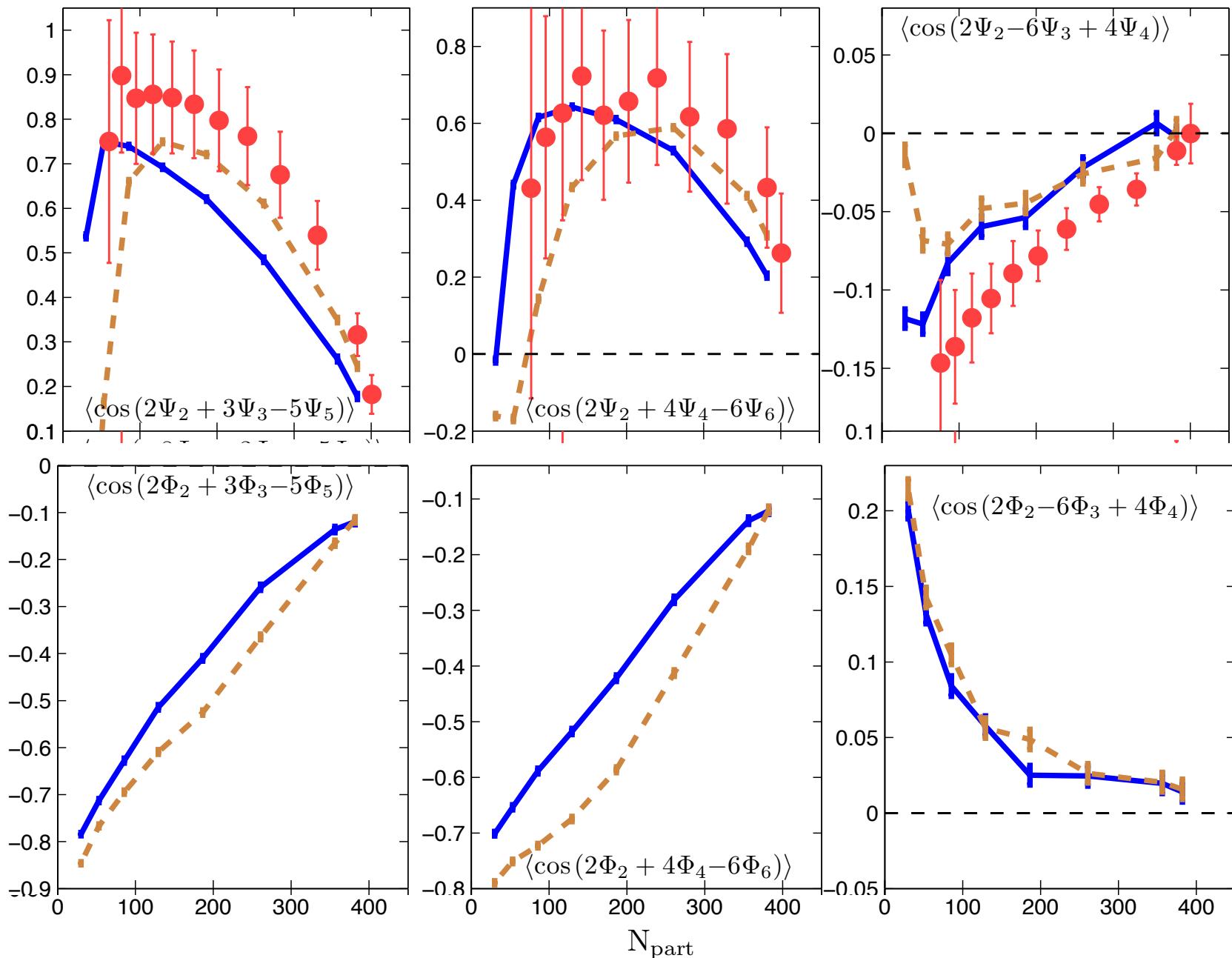
Initial geometry
+ hydrodynamic

- ATLAS data
- MC-Glb., $\eta/s = 0.08$
- - - MC-KLN, $\eta/s = 0.2$

Zhe & Heinz

geometry only

- MC-Glb.
- - - MC-KLN,



Over-constraining the transport properties

Event-by-event v_n distributions

Gaussian model of v_n fluctuations

- Flow vector $\frac{dN}{d\phi} \propto 1 + 2 \sum_n v_n \cos n(\phi - \Phi_n)$ $\vec{v}_n = (v_n \cos n\Phi_n, v_n \sin n\Phi_n) = \vec{v}_n^{\text{RP}} + \vec{p}_n^{\text{fluc}}$

- Gaussian model

arXiv: 0708.0800

arXiv: 0809.2949

$$p(\vec{v}_n) \propto \exp\left(\frac{-(\vec{v}_n - \vec{v}_n^{\text{RP}})^2}{2\delta_n^2}\right)$$

$$p(v_n) \propto v_n \exp\left(\frac{-(v_n^2 + (v_n^{\text{RP}})^2)}{2\delta_n^2}\right) I_0\left(\frac{v_n v_n^{\text{RP}}}{\delta_n^2}\right) \quad \text{Bessel-Gaussian function}$$

$$v_n^{\text{RP}} = 0 \Rightarrow p(v_n) \propto v_n \exp\left(\frac{-(v_n^2)}{2\delta_n^2}\right) \quad \text{For pure fluctuations}$$

- Multi-particle cumulants in Gaussian fluctuation limit

$$v_n\{2\} = \sqrt{(v_n^{\text{RP}})^2 + 2\delta_n^2}$$

$$v_n\{4\} = v_n\{6\} = v_n\{8\} = v_n^{\text{RP}}$$

In general $v_2\{2\}$ and $v_2\{4\}$ can be different even in the absence of nonflow

- Various estimators of the fluctuations:

$$\sqrt{\frac{v_n\{2\}^2 - v_n\{4\}^2}{v_n\{2\}^2 + v_n\{4\}^2}}$$

$$v_n^{\text{RP}} = 0 \Rightarrow$$

1

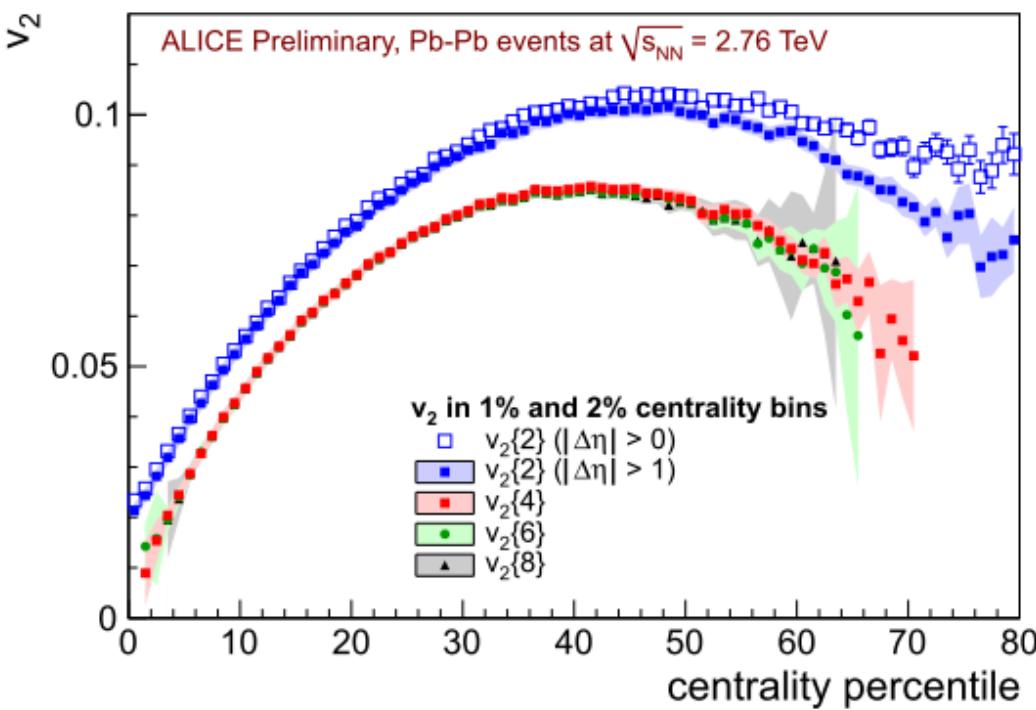
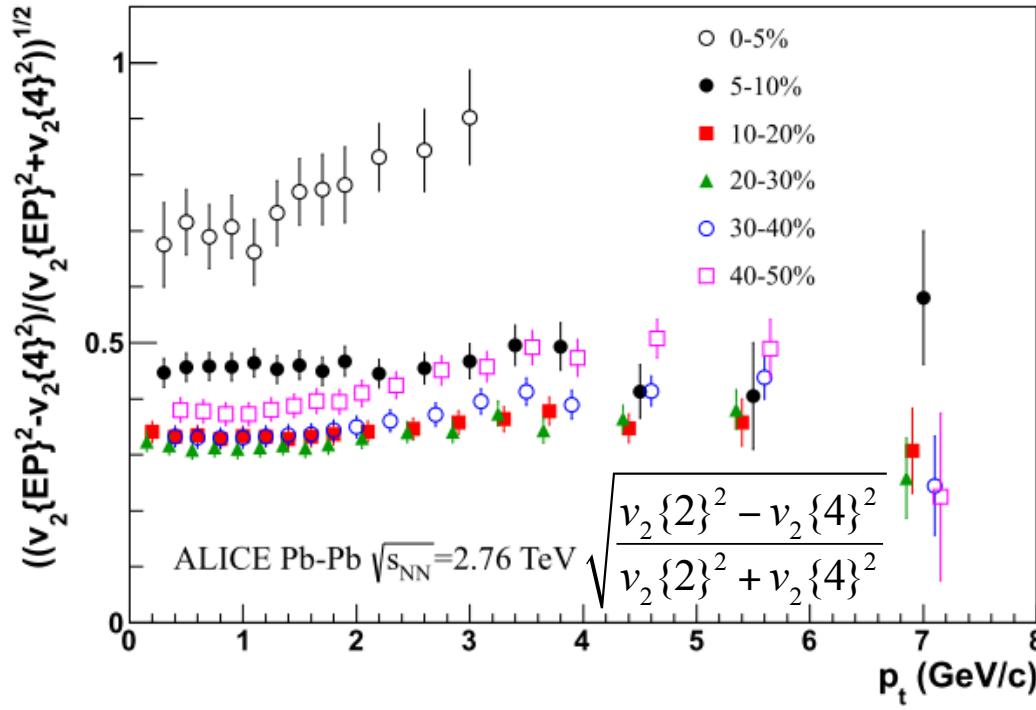
$$\sqrt{\frac{v_n\{2\}^2 - v_n\{4\}^2}{2v_n\{4\}^2}} = \frac{\delta_n}{v_n^{\text{RP}}}$$

