Spatiotemporal Dynamics of Optical Pulse Propagation in Multimode Fibers

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- www.osa.org/tgwebinars/#tab_ondemand
- Spatiotemporal Dynamics of Optical Pulse Propagation in Multimode Fibers
 - Prof. Frank Wise of Cornell University
 - Theoretical and experimental studies of the basic properties and spatiotemporal behavior of complex nonlinear dynamics in multimode fiber will be presented.





Spatiotemporal Dynamics of Optical Pulse Propagation in Multimode Fibers



F. W. Wise Department of Applied Physics Cornell University





Introduction to nonlinear pulse propagation

Recent progress in multimode nonlinear propagation

- Solitons in multimode GRIN fiber: formation and fission
- Multimode continuum generation
- Spatiotemporal dispersive waves
- Spatiotemporal modulation instability
- Beam self-cleaning

Future directions / toward applications





- Pulse propagation in multimode fiber is spatiotemporally complex
- 4D vector field



Our job is to figure out basic processes, building blocks, and "rules"





Introduction to Nonlinear Wave Propagation





 $n = n(\omega)$ $v(\omega) = c/n(\omega)$







Dispersive phase accumulation







anomalous dispersion $\lambda > 1300$ nm for silica



Dispersive phase accumulation







normal dispersion $\lambda < 1300$ nm for silica







self-phase modulation produces new frequencies







cross-phase modulation produces new frequencies





$$\frac{\partial A(z,t)}{\partial z} + i \frac{\beta^{(2)}}{2} \frac{\partial^2 A(z,t)}{\partial t^2} = i\gamma |A(z,t)|^2 A(z,t)$$



(anomalous) dispersion cancels nonlinearity for

$$A(t) = A_0 \operatorname{sech}(t/\tau_p) \exp(iz/z_{sol})$$





$$\frac{\partial A(z,t)}{\partial z} + i \frac{\beta^{(2)}}{2} \frac{\partial^2 A(z,t)}{\partial t^2} = i\gamma |A(z,t)|^2 A(z,t)$$



(anomalous) dispersion cancels nonlinearity for

$$A_0 \tau_p = \sqrt{\frac{\beta_2}{\gamma}}$$









Linear wave propagation

pulse spreads owing to group-velocity dispersion



beam spreads owing to diffraction







nonlinear phase shift produces self-focusing



nonlinear phase shift produces self-focusing





• diffraction balances self-focusing for $P = P_{cr} \sim 5 \text{ MW in glass}$







• diffraction balances self-focusing for $P = P_{cr} \sim 5 \text{ MW in glass}$



2D: unstable against collapse





A continuous wave breaks into temporal components





Why are solitons so important?

In general, waves in nonlinear media are unstable



Modulation Instability



- A beam breaks into its component solitons
- Stable products of instability are "eigenmodes" of nonlinear systems







If they exist, solitons are important

- as stable wave packets (sometimes nonlinear attractors)
- as components of arbitrary fields

In 1D solitons underlie

- modelocked lasers
- continuum generation
- breathers, Peregrine soliton
- rogue waves
- ...

2D and 3D: solitons are unstable



Multimode waveguides: between 1- and 3-D





https://commons.wikimedia.org/wiki/File:Optical_fiber_types.svg





- Little work on multimode nonlinear pulse propagation before 2013
- Recent theoretical, computational advances e.g., transfer matrix, principal modes,...

Relevance to multicore fibers



Huang et al., Opt Exp 2014



Why study propagation in multimode fiber now?



Laser/ amplifier / transmission applications

- Spatial division multiplexing in telecom

Agrell et al., J Opt 2016

Imaging through multimode fiber/ complex media



Ploschner et al., Nature Photon 2015









$$n^{2}(\rho) = n_{0}^{2} \left[1 - 2\Delta \left(\frac{\rho}{R} \right)^{\alpha} \right], \quad \rho \leq R$$
$$= n_{0}^{2} (1 - 2\Delta), \quad \rho > R$$













- Propagation constants equally-spaced
- Velocities of modes vary much less than in step-index fiber



532 nm



What should we measure?



