Combined Refraction/Diffraction Model for Spectral Wave Conditions

REF/DIF S

Version 1.1

Documentation and User's Manual

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Introduction

The weakly nonlinear combined refraction and diffraction model described here, denoted **REF/DIF** S, simulates the behavior of a random sea over irregular bottom bathymetry, incorporating the effects of shoaling, refraction, energy dissipation and diffraction.

Although the model is developed to simulate a random sea state it can also be used to model the behavior of monochromatic waves. Additionally, a preprocessor is included which generates a two-dimensional spectrum consisting of a frequency spectrum in conjunction with a directional spreading function, discretized to yield separate wave components characterized by a frequency and a direction. These individual wave components make up the essential part of the input to the model REF/DIF S.

The REF/DIF S model has the ability to stack up parabolic models for the individual wave components, propagating them simultaneously through the domain. This approach makes it possible to calculate statistical wave parameters after each forward step in space in the parabolic scheme.

Taking advantage of this feature, REF/DIF S also incorporates a new statistical breaking model. The significant wave height can be computed using solutions for all wave components and can be used to predict the energy dissipation due to breaking.

The model REF/DIF S includes most features present in the model REF/DIF 1 developed by Kirby and Dalrymple (1992). The reader is referred to that document for a discussion of the historical development of the model. The model REF/DIF S is constructed in parabolic form, and thus there is a restriction of the model to cases where the propagation direction is within $\pm 45^{\circ}$ of the assumed mean wave direction.

The REF/DIF S model is described in detail in this manual. The first chapter includes a short summary of the theoretical background. First, the discretization of a two-dimensional spectrum into discrete wave components is described, then the development of wave models for these individual components is described. The assumptions adopted by the model are pointed out. Then, the energy dissipation model and the breaking model are described, followed by comments on the calculation of radiation stresses and the numerical scheme.

The second chapter includes the operating manual. An overview of the operating manual can be found in Section 2.2. Chapter 3 documents the application of the model to actual examples and provides explicit descriptions of the input and output. Appendix A includes the REF/DIF S program listing, in Appendix B the listing of a plotting routine is given. This program can be used to construct contour plots of the wave field. Appendix C includes the fweb documentation of the SPECGEN program listing. This program discretizes a

standard two-dimensional design spectrum and creates the input file for REF/DIF S. Appendix D includes frequently asked questions and the answers to those questions.

1 THEORETICAL BACKGROUND

1.1 Wave Climate

Although REF/DIF S can be used with monochromatic wave trains, it is intended to be used with two-dimensional wave spectra. The input frequency spectrum in conjunction with a directional spreading function have to be divided into discrete wave components characterized by a certain frequency and direction. Thus the water surface elevation is represented as

 $\eta(x,y,t) = \sum_{all\,f} \sum_{all\,\theta} \left\{ \frac{A(x,y;f,\theta)}{2} e^{i\psi} + c.c. \right\}$ (1)

where f is the frequency, θ is the direction of any individual wave component and

$$\psi = \int \mathbf{k} \cdot d\mathbf{x} - \omega t \tag{2}$$

Discrete amplitudes and directions of the input waves have to be provided by the user on the offshore grid row, corresponding to x = 0. The wave is prescribed by the program as

$$A(0,y)_n = A_{n_0} e^{i\ell y} \tag{3}$$

where A_{n_0} is the given wave amplitude for the wave component n and ℓ is the wavenumber in the y-direction and is related to the wavenumber k by the relationship $\ell = k_n \sin \theta_n$, where θ_n is the angle made by the wave component n to the x-axis. Figure 1 illustrates the sign convention for the angle θ_n .

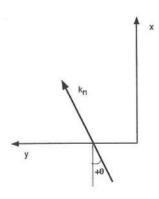


Figure 1: Sign Convention for Wave Angle θ_n .

The preprocessor **SPECGEN** subdivides a standard TMA spectrum and a corresponding wrapped normal directional spreading function into a desired number of discrete components and creates the input data for **REF/DIF** S. The *fweb* documentation of the preprocessor **SPECGEN** is given in Appendix C.

As mentioned earlier, computations are performed simultaneously for all wave components with each spatial step in the computational domain. Thus, after each step in the x-direction, the complex amplitudes $A(x,y)_n$ for all wave components are known. The significant wave height $H_{1/3}$ can then be computed in the following manner.

$$H_{1/3}(x,y) = \sqrt{8\sum_{n=1}^{N} |A(x,y)_n|^2}$$
(4)

where N is the total number of wave components and $A(x, y)_n$ is the amplitude of the wave component n.

The significant wave height $H_{1/3}$ is part of the output from the program and is also used in several places in the computational scheme, especially the wave breaking model to be discussed later in more detail.

1.2 Wave Models for Individual Spectral Components

1.2.1 Mild slope equation

The problem of water waves propagating over irregular bathymetry in arbitrary directions is a three-dimensional problem and involves complicated nonlinear boundary condition. Very few solutions to the three dimensional problem exist and those that do are only for flat bottoms. In one horizontal dimension, sophisticated models by Chu and Mei (1970) and Djordjevic and Redekopp (1978) predict the behavior of Stokes waves over slowly varying bathymetry. In order to simplify the problem in three dimensions, Berkhoff (1972), noted that the important properties of linear progressive water waves could be predicted by a weighted vertically integrated model. (The vertical integration reduces the problem to only the two horizontal dimensions, x and y.)

Berkhoff's equation is known as the mild slope equation. It is written in terms of the surface displacement, $\eta(x,y)$. The equation, in terms of horizontal gradient operator, is

$$\nabla_h \cdot (CC_q \nabla_h \eta) + k^2 CC_q \eta = 0 \tag{5}$$

Here,

$$C = \sqrt{(g/k) \tanh kh}$$
, the wave celerity,
 $C_g = C\{1 + 2kh/\sinh 2kh\}/2$, the group velocity,

where the local water depth is h(x,y) and g is the acceleration of gravity. The local wave number, k(x,y), is related to the angular frequency of the waves, ω , and the water depth h by the linear dispersion relationship,

$$\omega^2 = gk \tanh kh \tag{6}$$

The model equation (5) is an approximation; however, it is quite good even for moderately large local bottom slopes (see Booij, 1983). In both deep and shallow water, it is exact. Numerous authors have applied the mild slope model to various examples, primarily using finite element techniques. See, for example, Jonsson and Skovgaard (1979), Bettess and Zienkiewicz (1977), and Houston (1981).

For the linear mild slope equation, Radder (1979) developed a parabolic model, which had several advantages over the elliptic form presented by Berkhoff. First, the boundary condition at the downwave end of the model area is no longer necessary and, secondly, very efficient solution techniques are available for the finite difference form of the model. Radder used a splitting matrix approach, which involves separating the wave field into a forward propagating wave and a backward propagating wave, and then neglecting the backward scattered wave (which is justified in most applications as only the forward propagating wave is used for design). Radder's approximation for derivatives transverse to the wave direction results in a restriction on his parabolic model: the waves must propagate within approximately 30° of the assumed wave direction. Booij (1981) also developed a splitting of the elliptic equation, but his procedure included more terms in the approximation to the lateral derivative and therefore his procedure enables the parabolic model to handle waves propagating within 45° of the assumed direction. Booij's procedure has been used in previous versions of the REF/DIF 1 model (up through version 2.3).

More recently, Kirby (1986b) has developed an extension to the Booij approximation based on a Minimax principle, which further extends the range of validity of the model equations. The wave-current version of the resulting model is included here and in Version 2.4 of **REF/DIF** 1, however coefficients resulting in Booij's model are used.

1.2.2 Diffraction models

In contrast to the mild slope model which is valid for varying bathymetry, researchers in the area of wave diffraction were developing models for constant bottom applications. For example, Mei and Tuck (1980) developed a simple parabolic equation for wave diffraction and applied it to the diffraction of waves by a slender island. Their equation is

$$\frac{\partial A}{\partial x} = \frac{i}{2k} \frac{\partial^2 A}{\partial y^2} \tag{7}$$

where A is a complex amplitude related to the water surface displacement by

$$\eta = Ae^{i(kx - \omega t)} \tag{8}$$

Yue and Mei (1980), using a multiple scales approach, developed a nonlinear form of this equation, which accurately predicts the propagation of a third order Stokes wave. A striking result of their numerical experiments was the development of Mach stem reflection due to the reflection of obliquely incident waves from a breakwater. This phenomenon is uniquely a nonlinear effect and not predictable from a linear refraction theory.

The parabolic model described below combines the essential features of the two approaches described above. The variable depth features of the mild-slope equation (along with extensions to include effects of wave-current interaction) are retained, but the model is developed in parabolic form and in terms of a complex amplitude A.

1.2.3 Nonlinear combined refraction/diffraction models

Kirby (1983), using a Lagrangian approach, and Kirby and Dalrymple (1983a), with a multiple scales technique, developed the predecessor to the REF/DIF S model, which bridged the gap between the nonlinear diffraction models and the linear mild slope equation. This model can be written in several forms depending on the application. The hyperbolic form, for time dependent applications, and the elliptic form, for steady state problems, require the use of boundary conditions on all sides of the model domain. This is a difficult requirement, as the reflected wave at a boundary is not generally known a priori. These models, however, have the advantage that there is no restriction on the wave direction.