II
∞

$$\frac{\sigma_n}{\langle v_n \rangle}$$

$$\sqrt{4/\pi - 1} = 0.52$$

Measuring v_2 fluctuations with Cumulants

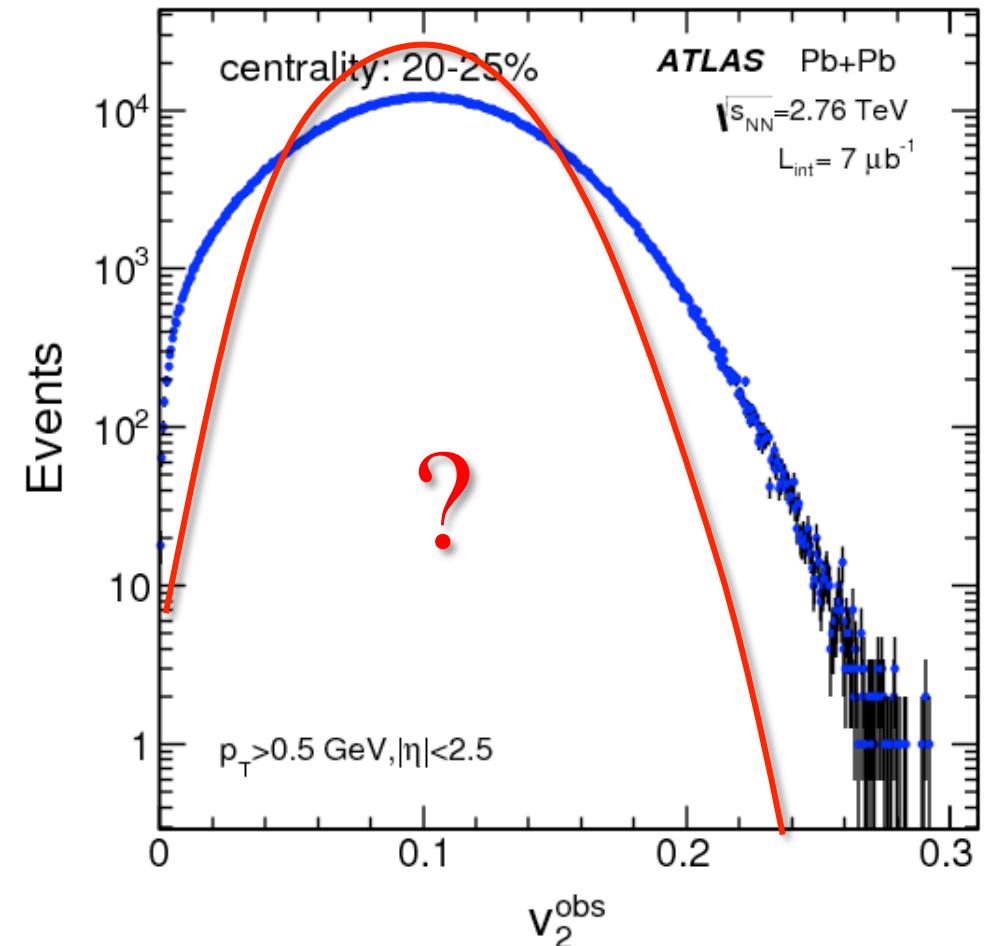
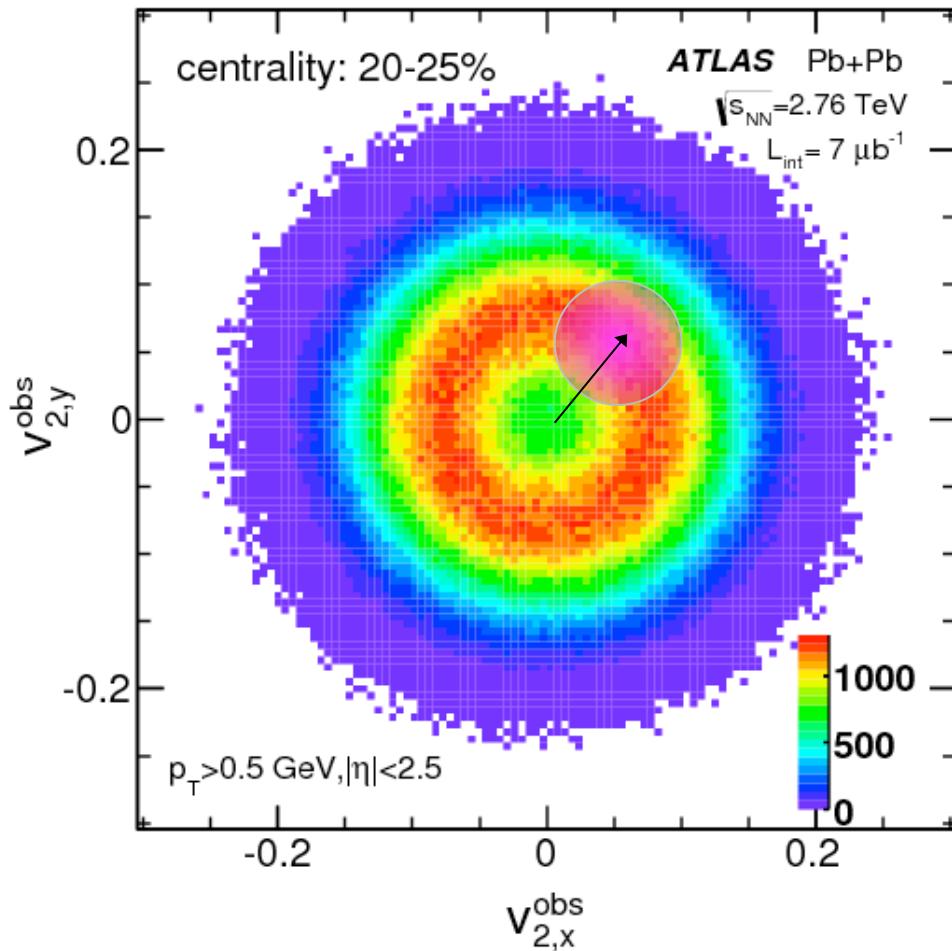


- Relative fluctuations increase with p_T (by 25%)
 - $v_2^{RP} \neq 0$ even in 0-5% central collisions
- Higher order cumulants such as $v_2\{6\}, v_2\{8\}$ all measure v_2^{RP}
 - Significant uncertainty in 0-5% centrality.
 - Fluctuation is described by Bessel-Gaussian function??

Flow vector and smearing

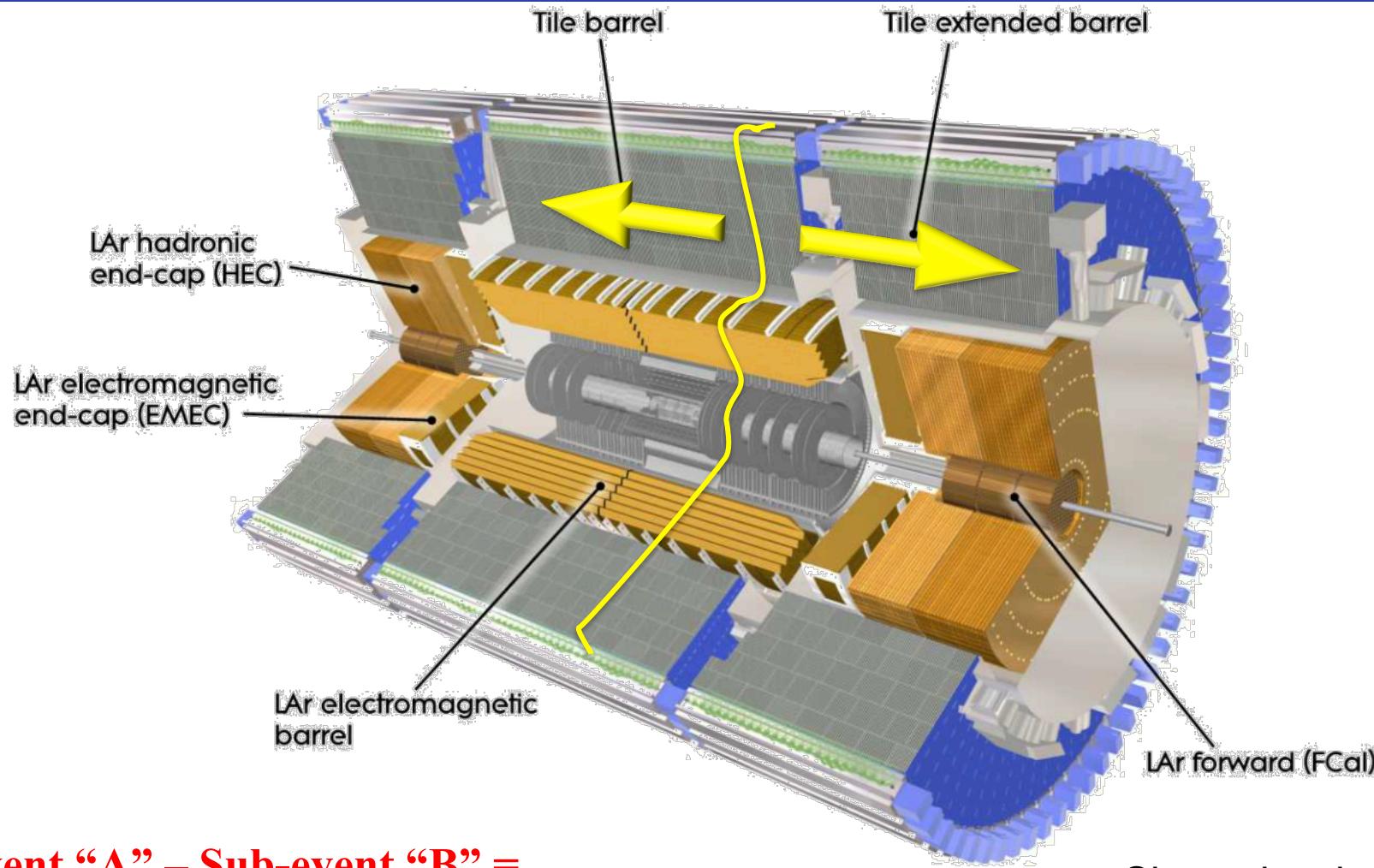
$$\frac{dN}{d\phi} \propto 1 + 2 \sum_n v_n^{obs} \cos n(\phi - \Phi_n^{obs})$$

$$\begin{aligned} \vec{v}_n^{\rightarrow obs} &= (v_{n,x}^{obs}, v_{n,y}^{obs}) = \vec{v}_n^{\rightarrow} + \vec{p}_n^{\rightarrow smear} \\ &\Downarrow \\ \vec{v}_n^{\rightarrow RP} &+ \vec{p}_n^{\rightarrow fluc} \end{aligned}$$



The key of unfolding is response function: $p(v_n^{obs} | v_n)$

Split the event into two: 2SE method



Sub-event “A” – Sub-event “B” =

Shown by simulation studies

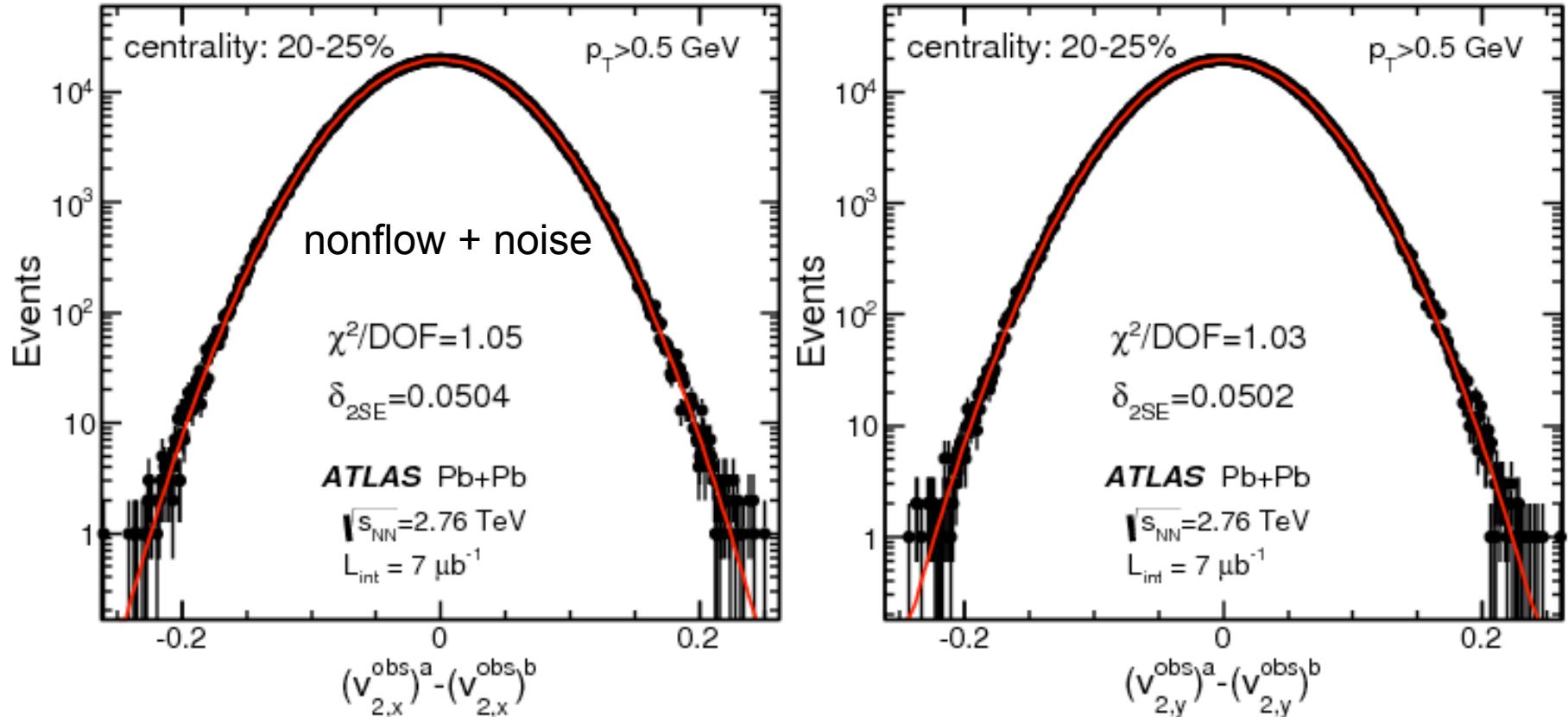
$$(\vec{v}_n^{\text{obs}})^a - (\vec{v}_n^{\text{obs}})^b = \text{nonflow} + \text{noise}$$