Broadband space-time diagnostic does not exist



- Record overall average spectrum to compare to calculated
- Image near-field on autocorrelator
- Compute spatiotemporal autocorrelation for comparison





Multimode Solitons



Linear propagation







Multimode soliton formation











Excite 3 lowest modes










- For E < 0.1 nJ pulse disperses</p>
- 0.5 nJ pulse energy







• Launch 0.5 nJ / 300 fs



• Coupled-mode theory and beam-propagation give similar results



Renninger et al., Nature Commun 2013

- Solitons with more modes require greater nonlinear phase / energy
- Solitons with up to 10 modes generated



Multimode soliton formation







Multimode soliton formation







Multimode soliton fission







Simulation

- Smaller peaks in AC from less-localized modes
- Intermodal energy transfer during, after fission

Experiment





Multimode soliton fission: experiment





Simulation

Experiment





- Fission produces multiple MM solitons and MM dispersive waves
- Fission is spatiotemporal
- Raman "focuses" energy into the low-order mode

Wright et al., Opt Express 2015





Continuum Generation





500 fs energy up to 1 μJ peak power up to MW 1550 nm































Spatial conditions determine the continuum











Perturbation of solitons (1D tutorial)

• Perturbed soliton adjusts to reach $A_0 \tau_p = \sqrt{\frac{\beta_2}{\gamma}}$

and radiates dispersive wave

Periodic perturbation (period = Z_c)
Resonant energy transfer when wave vectors match

$$(k_{sol} - k_{dis}) = 2m\pi/Z_c$$

$$\Omega_{res} = \frac{1}{\tau} \sqrt{\frac{8Z_o m}{Z_c} - 1}$$



Gordon, J Opt Soc Am B 1992





Spatiotemporal oscillations











Simulation, experiment and analytic theory agree well

Wright et al., Phys Rev Lett 2015



Oscillations about equilibrium as an instability: why more degrees of freedom matters









- Continuum is controllable through launched spatial modes
- Spatiotemporal oscillation leads to the generation of multimode dispersive waves
- Phenomenon understood in terms of multimode soliton dynamics

Wright et al., Nature Photon 2015 Wright et al., Phys Rev Lett 2015





Spatiotemporal Modulation Instability



Spatiotemporal modulation instability



Launch continuous wave or long pulse at normal dispersion









- Periodic self-imaging plays a role
- Instability occurs for either sign of dispersion

Longhi, Opt Lett 2003 Matera et al., Opt Lett 1993 Nazemosadat et al., JOSA B 2016



Spatiotemporal MI in GRIN fiber





Krupa et al., Phys Rev Lett 2016





 Geometric parametric instability: periodic self-imaging of field allows quasi-phase-matching of 4WM sidebands

Krupa et al., Phys Rev Lett 2016





Beam Self-Cleaning in Multimode Fiber





Krupa et al., arXiv 2016



Krupa et al., arXiv 2016



Krupa et al., arXiv 2016



Beam self-cleaning in GRIN fiber



Linear dynamics





Simulations show that Kerr nonlinearity underlies self-cleaning

Krupa et al., arXiv 2016



Lopez-Galmiche et al., Opt Lett 2016




- Continuum from spatiotemporal MI, geometric parametric instability, Raman, and other 4-wave mixing processes
- Self-cleaning confirmed
- Speckle-free output with moderate M²



- 80 μJ pulse energy
- Route to compact, bright, multi-octave continuum

Lopez-Galmiche et al., Opt Lett 2016



Z. Liu et al., 2016





- P < P_{cr}
- Negligible dissipation



Temporal coherence maintained

Z. Liu et al., 2016







- Kerr nonlinearity underlies self-cleaning
- Process independent of pulse duration





Implications / Future Directions



Solitons in few-mode fibers



Mode-resolved studies



Nicholson et al., JSTQE 2009





Classical wave condensation



Wave turbulence theory

random optical waves can "thermalize"



- initial incoherent field self-organizes to form large coherent structure
- equipartition of energy in higher-order modes

 2D + parabolic waveguide: condensation predicted theoretically



Aschieri et al., Phys Rev A 2011



- Optical wave turbulence studied in 1D systems
- True turbulence requires 3D



Effects of disorder and dissipation



 Introduce random mode coupling gain, loss



- Complex system
- Controllable and measurable
- Testbed for

cooperative phenomena self-organized critical behavior

Wright et al., arXiv 2016



Relevance to telecommunications



■ N modes → N channels



Multimode solitons versus independent channels

Strongly-coupled mode groups: Manakov solitons



Instabilities may limit transmission



Relevance to telecommunications



Multimode fibers are small-world networks

- Coupling is primarily between nearest neighbors
- "Shortcut" links can lead to a strong-coupling transition, manymode self-organization