A detailed comparison of results of the weakly-nonlinear model of Kirby and Dalrymple (1983a) to laboratory data was shown by Kirby and Dalrymple (1984). The laboratory test, conducted at the Delft Hydraulics Laboratory by Berkhoff, Booij and Radder (1982), consisted of determining the wave amplitude over a shoal on a sloping bottom. While results predicted by ray tracing techniques were shown by Berkhoff, Booij and Radder to be very poor, the agreement between the weakly-nonlinear model and the laboratory data was excellent. Comparisons between linear and nonlinear parabolic models clearly showed the importance of the nonlinear dispersion terms in the governing equations.

1.2.4 Wave-current interaction models

Booij (1981), using a Lagrangian approach, developed a version of the mild slope equation including the influence of strong ambient current. Kirby (1984) presented the corrected form of this mild slope model, the parabolic approximation of this form was applied by Kirby

and Dalrymple (1983b) and results for waves interacting with a current jet were obtained. Their equation is

$$(C_g + U)A_x + VA_y + i(\bar{k} - k)(C_g + U)A + \frac{\sigma}{2} \left\{ \left(\frac{C_g + U}{\sigma} \right)_x + \left(\frac{V}{\sigma} \right)_y \right\} A$$
$$- \frac{i}{2\sigma} \left((p - V^2)A_y \right)_y + i\sigma \frac{k^2}{2} D|A|^2 A = 0$$
(9)

where $\sigma = \omega - kU$ is the principle component of the intrinsic frequency, $p = CC_g$ and \bar{k} = reference wave number, taken as the average wave number along the y-axis, and U is the mean current velocity in the x-coordinate direction and V is in the y-direction. The nonlinear term includes D, which is

$$D = \frac{(\cosh 4kh + 8 - 2\tanh^2 kh)}{8\sinh^4 kh}$$

Kirby (1986a) rederived the above equation for a wide angle parabolic approximation, which allows the study of waves with larger angles of wave incidence with respect to the x-axis. This more accurate equation was used as the basis for earlier versions of $\mathbf{REF}/\mathbf{DIF}$ 1. The equation has been extended to include the more accurate minimax approximation (Kirby, 1986b): however the present version of $\mathbf{REF}/\mathbf{DIF}$ S uses the previous model coefficients. The revised governing equation for the refraction, diffraction and shoaling of the discrete wave component n is given by

$$(C_{gn} + U)(A_n)_x - 2\Delta_1 V(A_n)_y + i(\bar{k}_n - a_0 k_n)(C_{gn} + U)A_n$$

$$+ \left\{ \frac{\sigma_n}{2} \left(\frac{C_{gn} + U}{\sigma_n} \right)_x - \Delta_1 \sigma_n \left(\frac{V}{\sigma_n} \right)_y \right\} A_n + i\Delta'_n \left[\left((CC_g)_n - V^2 \right) \left(\frac{A_n}{\sigma_n} \right)_y \right]_y$$

$$- i\Delta_1 \left\{ \left[UV \left(\frac{A_n}{\sigma_n} \right)_y \right]_x + \left[UV \left(\frac{A_n}{\sigma_n} \right)_x \right]_y \right\} + \frac{w_n}{2} A_n + \alpha A_n$$

$$+ \frac{-b_1}{k_n} \left\{ \left[\left((CC_g)_n - V^2 \right) \left(\frac{A_n}{\sigma_n} \right)_y \right]_{yx} + 2i \left(\sigma_n V \left(\frac{A_n}{\sigma_n} \right)_y \right)_x \right\}$$

$$+ b_1 \beta_n \left\{ 2i\omega_n U \left(\frac{A_n}{\sigma_n} \right)_x + 2i\sigma_n V \left(\frac{A_n}{\sigma_n} \right)_y - 2UV \left(\frac{A_n}{\sigma_n} \right)_{xy} \right\}$$

$$+ \left[\left((CC_g)_n - V^2 \right) \left(\frac{A_n}{\sigma_n} \right)_y \right]_y \right\} - \frac{i}{k_n} b_1 \left\{ (\omega_n V)_y + 3(\omega_n U)_x \right\} \left(\frac{A_n}{\sigma_n} \right)_x$$

$$- \Delta_2 \left\{ \omega_n U \left(\frac{A_n}{\sigma_n} \right)_x + \frac{1}{2} \omega_n U_x \left(\frac{A_n}{\sigma_n} \right) \right\} + ik\omega_n U(a_0 - 1) \left(\frac{A_n}{\sigma_n} \right) = 0$$

$$(10)$$

where

$$\beta_n = \frac{(k_n)_x}{k_n^2} + \frac{(k_n ((CC_g)_n - U^2))_x}{2k_n^2 ((CC_g)_n - U^2)}$$

$$\Delta_1 = a_1 - b_1$$

$$\Delta_{2} = 1 + 2a_{1} - 2b_{1}$$

$$\Delta'_{n} = a_{1} - b_{1} \frac{\overline{k}_{n}}{k_{n}}$$
(11)

The model coefficients a_0 , a_1 and b_1 depend on the aperture width chosen to specify the minimax approximation; see Kirby (1986b). The combination

$$a_0 = 1$$
 $a_1 = -0.5$
 $b_1 = 0$ (12)

recovers Radder's approximation, while the choices

$$a_0 = 1$$
 $a_1 = -0.75$
 $b_1 = -0.25$ (13)

recover the approximation of Booij (1981) which is still used in Version 1.0 of **REF/DIF** S. The values of a_0 , a_1 and b_1 used for the minimax approximation depend on the range of angles to be considered; Kirby (1986b) found that the values for a maximum angular range of 70° gave reasonable results over the range of angles typically used. The corresponding coefficient values are

$$a_0 = 0.994733$$
 $a_1 = -0.890065$
 $b_1 = -0.451641$ (14)

These will be incorporated in future versions of REF/DIF S after further testing.

Equation (10) is the model equation used in **REF/DIF** S to calculate individual wave components.

In the previous two equations, the dispersion relationship relating the angular frequency of the wave, the depth and the wave number is changed to reflect the Doppler shift due to currents. The new form of Equation (6) is

$$\sigma_n^2 = (\omega_n - k_n U)^2 = g k_n \tanh k_n h \tag{15}$$

where the absolute frequency, ω_n , is related to the intrinsic frequency, σ_n , by

$$\omega_n = \sigma_n + k_n U \tag{16}$$

where the assumption that the wave is primarily travelling in the x-direction has been used.

1.3 Assumptions

The REF/DIF S model, in parabolic form, has a number of assumptions inherent in it and it is necessary to discuss these directly. These assumptions are:

- 1. Mild bottom slope. The mathematical derivation of the model equations assumes that the variations in the bottom occur over distances which are long in comparison to a wave length. For the linear model, Booij (1983) performed a comparison between an accurate numerical model and the mild slope model for waves shoaling on a beach. He found that for bottom slopes up to 1:3 the mild slope model was accurate and for steeper slopes it still predicted the trends of wave height changes and reflection coefficients correctly.
- 2. Weak nonlinearity. Strictly, the model is based on a Stokes perturbation expansion and is therefore restricted to applications where Stokes waves are valid. A measure of the nonlinearity is the Ursell parameter which is given as

$$U = HL^2/h^3 \tag{17}$$

When this parameter exceeds about 40, the Stokes solution is no longer valid. In order to provide a model which is valid in much shallower water, a heuristic dispersion relationship developed by Hedges (1976) was provided as an option in the previous model REF/DIF 1. This relationship between the frequency and the water depth is

$$\sigma^2 = gk \tanh\left(kh(1+|A|/h)\right) \tag{18}$$

In shallow water, this equation matches that of a solitary wave, while in deep water it asymptotically approaches the linear wave result, neglecting real amplitude dispersive effects. For this reason, a model with a dispersion relationship which is a smooth patch between the Hedges form (valid in shallow water) and the Stokes relationship (valid in deep water) is used. This hybrid model is described in Kirby and Dalrymple (1986b).

A modification has been made to this description as nonlinear effects become increasingly important in situations where strong wave breaking occurs. The significant wave height at each row is therefore incorporated into the expression so that the dispersion relation for a wave component n is

$$\sigma_n^2 = gk_n \tanh\left(k_n h(1 + H_{1/3}/2h)\right)$$
 (19)

A patch between this modified form (valid in shallow water) and the Stokes form (valid in deep water) is again used to obtain a smooth dispersion relationship.

There are, as a result of the different dispersion relationships possible, three options in **REF/DIF** S: (1), a linear model, (2), a Stokes-to-Hedges nonlinear model, and (3), a Stokes model, which incorporates the third order self-interaction of each component in the wave train. Of these options, the second will cover a broader range of water depths and wave heights than the others and is also modified to correct results for breaking cases.

3. The wave direction is confined to a sector $\pm 45^{\circ}$ to the principal assumed wave direction, due to the use of model coefficients recovering the approximation by Booij (1981).

1.4 Energy Dissipation

1.4.1 General form

Energy dissipation in the model occurs in a number of ways depending on the situation being modelled. An energy loss term, due to Booij (1981) and expanded by Dalrymple et al. (1984a), permits the model to treat bottom frictional losses due to rough, porous or viscous bottoms and surface films. The linear form of the mild slope equation with dissipation is

 $\frac{\partial A_n}{\partial x} = \frac{i}{2k_n} \frac{\partial^2 A_n}{\partial y^2} - \frac{w_n}{2C_{gn}} A_n \tag{20}$

where the dissipation factor, w_n , is given by a number of different forms depending on the nature of the energy dissipation. The factor w_n is the energy dissipation divided by the energy for a wave component n and has the units of time⁻¹.