arxiv:1304.1471

$$(\vec{v}_n^{\text{obs}})^a + (\vec{v}_n^{\text{obs}})^b = \text{nonflow} + \text{noise} + 2\vec{v}_n$$

Width of $(\vec{v}_n^{\text{obs}})^a - (\vec{v}_n^{\text{obs}})^b$ $\xrightarrow{1/\sqrt{2}}$ Width of $(\vec{v}_n^{\text{obs}})^a$ $\xrightarrow{1/\sqrt{2}}$ Width of \vec{v}_n^{obs}

Obtaining the response function



Response function is a 2D Gaussian around truth

Nonflow is Gaussian!

$$p(\vec{v}_n^{\text{obs}} | \vec{v}_n) \propto e^{-\frac{|\vec{v}_n^{\text{obs}} - \vec{v}_n|^2}{2\delta^2}} \quad \delta = \delta_{\text{2SE}}/2$$

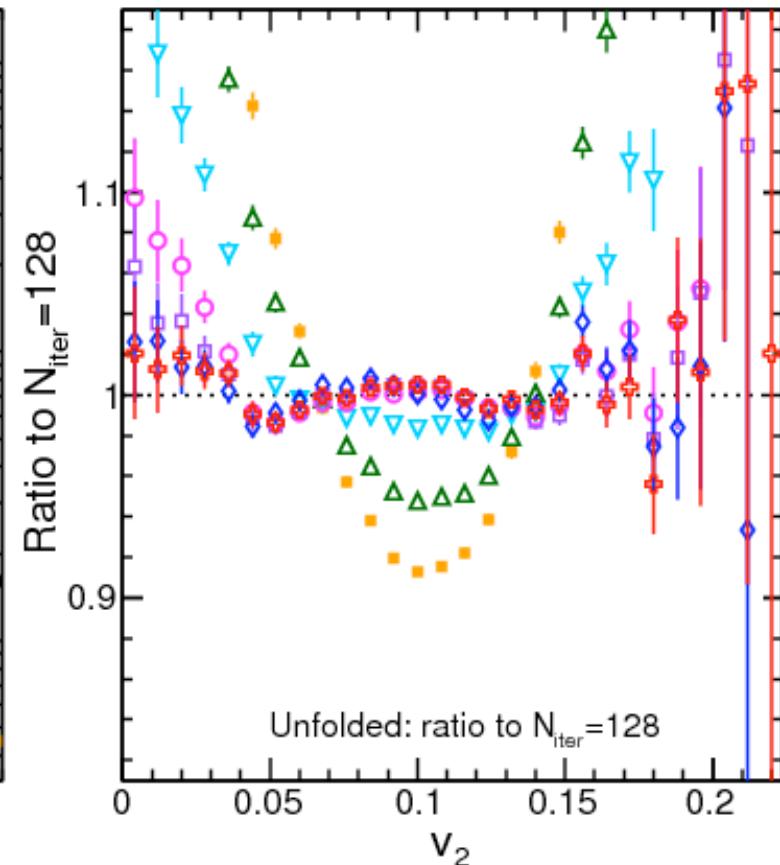
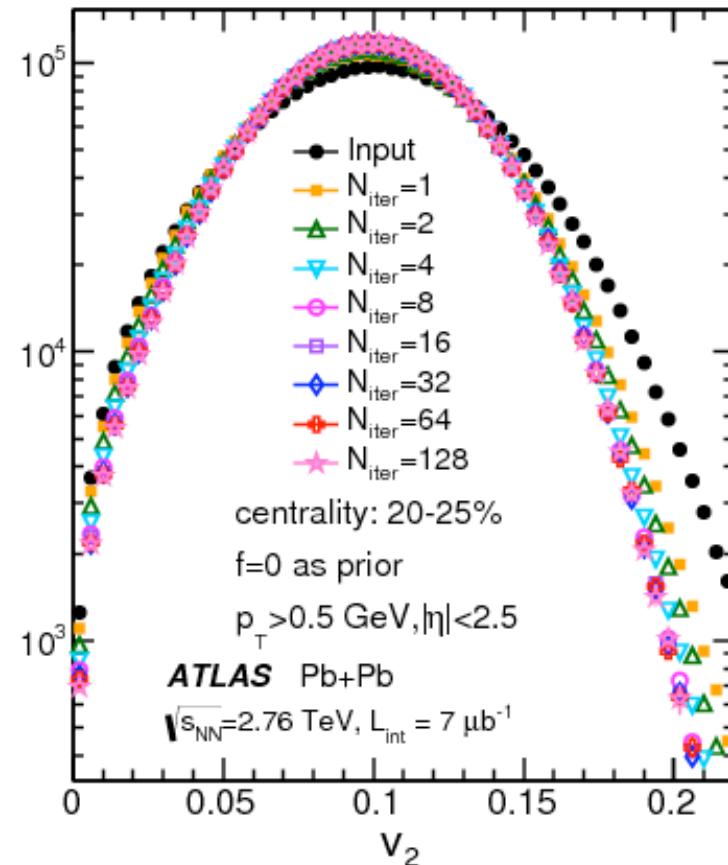
$$p(v_n^{\text{obs}} | v_n) \propto v_n^{\text{obs}} e^{-\frac{(v_n^{\text{obs}})^2 + v_n^2}{2\delta^2}} I_0\left(\frac{v_n^{\text{obs}} v_n}{\delta^2}\right)$$

Data driven method

Unfolding performance: v_2 , 20-25%

$$Av_n = v_n^{\text{obs}}$$

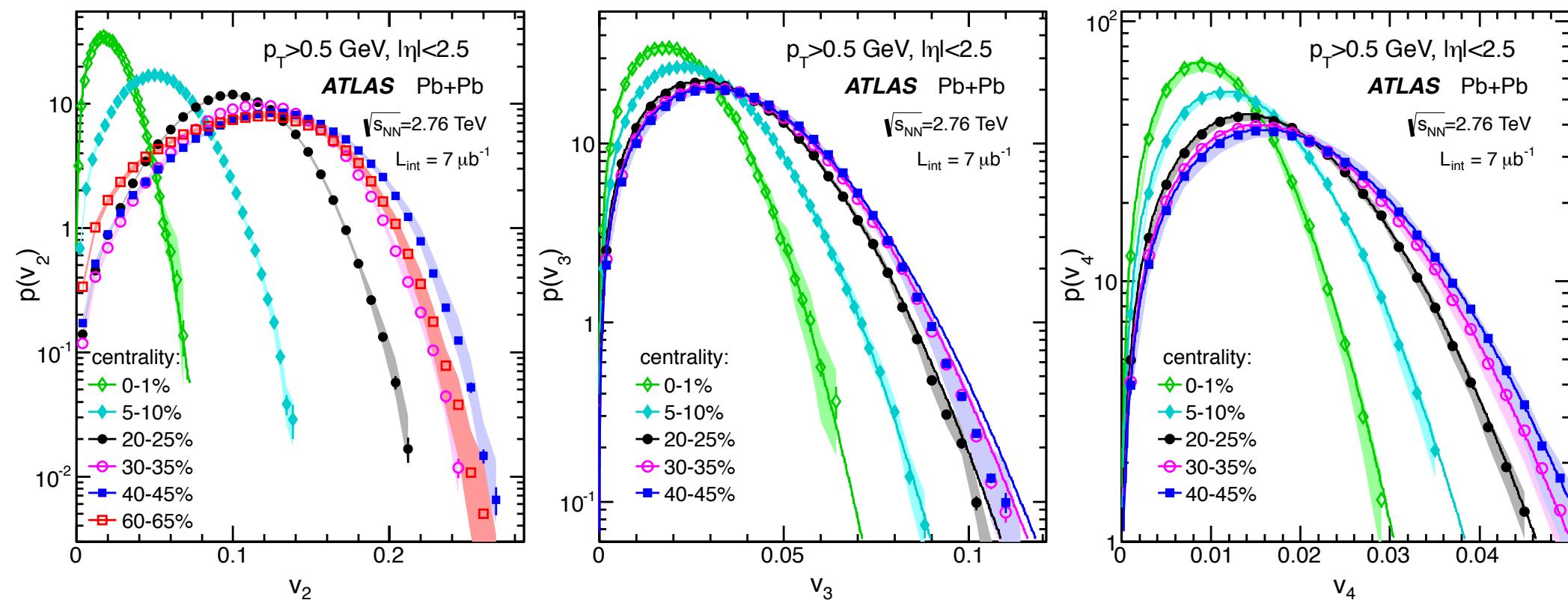
$$A_{ji} = p(v_{n,j}^{\text{obs}} | v_{n,i})$$



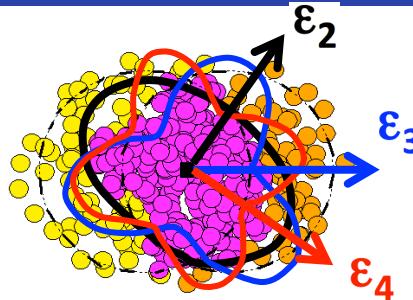
- Standard Bayesian unfolding technique
 - Converges within a few % for $N_{\text{iter}} = 8$, small improvements for larger N_{iter} .
 - Many cross checks show good consistency
 - Unfolding with different initial distributions
 - Unfolding using tracks in a smaller detector
 - Unfolding based on the EbyE two-particle correlation.
 - Closure test using HIJING+flow simulation
- Details in arxiv:1305.2942
- arxiv:1304.1471

$p(v_2)$, $p(v_3)$ and $p(v_4)$ distributions

- Measured in broad centrality over large v_n range
 - The fraction of events in the tails is less than 0.2% for v_2 and v_3 , and $\sim 1\text{-}2\%$ for v_4 .