A small-world network Strogatz, Nature 2001

 Need to understand many-mode nonlinear interactions Mode-dependent gain and loss Mode-dependent, longitudinally-varying disorder



A multimode fiber laser is a new environment for nonlinear waves. It adds

- spatially-dependent gain, saturable absorption
- spatial and spectral filtering





Multimode fiber lasers can have much higher energy than single-mode fiber lasers

Larger mode area



Multimode soliton lasers



Multimode fiber lasers can have much higher energy than single-mode fiber lasers

Larger mode area

$$E \sim A_{eff}$$

single mode fiber large-mode-area microstructure fiber single higher-order mode multimode fiber $\begin{aligned} A_{eff} &= 50\text{-}100 \ \mu\text{m}^2 \\ A_{eff} &\sim 5,000 \ \mu\text{m}^2 \\ A_{eff} &\sim 3,000 \ \mu\text{m}^2 \\ A_{eff} &> 30,000 \ \mu\text{m}^2 \end{aligned}$

(1550 nm)





Multimode fiber lasers can have much higher energy than single-mode fiber lasers

- Larger mode area
- Modal dispersion

 $E \sim \sum$ dispersion





Multimode soliton lasers



Multimode fiber lasers can have much higher energy than single-mode fiber lasers

- Larger mode area
- Modal dispersion
- New (spatiotemporal) pulse evolutions

Role of spatiotemporal instabilities?

Ultimate limit from self-focusing





- Multimode fiber supports a variety of new spatiotemporal phenomena
- Initial results indicate that multimode solitons will help understand complex dynamics
- Relevance of nonlinear dynamics to applications
 - High-power, multi-octave continuua
 - Connection to optics of complex media
 - Space-division multiplexing in telecommunications
 - Laser / amplifier / transmission applications



























Reserve slides





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GRIN



In GRIN fiber, modes have similar group velocities





$$\frac{\partial A}{\partial z} = \frac{i}{2k_0} \left(\frac{\partial^2 A}{\partial x^2} + \frac{\partial^2 A}{\partial y^2} \right) - i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial t^2} - i \frac{k_0 \Delta}{R^2} (x^2 + y^2) A + i\gamma |A|^2 A$$

diffraction

dispersion

index profile

Kerr





$$\frac{\partial A}{\partial z} = \frac{i}{2k_0} \left(\frac{\partial^2 A}{\partial x^2} + \frac{\partial^2 A}{\partial y^2} \right) - i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial t^2} - i \frac{k_0 \Delta}{R^2} (x^2 + y^2) A + i\gamma |A|^2 A$$

Gross-Pitaevskii equation





"GMMNLSE"

$$\partial_z A_p(z,t) = i \left(\beta_0^{(p)} - \Re\left[\beta_0^{(0)}\right]\right) A_p - \left(\beta_1^{(p)} - \Re\left[\beta_1^{(0)}\right]\right) \frac{\partial A_p}{\partial t} + \sum_{m=2}^3 i^{m+1} \frac{\beta_m}{m!} \partial_t^m A_p$$

modal wavenumber mismatch modal velocity mismatch group velocity dispersion

$$+i\frac{n_{2}\omega_{o}}{c}\left(1+\frac{i}{\omega_{o}}\partial_{t}\right)\sum_{l,m,n}\left\{(1-f_{R})S_{plmn}^{k}A_{l}A_{m}A_{n}^{*}+f_{R}A_{l}S_{plmn}^{R}\int_{-\infty}^{t}d\tau A_{m}(z,t-\tau)A_{n}^{*}(z,t-\tau)h_{R}(\tau)\right\}$$
shock
Kerr
Raman

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A. Mafi, "Pulse Propagation in a Short Nonlinear Graded-Index Multimode Optical Fiber," J. Lightwave Technol. 30, 2803–2811 (2012).







Fig. 1. Relation between the LP modes and the real waveguide modes HE_{11x} , HE_{11y} , TE_{01} , TM_{01} , HE_{21a} , and HE_{21b} of the six-mode FMF.

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