1.4.2 Laminar surface and bottom boundary layers

At the water surface and at the bottom, boundary layers occur due to the action of viscosity. For a contaminated surface, resulting from surface films (natural or otherwise), a significant amount of damping occurs, which is dependent on the value of the fluid viscosity. From Phillips (1966), the surface film damping is

$$w_n = \frac{\sigma_n k_n \sqrt{(\nu/2\sigma_n)}(1-i)}{\tanh k_n h} \tag{21}$$

where ν is the kinematic viscosity. The term under the square root sign is related to the thickness of the boundary layer, which is generally small. At the bottom, the boundary layer damping is

 $w_n = \frac{2\sigma_n k_n \sqrt{(\nu/2\sigma)(1-i)}}{\sinh 2k_n h} \tag{22}$

By setting the input switch, iff(3)=1 in the REF/DIF S model, surface and bottom damping are computed at all locations in the model.

1.4.3 Turbulent bottom boundary layer

In the field, the likely wave conditions are such that the bottom boundary layer is turbulent. In this case, an alternative specification of the energy dissipation must be provided. Utilizing a Darcy-Weisbach friction factor, f, the dissipation term can be shown to be

$$w_n = \frac{2\sigma_n k_n f |A_n|}{3\pi \sinh 2k_n h \sinh k_n h} \tag{23}$$

See Dean and Dalrymple (1984). In order to implement this damping term, the value of f = 0.01 was assumed. In **REF/DIF** S, if the input data switch *iff*(1)=1, then turbulent damping is computed at all locations.

1.4.4 Porous sand

Most sea bottoms are porous and the waves induce a flow into the bed. This results in wave damping due to the Darcy flow in the sand. For beds characterized by a given coefficient of permeability, C_p , the damping can be shown to be

$$w_n = \frac{gk_nC_p}{\cosh^2 k_n h} \tag{24}$$

The coefficient of permeability, C_p , has the units of (m^2) and is of order 4.5 x 10^{-11} m². Liu and Dalrymple (1984) show, for very permeable sands, that the damping is inversely related to C_p and a different w term must be used. However, this case is not likely to occur in nature. Porous bottom damping is computed in **REF/DIF** S when iff(2) = 1.

1.5 The Breaking Model

In contrast to the previous versions of REF/DIF 1 individual parabolic models for separate wave components are stacked up in REF/DIF S making it possible to compute statistical properties after each spatial step in the parabolic scheme. Thus the significant and root-mean-square (rms) wave heights at each point are computed after each forward step in space using the computed complex wave amplitudes.

A breaking model by Thornton and Guza (1983) is used. They showed that the energy dissipation can be expressed as

$$\frac{\partial EC_{gn}}{\partial x} = -\epsilon_b \tag{25}$$

where the energy E and the bore dissipation ϵ_b can be expressed as

$$E = \frac{1}{8}\rho g H_{rms}^2 \tag{26}$$

$$\epsilon_b = \frac{3\sqrt{\pi}}{16} \frac{\rho g \bar{f} B^3}{\gamma^4 h^5} H_{rms}^7 \tag{27}$$

Here, \bar{f} is the peak frequency of the frequency spectrum, h is the local water depth, B and γ are constants. In the testing of the model B=1 and $\gamma=0.6$ were used. The rms wave height is related to the significant wave height by

$$H_{rms} = \frac{1}{\sqrt{2}} H_{1/3} \tag{28}$$

The energy dissipation is built directly into the algorithm using an additional breaking term. No criterion for turning breaking on or off has to be stated as the breaking term is always in effect. Thus, the linear form of the mild slope equation with dissipation and breaking terms is now

$$\frac{\partial A_n}{\partial x} = \frac{i}{2k} \frac{\partial^2 A_n}{\partial y^2} - \frac{w_n}{2C_a} A_n - \alpha A_n \tag{29}$$

The coefficient α is a function of the above mentioned bore dissipation ϵ_b and is very small in magnitude when breaking does not occur but takes on significant values once breaking starts to occur, and thus energy starts to be dissipated out of the system.

$$\alpha = \frac{4\epsilon_b}{\rho g H_{rms}^2} = \frac{3\sqrt{\pi}}{4} \frac{\bar{f}B^3}{\gamma^4 h^5} H_{rms}^5$$
 (30)

No modifications have been made to Thornton and Guza's breaking model to account for directional effects. Only the change in the energy flux in the x-direction is considered and the energy flux in the y-direction is assumed constant. Also, no modifications have been made to account for the possibility that the breaking waves stop breaking and reform when they encounter an increase in water depth. These assumptions will require further future study.

The results of experiments performed by Vincent and Briggs (1989) show that in cases with strong wave breaking over a shoal nonlinear effects prevent the development of a focus due to the increase in the speed of the waves breaking over the crest. An attempt to account for this effect has been made by modifying the approximate nonlinear formulation by Kirby and Dalrymple (1986b) using the statistical wave height in shallow water. This formulation (Eqn. 19) has been discussed in Section 1.3.

1.6 Radiation Stresses

Radiation stresses for waves propagating at an angle θ to the x-direction are computed after each forward step in x. For a monochromatic wave the radiation stress in the y-direction due to the excess momentum flux in the x-direction is given by

$$S_{xy} = E \frac{C_g}{C} \sin \theta \cos \theta \tag{31}$$

where the wave energy $E = \frac{1}{2}\rho g |A|^2$.

For the spectral case a summation over all wave components is performed, thus S_{xy} is given by

$$S_{xy}(x,y) = \frac{1}{4}\rho g \sum_{n=1}^{N} \left(\frac{C_{gn}}{C_n}\right)(x,y) |A(x,y)_n|^2 \sin 2\theta(x,y)_n$$
 (32)

where $\theta(x,y)_n$ is the wave direction of the wave component n at the location (x,y) and is computed using the same algorithm as in REF/DIF 1.

Similarly the radiation stress in the x-direction due to the momentum flux in the x-direction is given by

$$S_{xx}(x,y) = \frac{1}{2}\rho g \sum_{n=1}^{N} |A(x,y)_n|^2 \left\{ \left(\frac{C_{gn}}{C_n} \right) (x,y) (1 + \cos^2 \theta(x,y)_n) - \frac{1}{2} \right\}$$
(33)

Finally the radiation stress in the y-direction due to the momentum flux in the y-direction is given by

$$S_{yy}(x,y) = \frac{1}{2}\rho g \sum_{n=1}^{N} |A(x,y)_n|^2 \left\{ \left(\frac{C_{gn}}{C_n} \right) (x,y) (1 + \sin^2 \theta(x,y)_n) - \frac{1}{2} \right\}$$
(34)

The values of these three components are written into a seperate output file *outrad.dat* after each spatial step.

1.7 Numerical Development

1.7.1 Crank-Nicolson Technique

The parabolic model is conveniently solved in finite difference form. In order to accomplish this, the study area bathymetry must be input as a grid with the (x, y) directions, divided into rectangles of Δx and Δy sizes. The complex amplitude A(x, y) will then be sought at each grid location and therefore, we can keep track of A by denoting its location, not by (x, y), but by (i, j) where $x = (i - 1)\Delta x$ and $y = (j - 1)\Delta y$. Now we have to determine the values of $A_{n_{i,j}}$ which satisfy Equation (10) for all wave components n from 1 to N, for all i between 1 and mr and for all j between 1 and nr. The procedure involves expressing all the derivatives in the (x, y) directions in terms of the complex amplitude at various grid points. For example, the forward difference representation of

$$\frac{\partial A_n}{\partial x} = \frac{A_{n_{i+1,j}} - A_{n_{i,j}}}{\Delta x} \text{ at location } i, j, \text{ for component } n$$
 (35)

If a forward difference is used for the x-direction and a central difference representation centered at i is used for the second derivatives in the lateral direction for all the derivatives in Equation (10), then an explicit finite difference equation results for $A_{n_{i+1,j}}$. This equation

can be solved directly for all the $A_{n_{i+1,j}}$, j=1,2,3...nr, for a given i and n, provided appropriate lateral boundary conditions are prescribed. This explicit representation is not as accurate as an implicit scheme and therefore an implicit Crank-Nicolson procedure is used for the amplitude calculations. For each wave component n and a given i row, the Crank-Nicolson scheme can be written

$$a_{n_i} A_{n_{i+1,j+1}} + b_{n_i} A_{n_{i+1,j}} + c_{n_i} A_{n_{i+1,j-1}} = B_{n_{i,j}}$$
(36)

The right hand side can be expressed in terms of known quantities:

$$B_{n_{i,j}} = d_{n_j} A_{n_{i,j+1}} + e_{n_j} A_{n_{i,j}} + f_{n_j} A_{n_{i,j-1}}$$
(37)

The amplitudes at (i + 1) are unknown, whereas the amplitudes at i are known from the previous calculations or boundary conditions. The coefficients a_{n_j} , b_{n_j} , c_{n_j} , d_{n_j} , e_{n_j} and f_{n_j} are complex valued, complicated functions that include nonlinear terms.

Writing the above equation for all j and for fixed i results in n tridiagonal matrix equations. These n equations can be combined resulting in a three-dimensional coefficient matrix, a two-dimensional unknown matrix and a two-dimensional known right hand side. A schematic representation of this equation is given in Figure 2. The solution for the amplitudes A_n for all j is performed at fixed i for all wave components. The tridiagonal matrix solution procedure by Carnahan et al. (1969), which is modified to handle complex variables, is used for the solution. Due to the nonlinearity of the finite difference equation, nonlinear terms are approximated on a first pass by using the $A_{n_{i,j}}$ values. Once the $A_{n_{i+1,j}}$ terms are computed, the equation is solved again for $A_{n_{i+1,j}}$ using now the just-calculated values in the nonlinear terms. This two-pass iterative method insures that the nonlinearities in the model are treated accurately (Kirby and Dalrymple, 1983a).