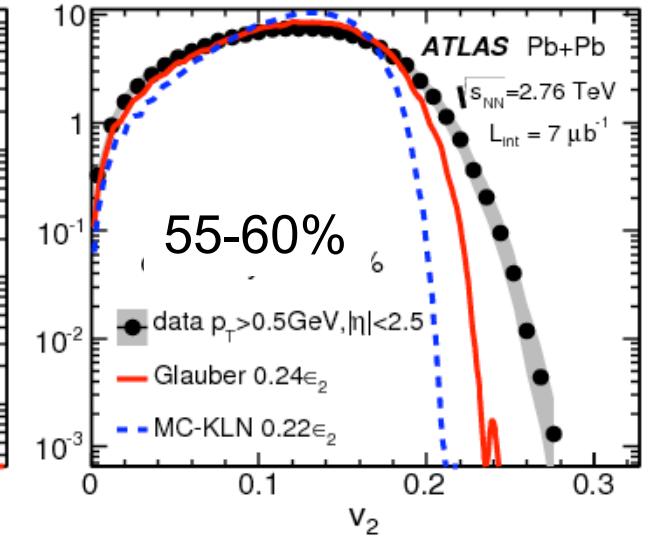
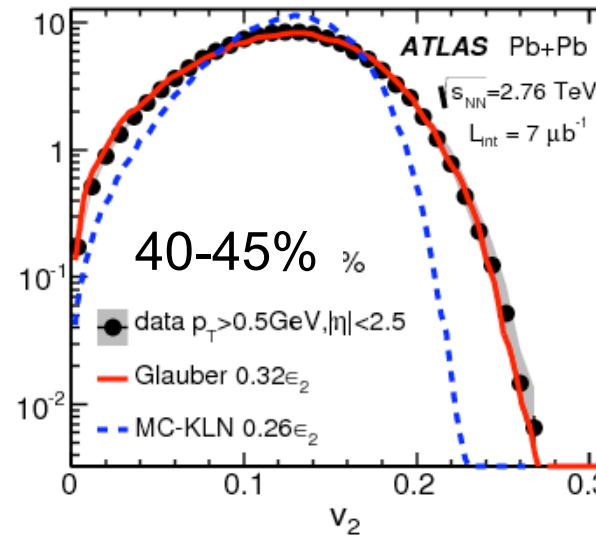
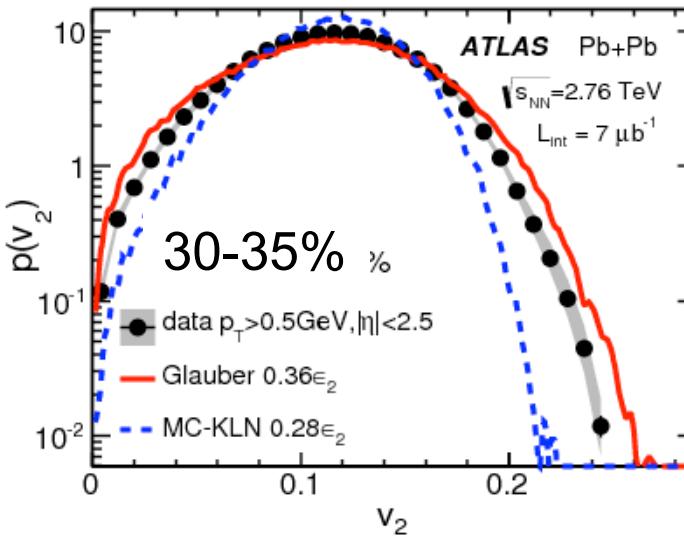
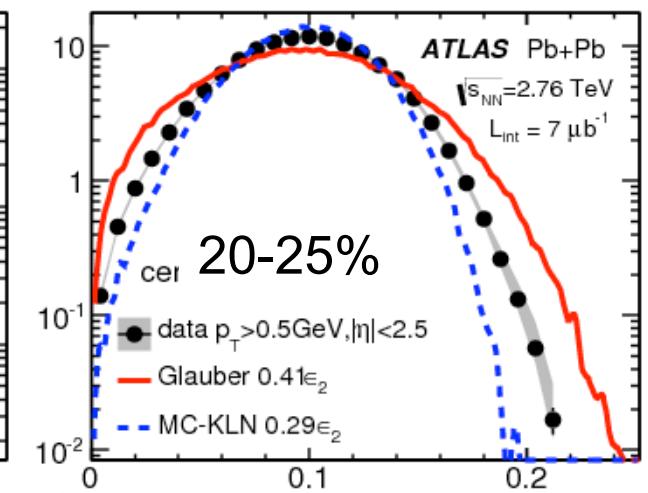
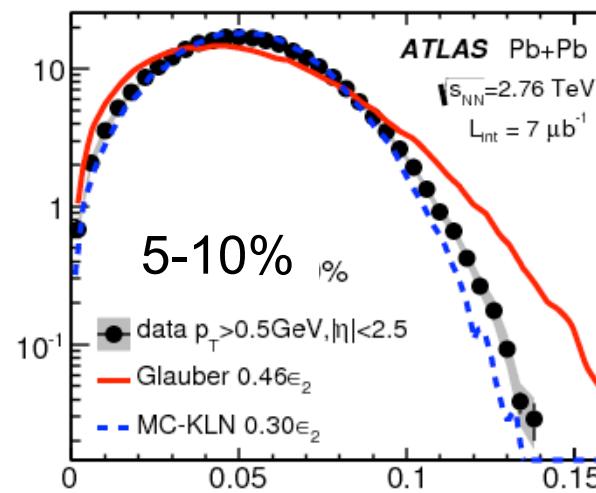
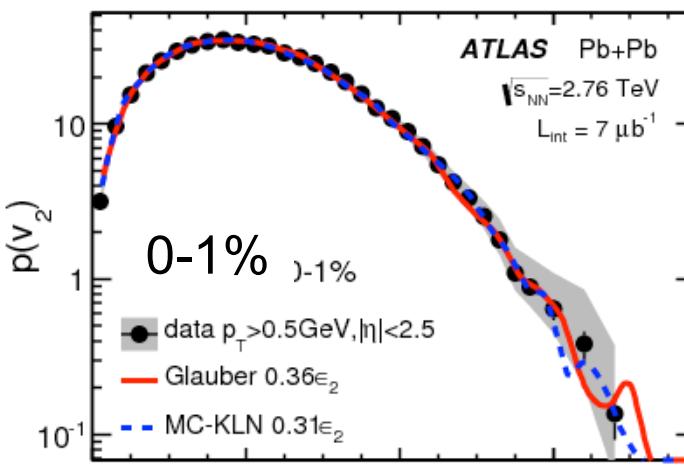
Compare with initial geometry models



$$v_n \propto \epsilon_n = \frac{\sqrt{\langle r^n \cos n\phi \rangle^2 + \langle r^n \sin n\phi \rangle^2}}{\langle r^n \rangle}$$

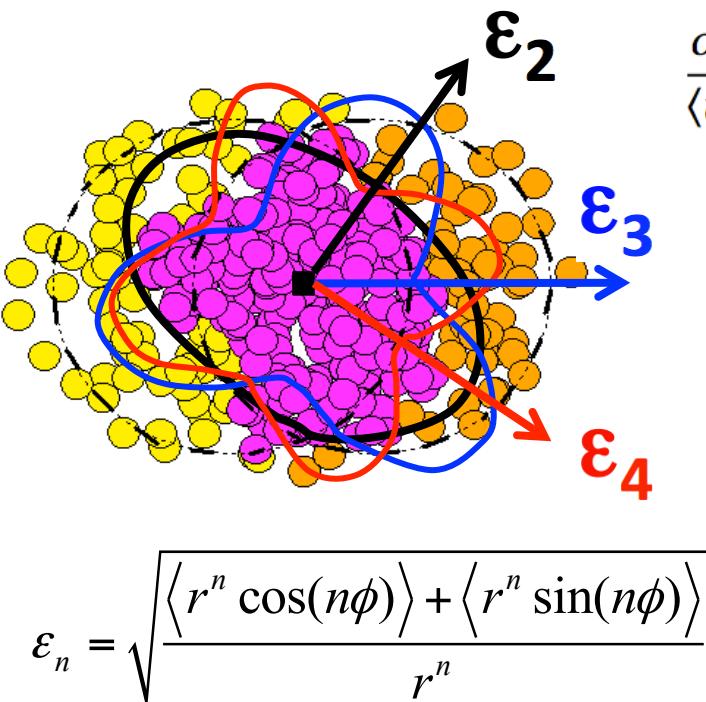
Glauber and CGC mckln

Rescale ϵ_n distribution to the mean of data

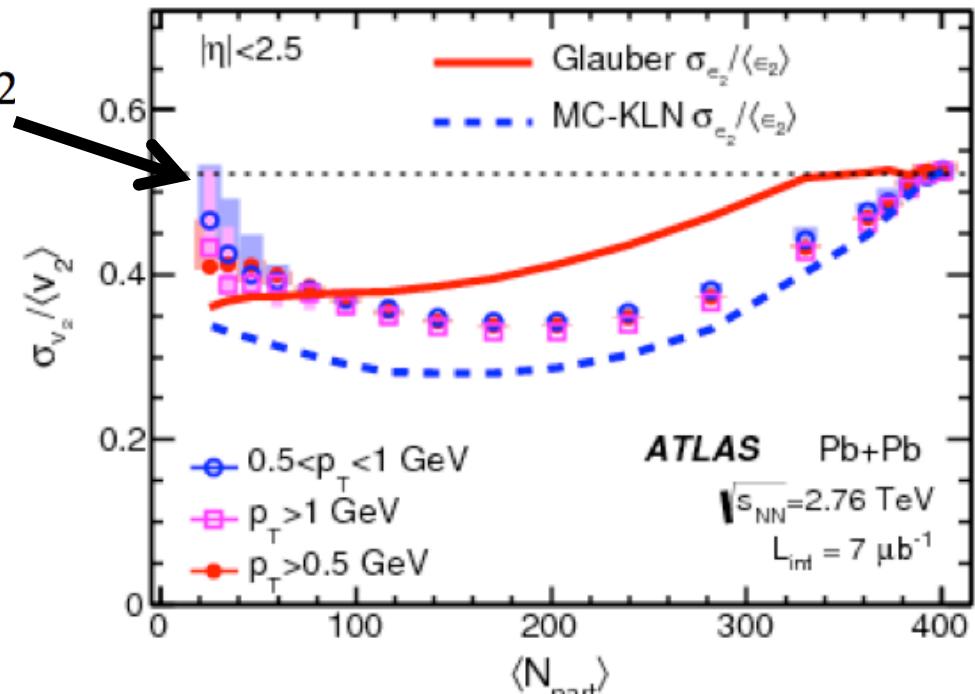


Both models fail describing $p(v_2)$ across the full centrality range

Compare with initial geometry models

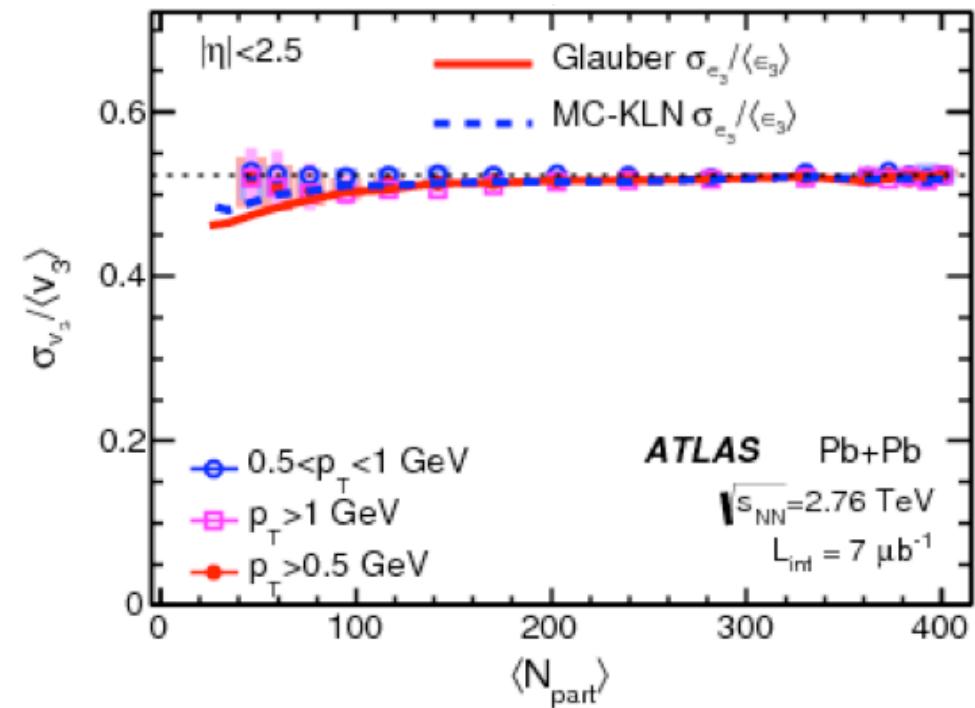


$$\sigma_{v_n} = \sqrt{\frac{4}{\pi} - 1} = 0.52$$



- Test relation

$$\frac{\sigma_{\varepsilon_n}}{\langle \varepsilon_n \rangle} = \frac{\sigma_{v_n}}{\langle v_n \rangle}$$