Furthermore, the model is set up to take advantage of a vector processor. The solution of the stacked tridiagonal matrix equations for all wave components is obtained using vectorization in n. This vectorization involves computing the amplitudes $A_{n_{i+1,j}}$ simultaneously for all n and at any j. This method greatly reduces the amount of computational time needed to solve the model equations.

The solution proceeds by moving one grid row in the x-direction (incrementing i by one) and, using the two-pass implicit-implicit technique, determining the complex amplitude $A_{n_{i+1,j}}$ for all wave components and for all the values of j on this row. Marching in the wave direction these calculations are repeated until all the A_n are known for all i and j. While it appears that the Crank-Nicolson procedure could be time consuming, given that there is a matrix inversion for each grid row, the coefficient matrix size is only 3 by n and the matrix inversion procedure is, in fact, very fast. The procedure is economical on storage as only the values for the rows i and i+1 are necessary at each calculation.

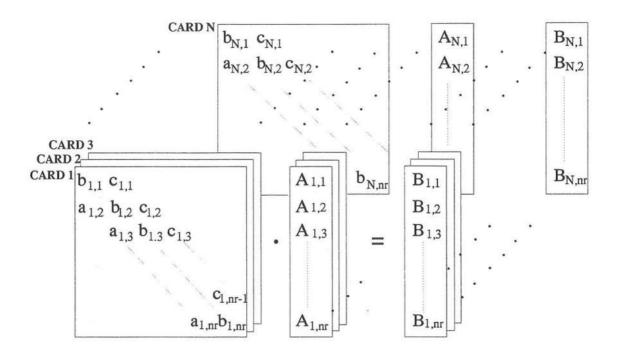


Figure 2: Vectorized Tridiagonal Matrix Equation

The amount of computations required is also greatly reduced by taking advantage of the physical properties of the domain. It is realized that the coefficients a_{n_j} , b_{n_j} , c_{n_j} , d_{n_j} , e_{n_j} and f_{n_j} include terms that are only related to the frequency of a wave component (e.g., the wavenumber k and wave celerity C). These terms only have to be computed for wave components with different frequencies as they will be the same for two waves with the same frequency and different directions.

1.7.2 Initial and Lateral Boundary Conditions

The initial condition is vital for the parabolic model. The furthest seaward grid row, corresponding to i=1, is taken as constant depth and the discretized incident wave spectrum is prescribed here. The individual wave components are then propagated over the bathymetry by the model.

As in the solution of any differential equation in a domain, the lateral boundary conditions are important. There are several ways to treat the boundaries; however, none of the presently existing boundary conditions result in the total transmission of scattered waves. Therefore, for the REF/DIF S model, a totally reflecting condition is generally used for each side (j = 1 and n). This requires that the specification of the model grid be done with care, as the reflection of the incident wave from the lateral boundaries can propagate into

the region of interest rapidly and result in erroneous results.

In general, the width of the model should be such that no reflection occurs until far downwave of the region of interest. As a precaution, a graphical representation of the computed wave field should be examined to determine where the reflection from the boundaries is important. By using the switch, *ibc*, partially transmitting boundaries can be used (Kirby, 1986c). In general, this boundary condition will result in less reflection in the model domain; however, since some reflection will occur, it is recommended that runs be carried out with the reflecting boundary conditions in order to assess the regions potentially affected by reflection from the model boundaries.

1.7.3 Subgrids

In order to reduce the amount of data input and yet provide the user the ability to prescribe the fine scale bathymetry in areas of interest, REF/DIF S utilizes a coarse scale user-specified reference grid and a fine scale subgrid, which can have many times the resolution of the reference grid. The principal purpose of the subgrid is to provide enough computational points to the numerical model to preserve accuracy. The user specifies the number of subgrid divisions in the y-direction with the parameter nd. If nd=1, then the subgrid spacing in the y-direction is the same as the reference grid. If nd=2, then the model uses twice as many computational points in the y-direction as there are in the reference grid. In the propagation direction, x, the model will automatically determine the subgrid spacing if ispace has been set to unity. Otherwise, the user provides the subgrid spacing using the input, mr, which permits variable spacing in the x-direction. For subgrids, the input flag, isp, must be set to one and an array, isd, must be specified.

1.7.4 Damping

When the large angle parabolic approximation is used as a basis for the computation of wave fields around islands, the presence of wave breaking, and resulting sharp lateral variations in wave height, leads to the generation of high-wavenumber spectral components in the computed complex amplitude A in the lateral (y) direction. Kirby (1986a) has shown that these components have propagation velocities which can become large in an unbounded fashion; as a result, they can propagate across the grid, filling the computational domain with high-wavenumber noise.

The high frequency noise generated by the breaking process is controlled by an algorithm which has recently been developed by Kirby (1993). Reference should be made to that document for a description of the algorithm. The damping is built into the computational algorithm and is turned on at all times. The damping coefficient *cdamp* is specified in

the code for **REF/DIF** S at the beginning of the subroutine *fdcalc* (after the statement functions). The effects of several values of damping coefficients on the wave fields are documented in Kirby (1993).

1.7.5 Breaking Model

The Crank-Nicolson scheme is used for most terms in the equation and involves representing x-derivatives around $(i+\frac{1}{2},j)$. If the breaking term, however, is implemented using Crank-Nicolson, it is observed to be stable everywhere in the domain except in the vicinity of a "thin film", where undesired effects occur. Therefore, a fully implicit approach is used in the vicinity of a "thin film". To achieve this a parameter β is introduced, such that the simplest form of the linear, constant depth equation

$$\frac{\partial A_n}{\partial x} = -\alpha A_n \tag{38}$$

can be written in finite difference form as

$$\frac{A_{n_{i+1,j}} - A_{n_{i,j}}}{\Delta x} = -\alpha \left[\beta A_{n_{i+1,j}} + (1 - \beta) A_{n_{i,j}} \right]$$
 (39)

The parameter β is defined as

$$\beta = \frac{1}{2} + \frac{1}{2} \left(\frac{h_{tf}}{h} \right)^3 \tag{40}$$

where h_{tf} is the water depth on the "thin film" and h is the local water depth. The parameter β reduces to 0.5 everywhere, giving a centered Crank-Nicolson scheme, but takes on unity on the "thin film", yielding a fully implicit step.

2 USER'S MANUAL

This chapter provides the operating manual for the program REF/DIF S. Chapter 3 provides sample documentation and calculations for example problems.

2.1 REF/DIF S Program History

REF/DIF S is essentially an enhancement to REF/DIF 1. REF/DIF S allows for the simultaneous computation of wave components with different frequencies and directions in order to simulate a random sea. Such a simulation would have to be made using sequential computations in REF/DIF 1. As wave amplitudes for all components are computed simultaneously at any grid location, statistical analysis can be performed. REF/DIF S also incorporates a statistical wave breaking model.

REF/DIF S should be regarded as a statistical model. It should be noted that some options valid for REF/DIF 1 are not included in the structure of REF/DIF S.

- 1. The desired two-dimensional spectrum is input as discrete wave components having frequencies and directions associated with them. Because of the large amount of wave components needed to describe a spectrum properly, it is fairly unlikely that a user would desire to input the complex amplitude of each wave on the offshore boundary. Therefore, this option is turned off. This corresponds to setting the switch *input* in the data file to unity.
- 2. Another option existing in REF/DIF 1 gives the user the opportunity to save the last row of computed amplitudes. Because of the large number of wave components present this option would result in a file that would allocate a large amount of disk space. Anticipating that this option would not be desired, it is turned off. This is achieved by setting *ioutput* to unity.
- 3. REF/DIF S expects an initial wave field in form of discrete wave components (iwave=1). These wave components may represent any desired two-dimensional spectrum. Therefore, there is no need for an option, iwave=2, which in REF/DIF 1 corresponds to a directional spreading of the input wave condition. So iwave has to be set to unity.
 - Even if any mistakes are made during input of these three switches, the program will automatically correct those and set the switches to the mentioned values.
- 4. The value of the tidal offset has to be a constant with respect to all frequencies. This was not a requirement in REF/DIF 1, where separate components are treated as completely independent simulations.

5. The number of waves per frequency *nwavs* is a constant, too. Thus there can only be a uniform number of directional wave components for each frequency component.

REF/DIF S needs the following additional information as input:

- 1. As a discretized spectrum is input into the model, the smooth spectrum cannot be reconstructed. Therefore, the peak frequency of the original spectrum that was discretized by a preprocessor has to be input by the user.
- 2. The output from REF/DIF S is given in terms of the statistical wave height. An option is built into the model in case the user might want to analyze any specific wave component. If turned on this option enables amplitude information about that specific component to be written into a separate output file specamp.dat.

Because of the possibility that the user wants to use an input file constructed for REF/DIF 1, these additional lines are not necessary in cases where only one wave component is specified. However, if more than one wave component is specified in the input file, REF/DIF S will require these additional lines to be present in the input file.

2.1.1 Changes Appearing in Version 1.1

Version 1.1 includes a revised input section for REF/DIF S which ensures that existing indat.dat files constructed for REF/DIF 1 to simulate monochromatic cases can be applied to REF/DIF S with no modification. In addition, this version includes modifications to the numerical implementation of the wave breaking algorithm to ensure proper interaction with the "thin film". Further minor revisions include modifications in the input of the preprocessor SPECGEN and enhancements to the documentation.

2.2 Overview of Operating Manual

This section provides a description of the program structure (section 2.3), followed by some notes on problems which are likely to be encountered during the installation and use of the program on different computer systems (section 2.4). Section 2.5 presents the two levels of grid information used by the program. Section 2.6 gives information about preprocessing spectral input. The input data file structure is then discussed. The program reads data in two essentially separate groups. The first group of data establishes the size of the model grid and gives the wave conditions to be studied; this group is discussed in section 2.7. The second group of data gives the reference grid data values and defines any user-specified subgrids; this is discussed in section 2.8. The structure of the program output is discussed in section 2.9. A listing of the program is included in Appendix A. A listing of a plotting

program used to interpret the model output is given in Appendix B. An *fweb* documentation of a preprocessor used to generate input data files is given in Appendix C. A list of frequently asked questions is given in Appendix D.

2.3 Program Outline and Flow Chart

The model REF/DIF S is organized in one main program refdifs and eleven subroutines. The program does not contain calls to any external package, and should be free standing on any system. The program is written in a subset of FORTRAN 77 and should be a legal code for any compiler which accepts the full FORTRAN 77 standard language. Exceptions might be **open** statements, which might have to be altered due to macine and site specific conditions.

REF/DIF S is structured in two levels; a main level, which reads in and checks input data and then starts the operation of the wave model level, and the model level itself, which performs the actual finite difference calculations. The flow charts for the two levels are given in Figures 3 and 4. A short description of each routine in the model follows.

- refdifs: Main program controls the calls to inref and inwave to read in data, and to model, which performs the actual calculations. No calculations are performed by this routine.
- 2. inref: Called by refdifs. inref reads in data controlling reference grid dimensions and the grid interpolation scheme from logical device number iun(5), and reads in the reference grid values of depth dr, x-direction velocity ur and y-direction velocity vr from logical device number iun(1). Some data checking is performed. If data is read in in English units, inref converts it to MKS units using the dconv factor. Output files are initialized. A description of the data files may be found in sections 2.5-2.8 of this manual. At the end of the subroutine, control is returned to refdifs.
- 3. inwave: Called by refdifs. inwave reads in data specifying the initial wave field along the first row of reference grid points. Data is read from logical device number iun(5). Conversion to MKS units is performed for data read in in English units. Control is returned to refdifs.
- 4. model: Called from refdifs. model controls execution of the computational part of the program. model performs the following series of operations:
 - (a) Initializes program by calculating the incident wave field on the first grid row for all wave components.

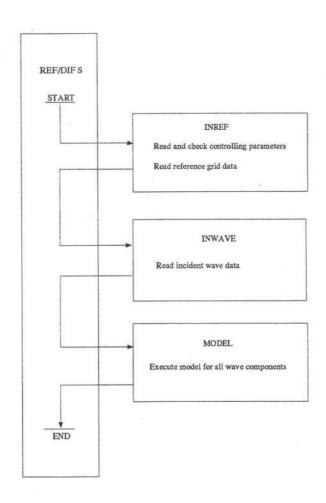


Figure 3: REF/DIF S: refdifs program level

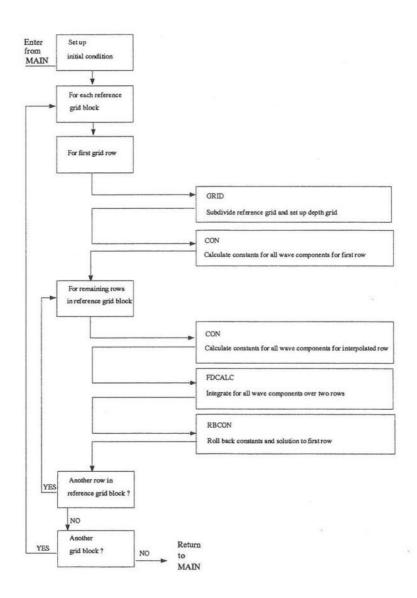


Figure 4: REF/DIF S: model subroutine level

- (b) Then, for each grid block in the reference grid:
 - Calls grid to perform the grid interpolation specified in the input data.
 - Calls con to calculate constants on the first row for all frequency components in the interpolated grid.
 - For each row in the interpolated grid:
 - Calls con to calculate constants on row for all frequency components.
 - Calls fdcalc to perform the numerical integration of the parabolic equation on row for all wave components.
 - Calls *rbcon* to roll back solution to previous row for all wave components.

The model execution is then complete. Control is returned to refdifs.

- 5. grid: Called by model. grid performs the required interpolation over a single grid block of the reference grid as specified in the input data. The interpolation is performed as described in Section 2.5. grid checks to see whether a user-specified subgrid feature should be read in, and reads it in from logical unit number iun(2) if needed. The interpolated depth grid is then corrected for tidal offset, and checked for surface-piercing features. These features are modified using the "thin-film" approach; see Kirby and Dalrymple (1986a). Control is returned to model.
- con: Called by model. con calculates various constants for all frequency components for a row in the reference grid created by grid. Control is returned to model.
- 7. fdcalc: Called from model. fdcalc performs the integration of the governing parabolic equation over the grid defined in grid. The coefficients of the finite difference form of the parabolic equation are developed according to the Crank-Nicolson method. A complete description of the equations and the treatment of nonlinearities may be found in Kirby (1986a) and Kirby and Dalrymple (1986b).
 - The parabolic equation is integrated for one row in the subdivided ir grid block and for all wave components. Control is then returned to model.
- 8. rbcon: Called by model. rbcon overwrites previous values of constants and solution with computed values and re-initializes arrays containing these computed values. So most recently calculated values are rolled back to the previous row for all wave components.
- 9. vwnum: Called by grid and con. vwnum performs a Newton-Raphson solution of the linear wave-current dispersion relation to obtain values of the wavenumber k for a grid row for all frequency components.

- 10. vtrida: Vectorizable utility routine which is called by fdcalc to perform the double sweep elimination to solve the implicit set of equations for all wave components.
- 11. diss: Called by con. diss calculates frictional dissipation coefficients based on values of the switches read in by inref for all wave components.

2.4 Special Installation Instructions

Several features of the program REF/DIF S may require some modification or customizing during program installation on various systems. REF/DIF S is written using the features of FORTRAN 77. The code is arranged to take advantage of automatic vectorization in the CRAY CF77 compiler; however, no explicit vector code is used.

If the program is to be used on machines with no upward limit placed on the size of the compiled, executable code, only two variables have to be checked during the initial installation of the program. These are the logical device numbers for the controlling input date file and the device number for line printer output. These numbers may vary from system to system. They appear in the program as:

- iun(4): logical device number for line printer output. (initialized in the main program refdifs).
- iun(5): logical device number for controlling input data file. (initialized near the top of subroutine inref).

The supplied version of REF/DIF S has those values set to iun(4)="*" and iun(5)=5. If a device number is to be given for iun(4), the choice of 6 for output should be fairly standard (one exception is the DEC-10, which uses the device number 5 for both input from file or screen and output to the line printer). The version of the program currently used on SUN workstations has the default device indicator "*" substituted for iun(4) everywhere. Output can be directed to a data file using the directed output switch in either UNIX or DOS.

Many systems require that access to disk files be initialized and terminated by the use of **open** and **close** statements in the program code. Since the parameter list of the **open** statement varies from system to system, the user should take care that the **open** statements are compatible with the system software being used. The **open** statements appear near the top of refdifs, inref and model. The corresponding **close** statements appear near the end of refdifs. The present model is internally configured to run on a SUN system.

When writing **open** statements, a choice of a fixed record length of 132 (line printer width) will safely bracket all the **format** specifications in the program.

REF/DIF S has at this time been run on SUN Sparcstation 2's, IBM RS 6000 and CRAY YMP machines, all using versions of UNIX operating systems.

2.5 Computational Grids and Grid Interpolation

The reference grid terminology is defined in Figure 5. The grid consists of a mesh of points with dimensions $mr \times nr$ in x and y. The values of mr and nr must be less than or equal to the parameters ixr and iyr whose values are set in the **parameter** statements. Reference grid data of depth dr and ambient current components ur and vr are defined at the grid points. The program assumes that the x, y coordinate system is established with the origin at grid point (ir,jr)=(1,1). In this manual, we make the distinction between the terms "grid row" ir, which is the row of points jr=1, nr at location ir, and "grid block" ir, which is the physical space between grid rows ir and ir+1.

The reference points are separated by spacings dxr and dyr which are uniform in the xand y-directions. The spacings dxr and dyr may have arbitrary, independent values.

Values of dr, ur and vr constitute the "reference grid data". Section 2.8 describes the required input data file for these quantities.