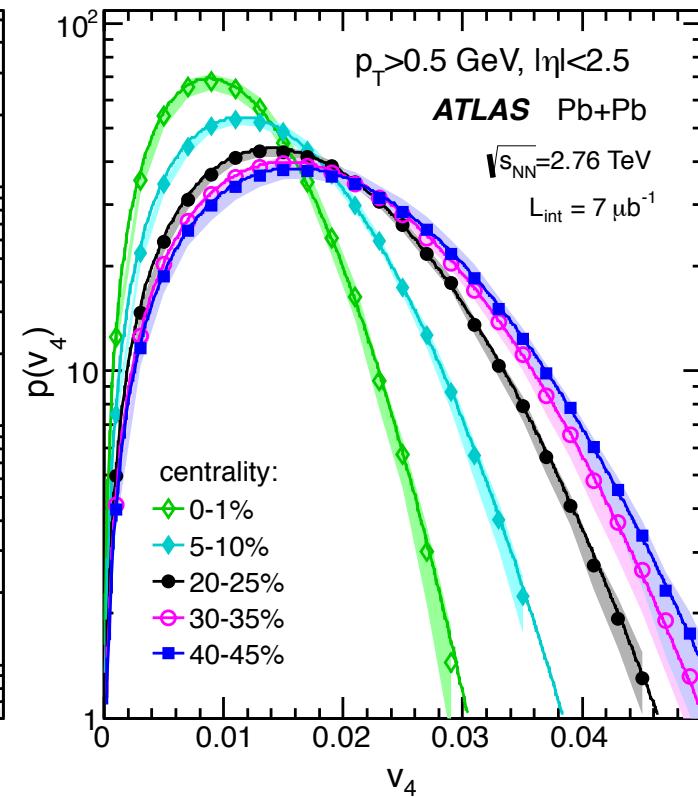
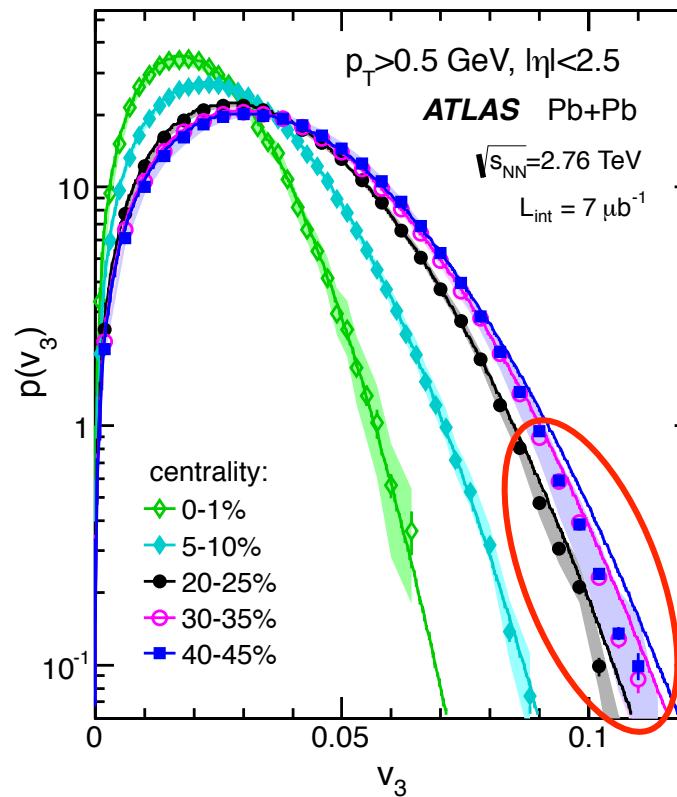
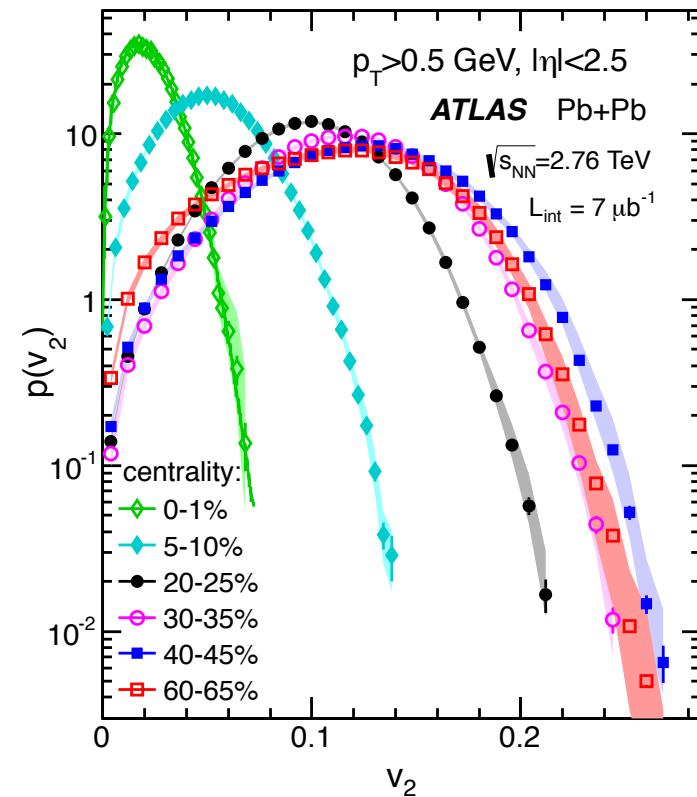


Both models failed.

$p(v_2)$, $p(v_3)$ and $p(v_4)$ distributions

Parameterize with pure fluctuation scenario:

$$p(v_n) \propto v_n \exp\left(\frac{-v_n^2}{2\delta_n^2}\right)$$

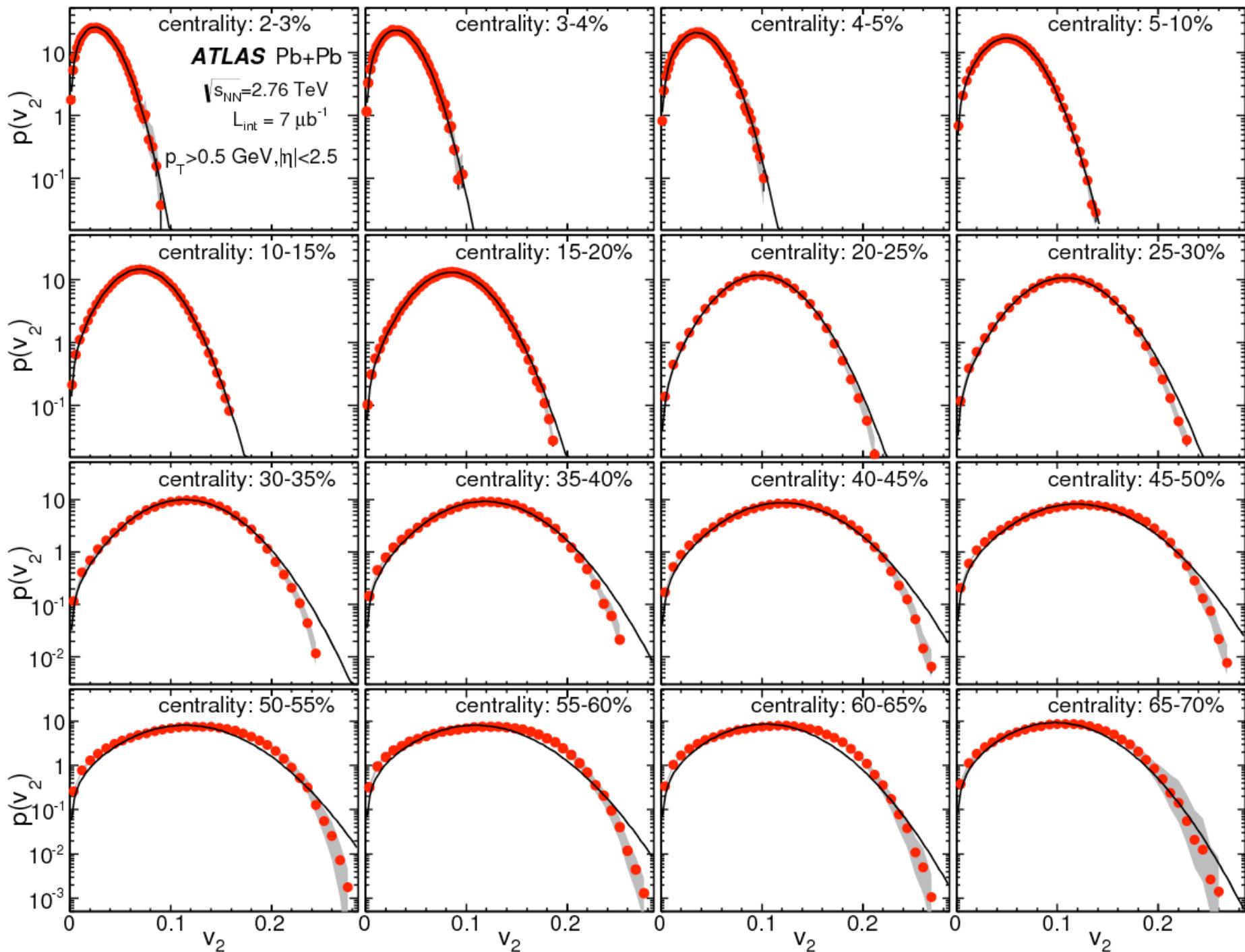


Only works in 0-2% centrality
Require non zero v_2^{RP} in others

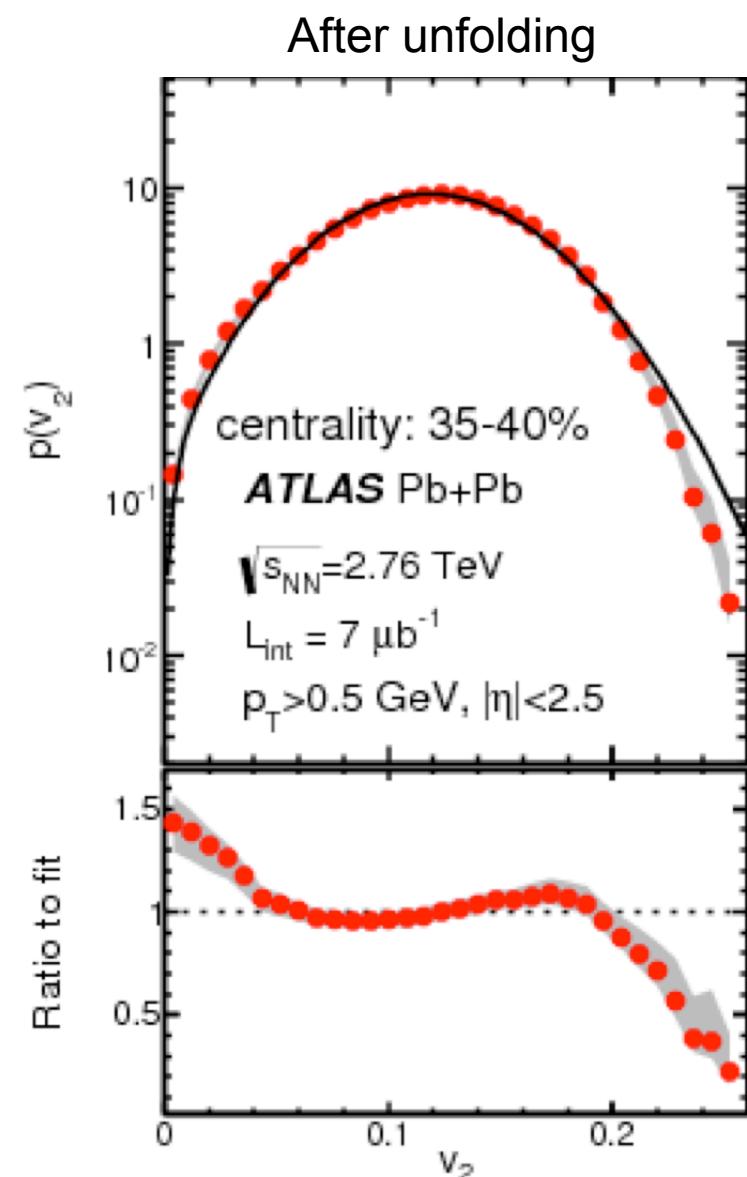
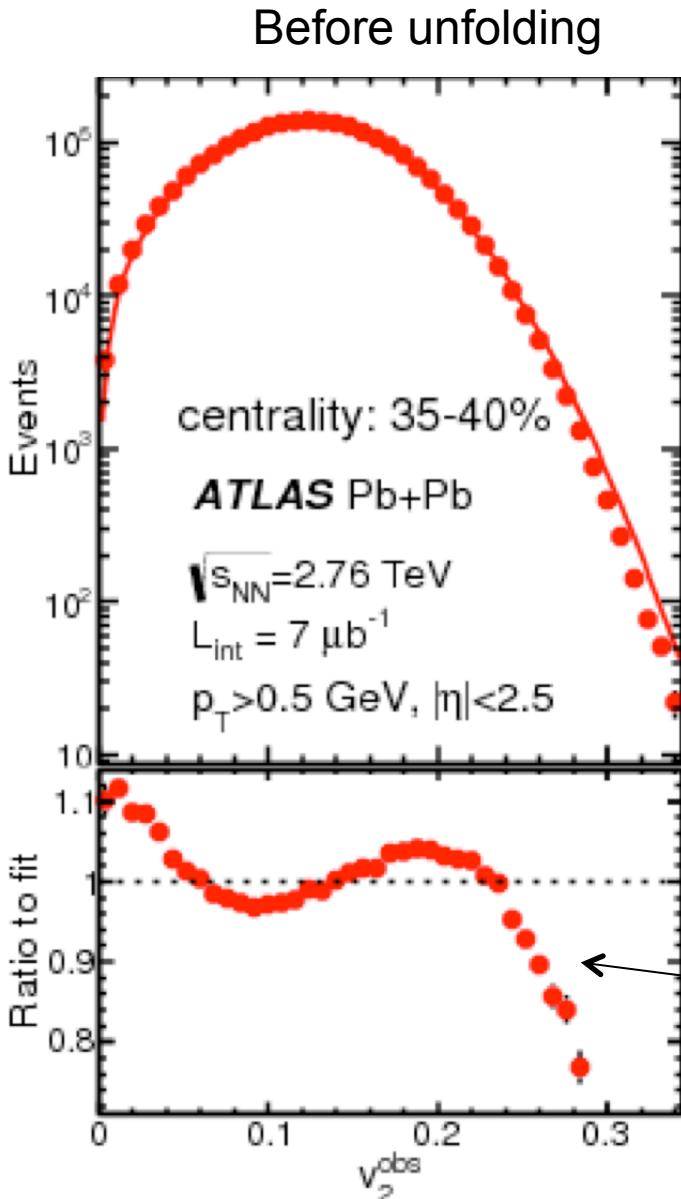
Deviations in the tails:
non zero v_3^{RP}

No deviation is observed,
however v_4 range is limited.

Bessel-Gaussian fit to $p(v_2)$

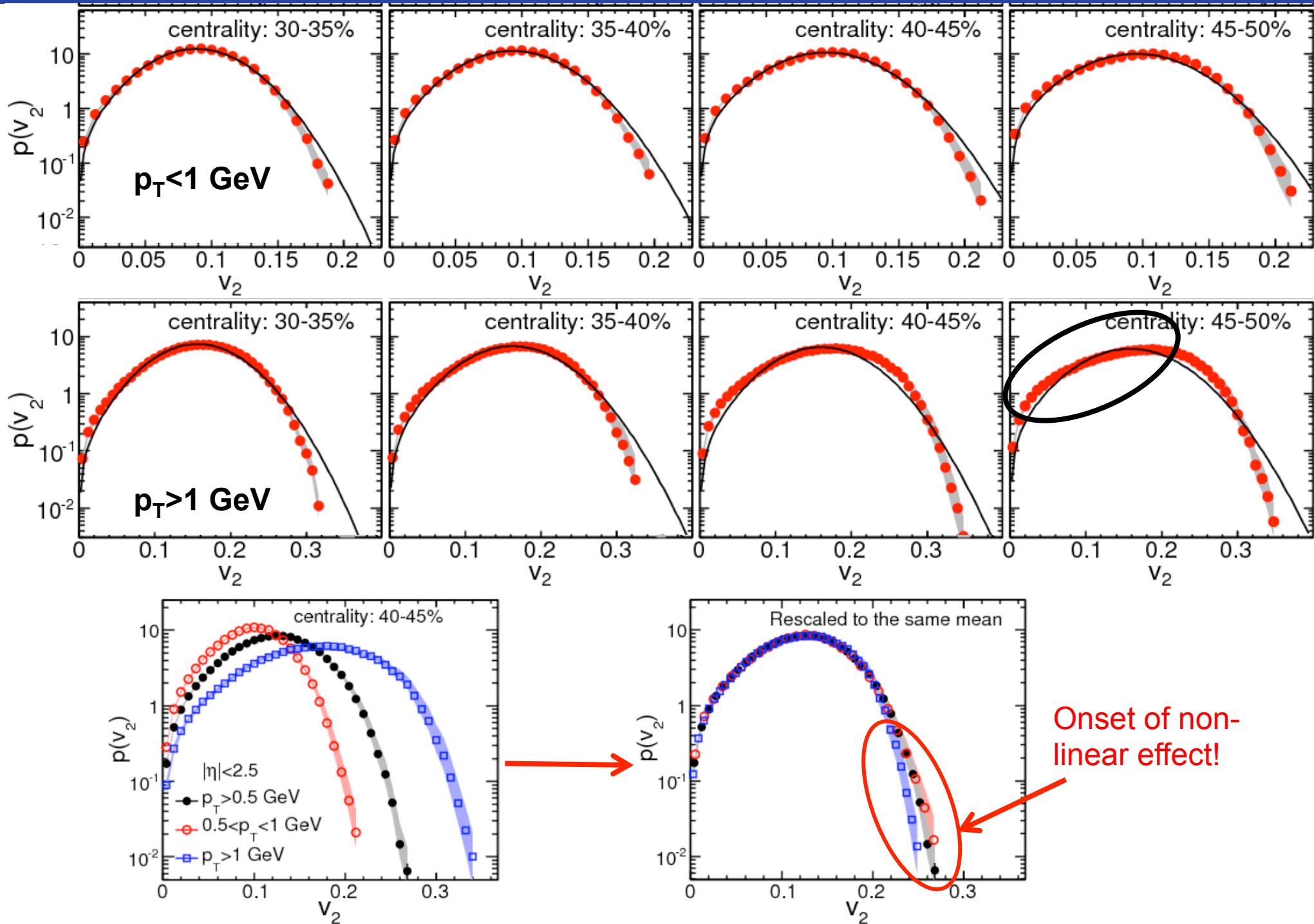


The deviation from B-G before/after unfolding



Deviation from BG is
already clear before
unfolding

Deviation grows with p_T



Are cumulants sensitive to these deviations?

$$v_2^{\text{calc}}\{4\}^4 \equiv -\langle v_2^4 \rangle + 2\langle v_2^2 \rangle^2 \approx (v_2^{\text{RP}})^4 ,$$

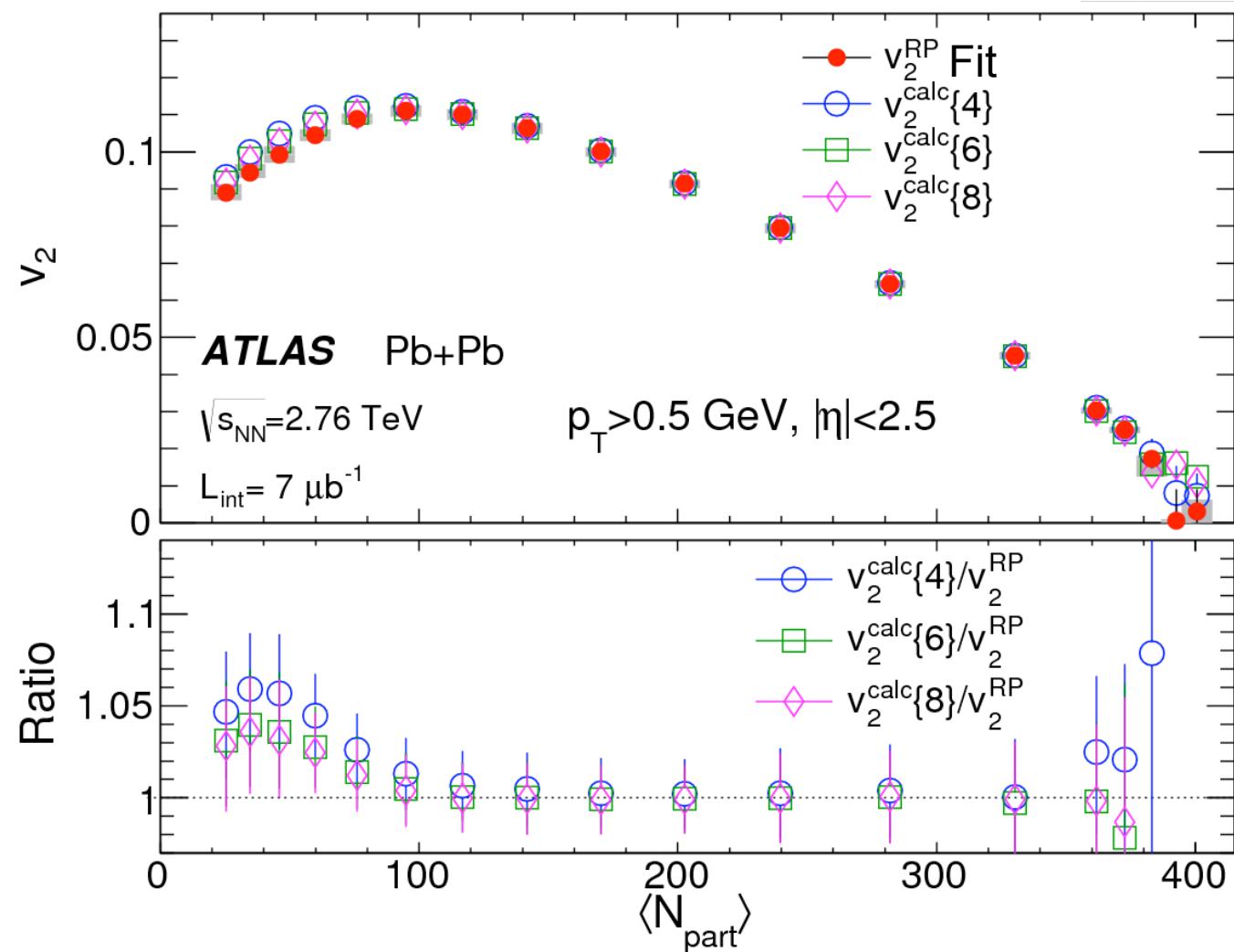
$$v_2^{\text{calc}}\{6\}^6 \equiv (\langle v_2^6 \rangle^2 - 9\langle v_2^4 \rangle\langle v_2^2 \rangle + 12\langle v_2^2 \rangle^3) / 4 \approx (v_2^{\text{RP}})^6 ,$$

$$v_2^{\text{calc}}\{8\}^8 \equiv -(\langle v_2^8 \rangle^2 - 16\langle v_2^6 \rangle\langle v_2^2 \rangle - 18\langle v_2^4 \rangle^2 + 144\langle v_2^4 \rangle\langle v_2^2 \rangle^2 - 144\langle v_2^2 \rangle^4) / 33 \approx (v_2^{\text{RP}})^8$$