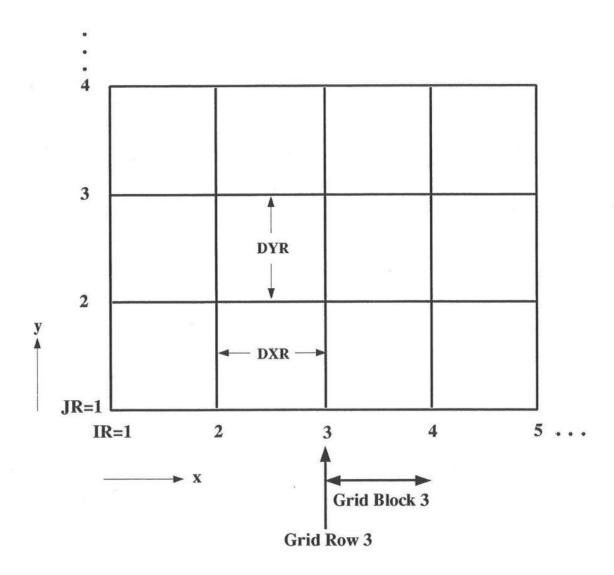


Figure 5: Reference grid notation

The computational grid for any particular application should be chosen with care. Since $\mathbf{REF}/\mathbf{DIF}$ S tries to use at least 5 points per wavelength of the shortest wave being modelled, the length of the computational domain in the propagation direction is restricted (with the given **parameter** statements). This range can be extended by increasing ixr and/or ix in the **parameter** statements. The width of the model domain should be chosen such that interference from the boundaries does not affect the study area. If the reflective boundary conditions are used, the extent of boundary influence is usually obvious, particularly when viewing two-dimensional plots of the significant wave height $H_{1/3}$.

2.5.1 Grid Subdivision

The only major feature of REF/DIF S which is not described in Chapter 1 is the ability to subdivide the given reference grid into a more finely subdivided computational grid. This would usually be done in cases where the reference grid spacing is too large to be used directly for calculations. In this case, the user may specify how the reference grid is to be subdivided, or the user may tell the program to attempt its own subdivisions.

The maximum reference grid dimensions have been set fairly arbitrarily and may be changed by modifying the **parameter** statements at the beginning of each routine in the program. The grid is large enough to comply with typical grids for various tidal computational models, which may be used to specify the ambient currents, and get small enough so that data values do not occupy too much internal storage. It is anticipated that the spacings dxr and dyr, if based on such a model, may be too large for an accurate integration of the parabolic model to be performed. Consequently, the model has been arranged so that the reference grid block between any two adjacent reference grid rows may be subdivided down to a finer mesh in order to provide sufficient resolution in the computational scheme. This subdivision process is performed internally in the program (with an exceptional feature to be described below), and may be controlled by the user or the program.

Remember that the computational scheme proceeds by marching in the x-direction, and, therefore, the only reference grid information required at any particular step is the data on row ir, where computation starts, and on row ir+1, where computation ends. The fine, subdivided mesh is set up on the intervening grid block. An example of a particular subdivision of a grid block is shown in Figure 6. Here, the choices of x-direction subdivision md(ir)=5 and nd=2 are illustrated, with md and nd being the number of spaces each reference grid cell is divided into, rather than the number of extra points being inserted.

Several restrictions are placed on the choice of nd and md's. First, the maximum dimensions of the subdivided grid cell is given by the parameters ix and iy. This implies that any md can be at most (ix-1) and nd can be at most (iy-1)/(nr-1). The maximum number of added spaces may be increased by increasing ix and iy in the **parameter** statements. Further, the y-subdivision specified by nd is applied uniformly along each grid row for the full extent to the reference grid. No provision is made for variable grid spacing in the y-direction. Grid spacing in the x-direction is arbitrary, so md's may differ arbitrarily for each grid block.

The user must specify the single value of nd in the input data. Two choices may be made regarding md's, however.

1. The user may let the program calculate md. The program proceeds by calculating

Reference Grid

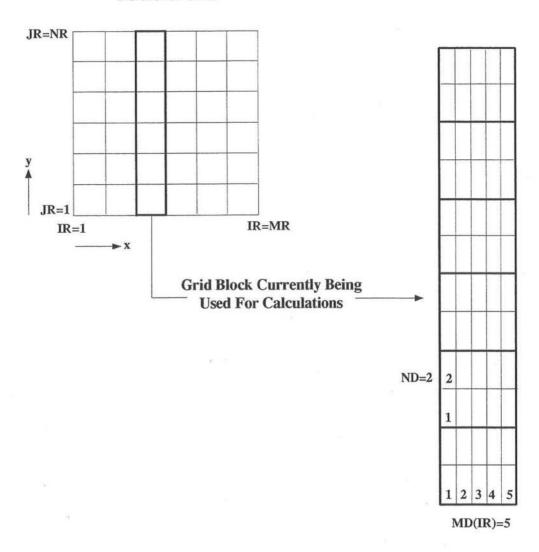


Figure 6: Sample grid subdivision

an average wavenumber along the reference grid row, and uses this to estimate the wavelength in the grid block. The program then chooses a subdivision so that at least 5 grid points per wavelength of the shortest wave being modeled will occur in the grid block. The program checks to see that this desired number of subdivisions does not exceed the maximum. If it does, the program reduces the number to the maximum and prints a warning message indicating that the grid block can not be subdivided finely enough. Computed results in this case must be regarded as being suspect. The number of subdivisions used is indicated on the output. The user chooses this option by setting the switch *ispace*=0 on input.

2. The user may specify each value of md(ir) from ir=1 to (mr-1). This is done by setting the switch ispace=1 on input, and the values of md are then read in from the input data file. Note that this choice is necessary if the user-defined subgrids discussed below are to be used, since the user will be inputting subgrids with prespecified spacings. As it is presently written, the program will only print a warning if it encounters subgrids with ispace=0 chosen; extensive garbling of input data may result.

After the subdivided grid block is established, the program uses this grid as the actual computation grid. New values of depth d and ambient current u and v are calculated at the extra grid points by fitting a twisted surface to the reference grid using linear interpolation. An example of the resulting bottom topography for a single grid cell is shown in Figure 7.

2.5.2 User-specified Subgrids

In some applications, an important topographic feature may be present at a subgrid scale within the reference grid. Examples include artificial islands, shoals, borrow pits, etc., which are superimposed on an otherwise slowly-varying topography which is represented by the sort of grid resolution appropriate to tidal models. An illustration of such a feature is shown in Figure 8, where a poorly resolved feature occupies portions of four reference grid cells. For cases such as this, the program includes the option for the user to input user-defined, subdivided grid cells in order to specify these features at the level of the computational grid.

The use of user-defined subgrids implies that the user will be choosing the grid spacing of the subgrids. Since grid spacings must be uniform across a grid block, the user should choose the option, *ispace*=1 and specify the values of md in the input file, *indat.dat*. Then, the program will look for a data array of the correct, pre-specified dimension when it reads the subgrid data from the file, subdat.dat (unit number iun(2)).

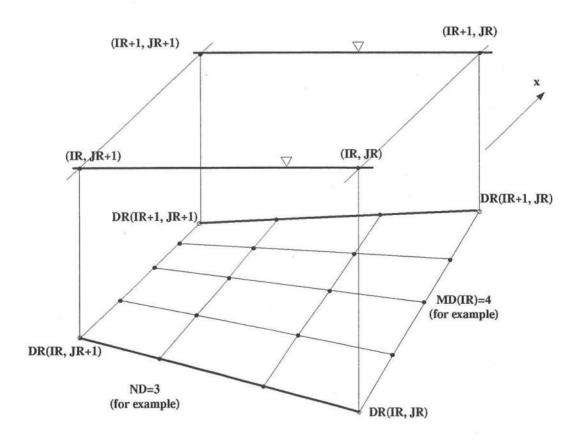


Figure 7: Interpolation of depth data

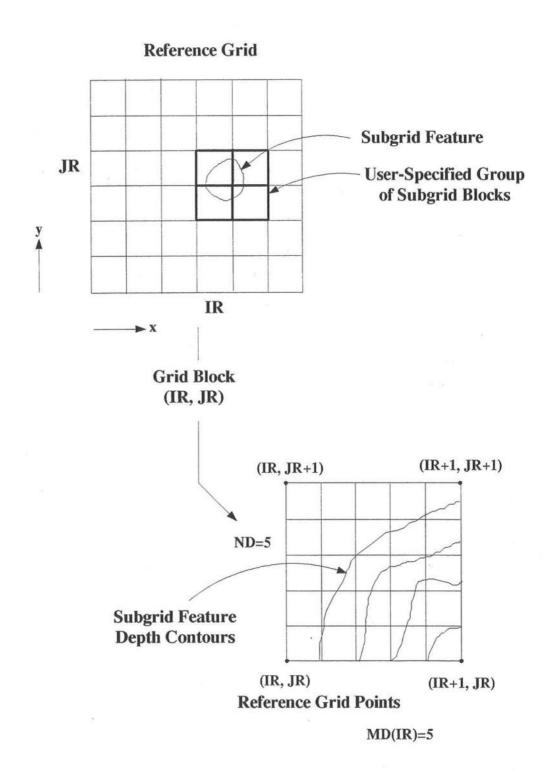


Figure 8: User-defined subgrids

Several aspects of the subgrid data should be noted. Data for depth d and current velocities u and v need to be specified at each of the subgrid points. These data are input in the same units, and with the same datum for depth as the data for the reference grid. Also, note that the subgrid includes the points on its boundaries; for example, the single subgrid shown in Figure 8 has a dimension of 6 by 6.

The data values on the outer borders of the subgrids should be setup to match with the linearly interpolated values in the region external to the user-defined subgrid. If the data do not vary linearly along the subgrid boundaries, there may be some mismatch between the subgrid boundaries and the external region. An exception to this rule occurs when two subgrids adjoin each other. Then, care should be taken that the data along the common boundary match. These common boundary point data are duplicated in the input data, since each subgrid data set includes the boundaries.

Instructions and formats for creating the subgrid data file *subdat.dat* are included in section 2.8.

2.6 Program Input: Preprocessing a Spectral Input

In many cases, the spectral density at the offshore boundary might be available as data and the user might want to use this information as input to the model. As can be seen in the next section, the program REF/DIF S requires the input of discrete wave components that describe the desired two-dimensional spectrum. In order to use spectral data at the off-shore boundary, the user has to discretize the spectrum and input the individual wave components in the manner described in the next section.

Several approaches to the task of discretizing a two-dimensional spectrum can be used. In order to clarify the procedure to be followed, the preprocessor SPECGEN is included in Appendix C. This program is specificly designed to create input files for three of the example cases in Chapter 3, which simulate experiments by Vincent and Briggs (1989). SPECGEN discretizes a TMA frequency spectrum (Hughes, 1984) used in conjunction with a wrapped normal directional distribution (Borgman, 1985). A similar procedure has to be followed for the discretization of any other design or measured spectrum.

SPECGEN divides the two-dimensional spectrum into bins of constant volume, therefore the spectrum is well resolved in the region of the peak frequency and the mean angle of incidence. As the volume of each bin is constant, the energy, and hence the amplitude, corresponding to each wave component is constant. The midpoints of each bin in frequency and in direction are chosen as the representative values for frequency and direction for that bin.

The discretization of the two-dimensional spectrum performed by SPECGEN results in

wave components of amplitude A with an associated frequency f and an angle of incidence θ with respect to the assumed propagation direction. SPECGEN outputs this information in the format of a indat.dat file. The format of this file is described in the next section.

In order to use any design or measured spectrum as input, the user has to discretize the spectrum and create discrete wave components of amplitude A with an associated frequency f and an angle of incidence θ . The method of discretization used is not relevant to **REF/DIF S**

Figure 9 shows plots of a divided TMA frequency spectrum and a divided wrapped normal directional spreading function. Discretization has been performed by SPECGEN.

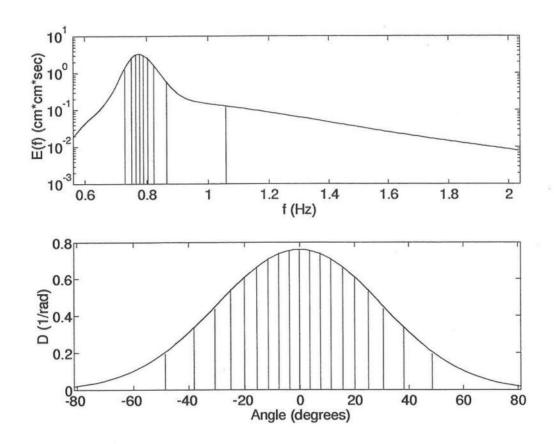


Figure 9: Divided Frequency Spectrum and Directional Spreading Function

2.7 Program Input: Model Control and Wave Data

This section discusses the structure of the input data file indat.dat read in through logical device number iun(5). This file contains all the information needed to control the operations of the program, and all of the input wave data. The file indat.dat is kept fairly standard

so that already existing data files for REF/DIF 1 can be used with no changes if a monochromatic wave is described and with only minor changes if more than one wave component is described. Data files for REF/DIF S can be used with REF/DIF 1 without any alterations. As the program structures are slightly different some switches for options existing in REF/DIF 1 are set to default values. The preprocessor SPECGEN will set the proper values of those switches. In addition, the default values are stated here and reference is made to Section 2.1 for details.

Reference grid values of depth dr and currents ur and vr are treated in the following section.

Data is read from iun(5) in both the *inref* and *inwave* subroutines. Each line of input data is in free format. The input file is set up as follows:

1. iun(1), iun(2), iun(3):

Logical device numbers for reference grid data, user-specified subdivision data, and reference grid output. Note that the values 5, 22 or 36 should not be chosen for iun(1), iun(2), iun(3) as these are default values chosen by the program for output files.

2. mr, nr:

Reference grid dimensions. Maximum values are ixr, iyr respectively.

3. iu, ntype, icur, ibc:

iu is switch for physical units; iu=1, MKS; iu=2, English. (program defaults to iu=1 if an error is made on input).

ntype is switch for nonlinearity; ntype=0, linear model; ntype=1, composite model; ntype=2, Stokes wave model.

icur is switch for input current data. icur=0, no currents input; icur=1, currents input.

ibc is the boundary condition switch. ibc=0, closed boundaries; ibc=1, open boundaries.

Both icur and ibc default to zero on input error.

4. dxr, dyr, dt:

Reference grid x-spacing and y-spacing (uniform over entire reference grid) and depth tolerance value.

5. ispace, nd:

ispace is switch controlling subdivisions; ispace=0, program attempts its own x-subdivisions. ispace=1, user specifies x-subdivisions.

nd is the number of y-direction subdivisions (needed in either case).

If *ispace*=1, the x-direction subdivisions are inserted here (one for each reference grid block):

(md(ir), ir=1, mr-1)

6. iff(1), iff(2), iff(3):

Dissipation switches; iff(1)=1; turn on turbulent boundary layer; iff(2)=1; turn on porous bottom damping; iff(3)=1; turn on laminar boundary layers. No damping if all values are zero.

7. isp:

Switch for user-specified sub-grid specifications. isp=0, no subgrids to be read; isp=1, subgrids will be read.

8. iinput, ioutput:

The switches *iinput* and *ioutput* control options that exist in REF/DIF 1, but that are not included in the structure of REF/DIF S. Nevertheless, this line of input is required for the user to be able to use this input file with REF/DIF 1, if desired. For more information, the reader is referred back to Section 2.1. For a proper run with REF/DIF S this line of the input file should be

1 1

Both iinput and ioutput default to unity on input error.

As iinput=1, the rest of indat.dat is as follows:

9. iwave, nfreqs:

iwave is the switch for wave field type. If iwave=1 the program expects discrete wave components to be input. Any other options are not incorporated into REF/DIF S. More information can be found in Section 2.1. Thus, set iwave=1. iwave defaults to unity on input error.

nfreqs is the number of frequency components to be run.

As *iwave*=1, the remainder of the file is as follows:

10. freqs(ifreq), tide:

Wave period for each frequency component and tidal offset, which is a constant for all frequencies.

11. nwavs:

Number of directional components per frequency, again a constant for all frequencies.

12. amp(ifreq, iwavs), dir(ifreq, iwavs):

Amplitude (not height) and direction (in degrees) for each component wave.

repeat items (11) and (12) nwavs times

repeat entire block from item (10) to item (12) nfreqs times

The following lines are needed in addition for the execution of **REF/DIF** S in cases where more than one wave component is desired.

13. fpeak:

Peak frequency of input frequency spectrum in Hz.

14. iopt:

iopt specifies whether computed amplitudes for any specific wave component are to be stored in a separate file specamp.dat. iopt=1 turns the option on, iopt=0 turns it off.

If iopt=0, the data file terminates here. No additional information is necessary. If iopt=1 the following information is needed.

15. ifcomp, idcomp:

This line specifies the wave component for which output will be produced. *ifcomp* specifies which frequency component is desired and *idcomp* specifies which directional component is desired.

Examples of data files are given in Chapter 3.

2.8 Program Input: Reference Grid and Subgrid Data

This section describes two files which provide arrays of data on the various grids.

Reference Grid Data File

This file (named refdat.dat) is accessed as logical device number iun(1). Its contents consist of the arrays of depth dr, x-velocity ur, and y-velocity vr at the reference grid points. This file is accessed only once per model run, and its entire contents are read in by subroutine inref. If icur=0, only the depth data dr need to be specified.

Data for this file should be written in the following format:

```
(then, if icur=1)

do 2 ir=1, mr
  write(unit,#) (ur(ir, jr), jr=1, nr)

continue

do 3 ir=1, mr
  write(unit,#) (vr(ir, jr), jr=1, nr)

continue

format(16f8.4)
```

The data may be in either MKS or English units; set the units switch iu in the iun(5) data file indat.dat accordingly.

Subgrid Data File

This file (named subdat.dat) is accessed as logical device number iun(2). If no user-defined subgrids are to be read in, this file may be omitted. The file consists of two parts; an integer array of 1's and 0's indicating which reference grid cells are to be defined by the user, and then a sequence of groups of arrays of d, u and v, one group for each subgrid.

The integer array isd is dimensioned (mr-1) by (nr-1), with one point for each spacing in the reference grid. The array contains a 0 if that cell is not to be user-defined, and a one if it is. For example, the (mr,nr)=(7,6) reference grid shown in Figure 7 has four cells which are to be read in as user-defined subgrids.

The array isd should be written in the data file first, using the format:

```
do 1 ir=1,mr-1
write(unit,#)(isd(ir,jr),jr=1,nr-1)
continue
format(15I4)
```

Now, a group of arrays of d, u and v must be entered into the data file, with one group corresponding to each value of 1 in isd. The program accesses the data file by x-row (ir) and then by y-column (jr). For the above example, the subgrid array groups should thus be stored in the order of the coordinate pairs (4,3), (4,4), (5,3), (5,4). The dimensions of each subgrid are given by m=md(ir)+1 and ns=nd+1. The borders of adjacent subgrids share the same common boundary points.

Each group of subgrid data should be written using a format similar to:

```
(then, if icur=1)
write(unit,#)((u(i,j),j=1,ns),i=1,m)
write(unit,#)((v(i,j),j=1,ns),i=1,m)
# format(16f8.4)
```

using the appropriate value of m for the grid block in question. Data may be in MKS or English units, depending on the value of iu in the input data file. The integer array isd is read by inref, and the individual subgrids are read by grid.

2.9 Program Output: Screen Output

This section discusses the output sent to the screen (or the chosen output device with logical unit number iun(4)). Screen output consists of three types: title and header information which is printed in order to identify the run and the operating parameters, run-time error messages (either warnings or terminal error messages), and the results of calculations.