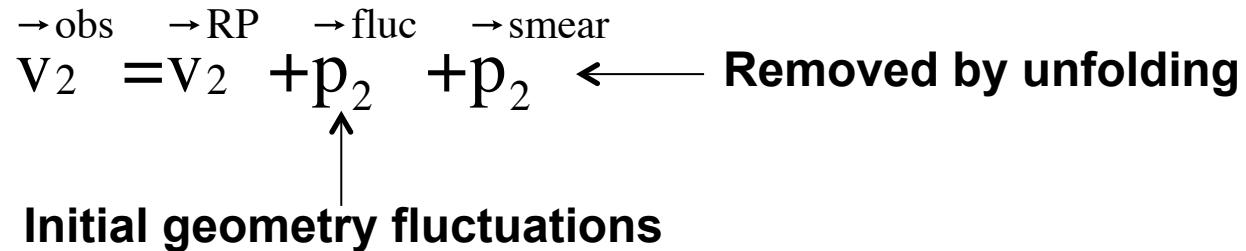
Cumulants are not sensitive to the tails, presumably because their values are dominated by the v_2^{RP} :

$$v_2\{6\}^6 \sim (v_2^{\text{RP}})^6 + (\Delta)^6$$

If $\Delta=0.5 v_2^{\text{RP}}$, $v_2\{6\}$ only change by 2%



$v_2\{4,6,8\}$ without unfolding



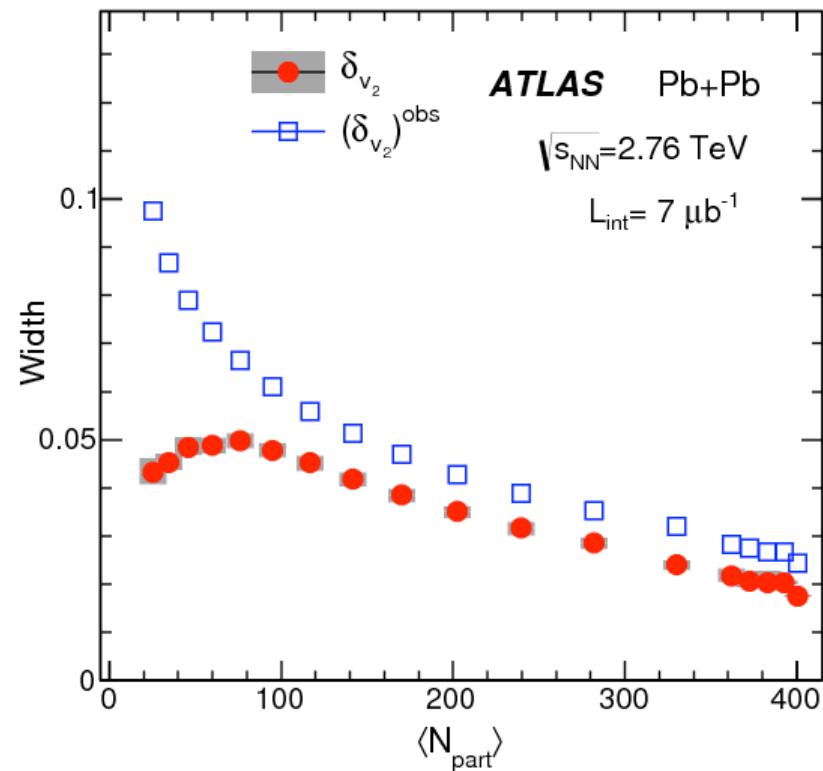
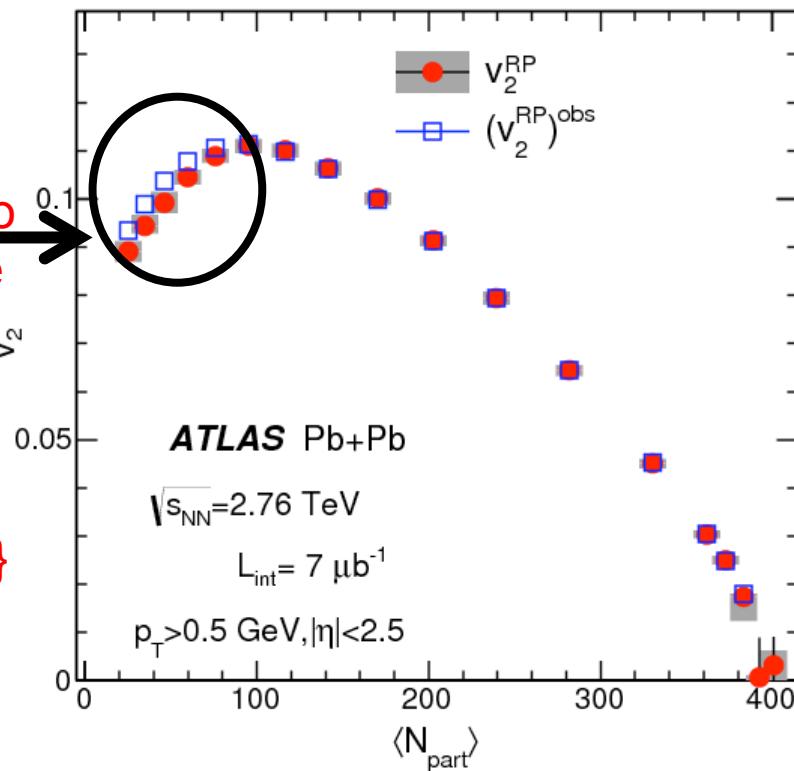
$$p(\vec{v}_2) \propto \exp\left(\frac{-(\vec{v}_2 - \vec{v}_2^{RP})^2}{2\delta_2^2}\right)$$

Additional Gaussian smearing won't change the v_2^{RP} .

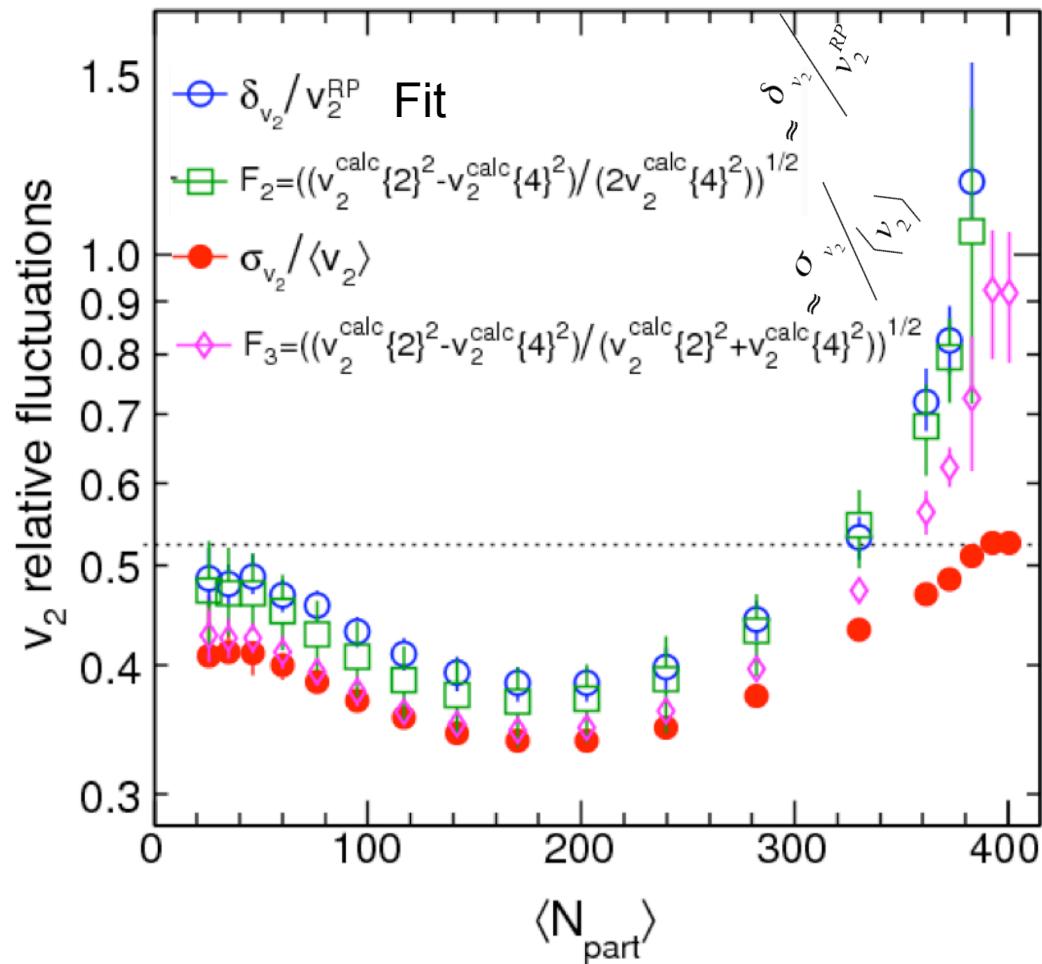
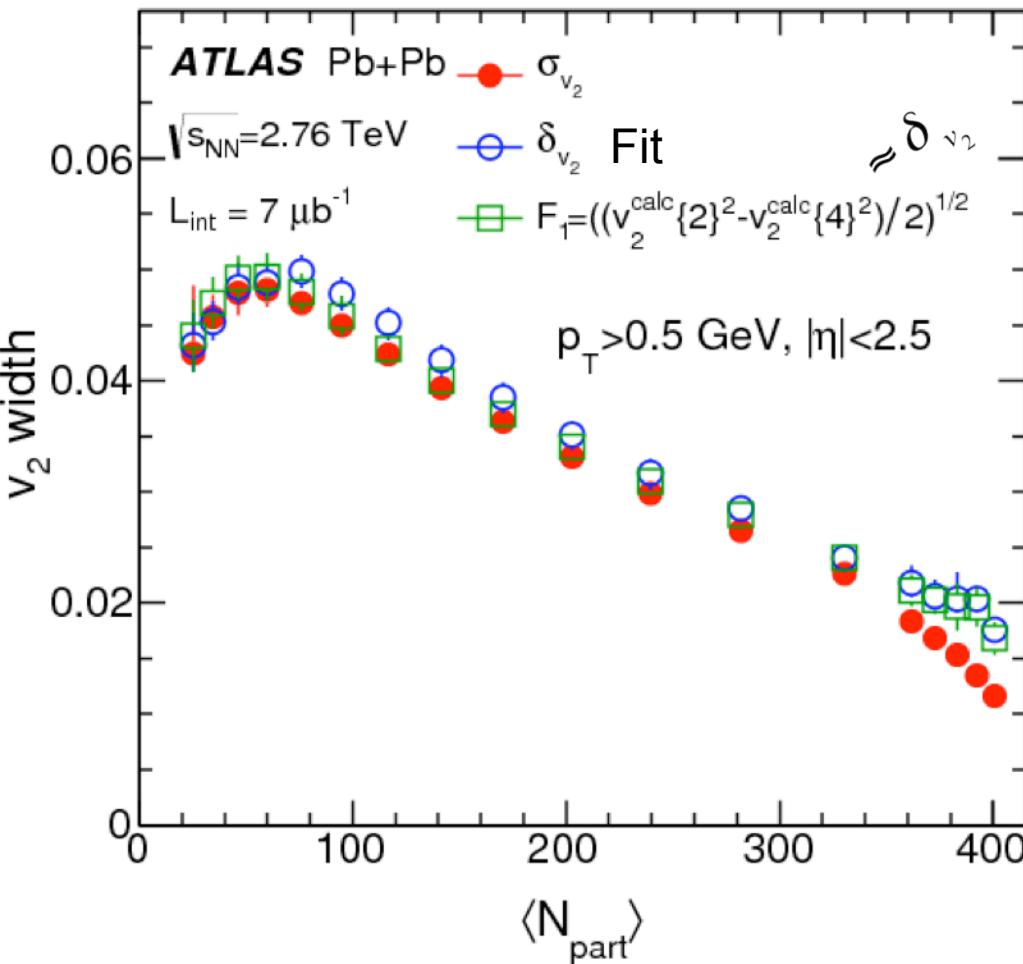
Indeed the response function is Gaussian

Response function
 (nonflow+noise) is no
 longer Gaussian, the
 real flow fluctuations
 may also be non-
 Gaussian.

 The meaning of $v_2\{4\}$
 is non-trivial in this
 limit (also in pPb)

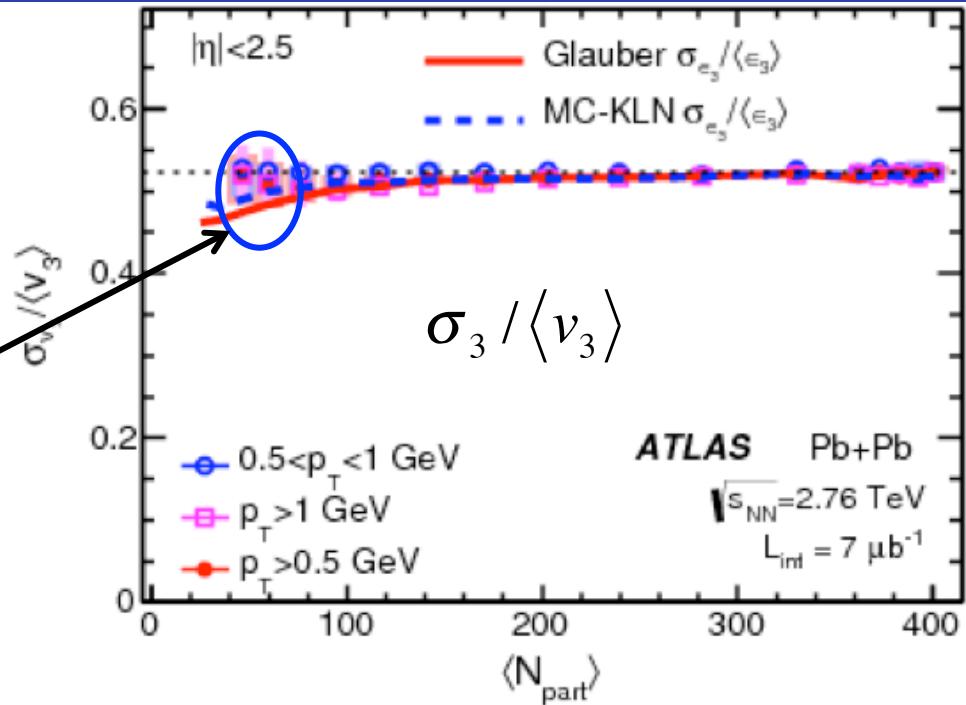
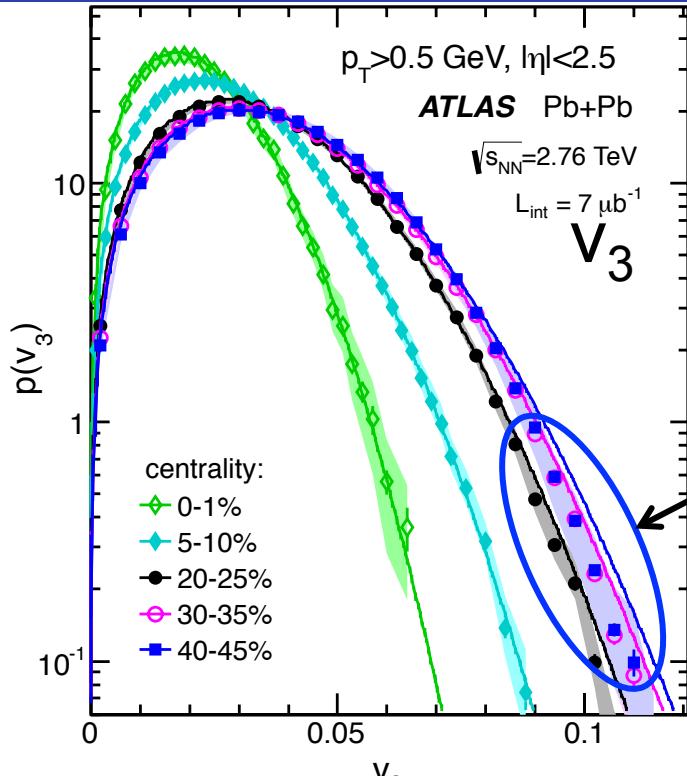


Extracting relative fluctuations



- Different estimator gives different answer, especially in central collisions
 - Expected since they have different limit.
 - Stick to one convention?

Flow fluctuation & $v_3\{4\}$



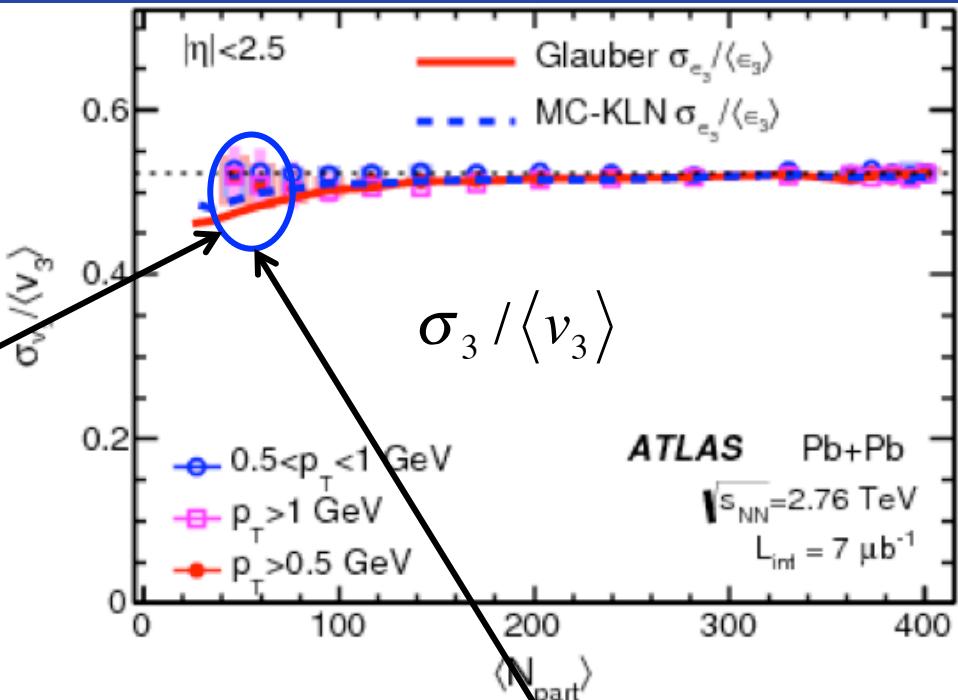
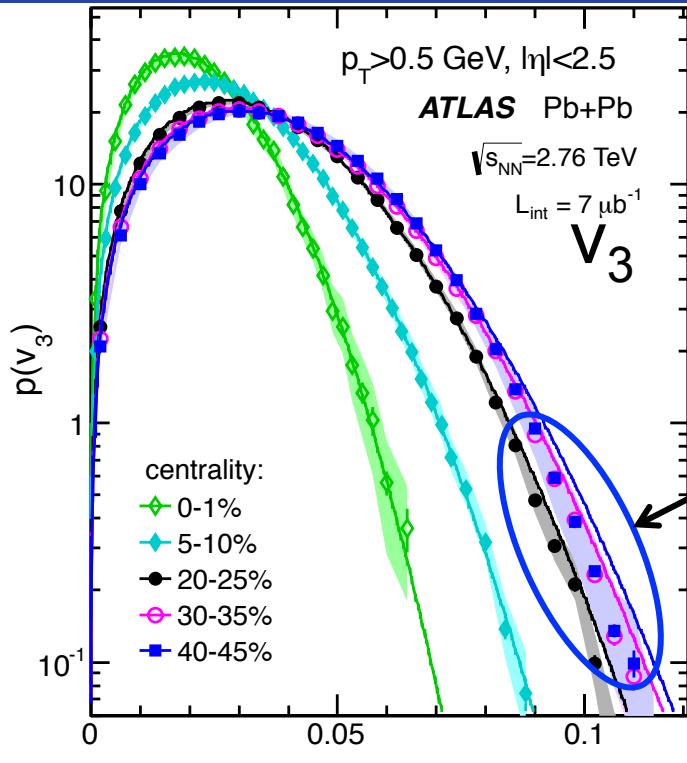
- Even a small deviation will imply a v_n^{RP} or $v_n\{4\}$ value comparable to δ_{vn}

$$v_n\{4\} = \left[2 \left\langle v_n^2 \right\rangle^2 - \left\langle v_n^4 \right\rangle \right]^{1/4}$$



a 4% difference gives a $v_n\{4\}$ value of about 45% of $v_n\{2\}$

Flow fluctuation & $v_3\{4\}$

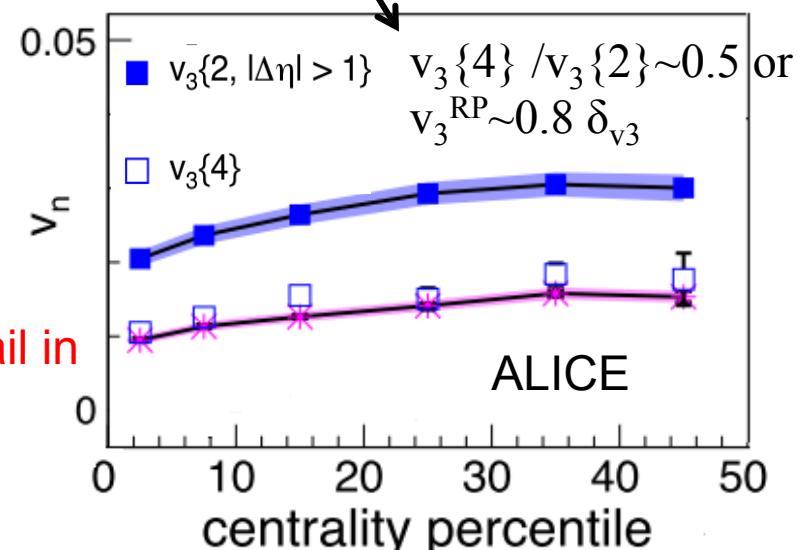


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a 4% difference gives a $v_n\{4\}$ value of about 45% of $v_n\{2\}$

Due to a non-Gaussian tail in the $p(v_3)$ distribution?



Summary

- Event-by-event fluctuation of the QGP and its evolution can be accessed via $p(v_n, v_m, \dots, \Phi_n, \Phi_m, \dots)$
- Detailed correlation measurement 2- and 3- event planes → the Fourier coefficients of $p(\Phi_n, \Phi_m)$ and $p(\Phi_n, \Phi_m, \Phi_L)$
 - Strong proof of mode-mixing/non-linear effects of the hydro response to initial geometry fluctuations.
 - New set of constraints on geometry models and η/s .
- First measurements of the $p(v_2)$, $p(v_3)$ and $p(v_4)$
 - Glauber and MC-KLN models ruled out
 - $p(v_2)$ show significant deviation of the fluctuation from Gaussian, also suggestive of strong non-linear effects.
 - $v_2\{4,6,8\}$ are not sensitive to these deviations, except in peripheral collisions.
 - $p(v_3)$ distribution suggests a non-zero v_3^{RP} .
- Look into other correlations.