Header Information

At the start of each run, a set of two title pages are printed. Page 1 identifies the program and then prints out messages identifying the parameters which set up the model run, as read in by subroutine *inref*. A sample title page 1 is shown in Figure 9.

Title page 2 gives the input wave conditions which were read in by *inwave*. Part of a sample title page 2 is given in Figure 10.

Dimensional quantities are printed on the title page in the units used for input. Title page 1 gives an indication whether quantities are in MKS or English units.

Run-time Error Messages

REF/DIF S performs some data checking and checking of calculations during a run. This checking may result in warnings or terminal errors which are beyond calculation errors which would lead to standard FORTRAN error messages. A list of possible errors and the resulting messages follow.

1. Reference grid dimensions were specified as being too large on input. mr > ixr and/or nr > iyr.

Message: dimensions for reference grid too large; stopping.

Refraction-Diffraction Model for Spectral Wave Conditions

REF/DIF"S Version 1.1

Center for Applied Coastal Research Department of Civil Engineering University of Delaware Newark, Delaware 19716

James T. Kirby and H.Tuba Ozkan, July 1992, February 1994

input section, reference grid values

reference grid dimensions mr=251 nr=275

reference grid spacings dxr= 0.1000 dyr= 0.1000

physical unit switch iu=1, input in mks units
icur=0, no current values read from input files
ibc=1, open lateral boundaries
ispace =0 chosen, program will attempt its own reference grid subdivisions

ntype = 1, stokes model matched to hedges model

y-direction subdivision according to nd= 1

switches for dissipation terms

- 0 turbulent boundary layer
- O porous bottom
- 0 laminar boundary layer

isp=0, no user defined subgrids
iinput = 1, program specifies initial row of a

Figure 10: Sample title page 1

input section, wave data values

iwave=1, discrete wave amps and directions

the model is to be run for 10 separate frequency components

Figure 11: Sample title page 2

Result: Program stops.

Error occurs in: inref

Action: Change the values of ixr and iyr in the parameter statements, such that $ixr \geq mr$ and $iyr \geq nr$.

2. User specifies a y-direction subdivision nd which will cause the number of y grid points n to exceed the maximum iy.

Message: y-direction subdivision nd="nd" too fine. maximum number of y grid points will be exceeded. execution terminating.

Result: Program stops.

Error occurs in: inref

Action: Change the values of iy in the **parameter** statements, such that $iy \ge (nr-1)nd+1$. Also, see Section 2.5.1.

3. User specifies an x-direction subdivision on one of the grid blocks ir which exceeds the maximum amount (ix-1). As a result, the dimension of the subdivided grid will be too large.

Message: x-direction subdivision md="md" too fine on grid block "ir", execution terminating.

Result: program stops

Error occurs in: inref

Action: Change the values of ix in the **parameter** statements, such that $ix \ge (md(ir)+1)$ at the grid block ir indicated in the error message. At this point the user should check if $ix \ge (md(ir)_{max}+1)$, where $md(ir)_{max}$ is the maximum number of x-direction subdivision on the grid blocks, to prevent reoccurance of this error. See Section 2.5.1.

4. A depth value occurs in the reference grid which differs from the average of its neighbors by more than the tolerance value dt specified on input. This is basically a data checking feature. Printed values are in meters.

Message: Depth "dr" (m) at reference grid location "ir,jr" differs from the average of its neighbors by more than "dt" (m). Execution continuing.

Result: None by program.

Error occurs in: inref

Action: Data in file refdat.dat should be checked for mistakes and corrected if wrong.

5. An ambient current value occurs which implies that the flow would be supercritical at the given location. This serves as both a check for anomalously large current values, and an indicator of possible subsequent computational problems.

Message: ambient current at reference grid location "ir,jr" is supercritical with froude number = "froude number", execution continuing

Result: None by program.

Error occurs in: inref

Action: Data in file refdat.dat should be checked for mistakes and corrected if wrong.

6. If the user specifies that predetermined subgrids are to be read in, while at the same time telling the program to performs its own subdivisions, the computed dimensions of the subgrid may be different than those of the subgrid included in the input. Runs requiring user-specified subgrids should choose the ispace=1 option. If an incompatible set of dimensions occurs, the program will either garble the input array or run out of data.

Message: Warning: input specifies that user will be supplying specified subgrids (isp=1), while program has been told to generate its own subgrid spacings (ispace=0). possible incompatibility in any or all subgrid blocks.

Result: None by program.

Error occurs in: inref

Action: Should restart unit with correct ispace, isp values. See Section 2.5.1.

 User specifies number of frequency components which exceeds the maximum amount icomp.

Message: Too many frequency components; stopping.

Result: Program stops.

Error occurs in: inwave

Action: Change the value of icomp in the **parameter** statements, such that $icomp \ge nfreqs$. The value of nnii in the **parameter** statements has to be changed, as well, such that $nnii \ge (icomp)(nnd)$.

8. User specifies number of directional components which exceeds the maximum amount nnd.

Message: Too many directional components; stopping.

Result: Program stops.

Error occurs in: inwave

Action: Change the value of nnd in the **parameter** statements, such that $nnd \ge nwavs$. The value of nnii in the **parameter** statements has to be changed, as well, such that $nnii \ge (icomp)(nnd)$.

9. While calculating its own subdivision spacings, the model may try to put more division in a reference grid block than is allowed by dimension ix. If this occurs, the program uses the maximum number of subdivisions allowed (ix-1), but prints a message indicating that the reference grid spacing is too large with respect to the waves being calculated. This problem may be circumvented by increasing the size of ix in parameter statements.

Message: model tried to put more spaces (md="md") than allowed in grid block "ir". Result: Program performs fixup and continues. Model resolution and accuracy may be poor.

Error occurs in: grid

Action: A finer reference grid should be used or the value of ix in the parameter statements should be increased, such that $ix \ge (md(ir)+1)$. Also, see Section 2.5.1.

10. While using the Stokes wave form of the model, ntype=2 the model may encounter large values of the Ursell number, indicating that the water is too shallow for that model to be appropriate. The cutoff point recognized by the program is $(A/h)/(kh)^2 = 0.5$.

Message: Warning: Ursell number = "U" encountered at grid location "i,j" during calculation of wave component "ii". Should be using Stokes-Hedges model (ntype = 1) due to shallow water.

Result: None by program.

Error occurs in: fdcalc

Action: The program should be terminated and re-run with the composite nonlinear model.

11. The Newton-Raphson iteration for wavenumber k may not converge at a grid location (i,j) of the computatinal grid in the specified number of steps. This may occur for waves on strong opposing currents.

Message: Wavenumber failed to converge for frequency component "ifreq" on column "i".

K = last iterated value of wavenumber

D = depth

T = period calculated from last iterated value of k

U = x-direction velocity

F = value of objective function (should be = 0 for convergence)

Result: Program continues with last iterated value of k. Computed results are of questionable accuracy.

Error occurs in: vwnum

Action: Check if the test case suits the assumptions of the program. Also, an error in the input files might be possible cause.

Results of Calculations

As part of the screen output, the model first prints the x-position for each reference grid row. Also, for each grid row, the values of significant wave height are printed for each reference grid location in 12(f10.4) format using the physical units (MKS or English) used for input.

2.10 Program Output: Stored Output

This section discusses the output stored in two disk files outdat.dat and outrad.dat (logical unit numbers iun(3) and 22, respectively) with a resolution of the reference grid. Also the

optional output file *specamp.dat* (logical unit number 36) will be discussed. All data is stored using free format. The data file *outdat.dat* has been used to generate the plots given in Chapter 3.

Output stored in outdat.dat

Information about the depth grid and the significant wave heights computed at each reference grid point are stored in the file outdat.dat. The program first stores the values of nr and mr, and the values of the reference grid y-locations.

- 1. nr, mr
- 2. yr(jr), jr=1, nr

Then, for each x-position ir in the reference grid, the program stores the following information.

- 3. xr(ir)
- 4. d(ir,jr), jr=1, nr
- 5. h13(ir,jr), jr=1, nr

The set of output (3-5) is repeated for each x-position ir. The data for depth d include the effect of the tidal offset (input as tide) and also will show the thin film where exposed, dry land has been replaced by a small water depth. h13 is the significant wave height at any reference grid location. The data are stored using the units (MKS or English) used for input.

Output stored in outrad.dat

Information about radiation stresses is stored in the file outrad.dat. The program first stores the dimensions of the reference grid nr and mr.

1. nr, mr

Then for each x-position in the reference grid, the following information is stored.

- 2. $S_{xx}(ir,jr), jr=1, nr$
- 3. $S_{xy}(ir,jr), jr=1, nr$
- 4. $S_{yy}(ir, jr), jr=1, nr$

(2-4) is repeated for each row ir. Here S_{xx} is the radiation stress in the x-direction due to the momentum flux in the x-direction, S_{xy} is the radiation stress in the y-direction due to the momentum flux in the x-direction and S_{yy} is the radiation stress in the y-direction due to the momentum flux in the y-direction. The radiation stresses are stored in (kg/s^2) if MKS units are used in the input and in (sl/s^2) if English units are used in the input.

Output stored in specamp.dat

Information about the reference grid and the complex amplitudes of any specified wave component can be stored in the file specamp.dat when the switch iopt is set to unity and the desired component is specified in the input. This option provides the user with the opportunity to examine the behavior of any wave component in more detail. The file specamp.dat is set up similar to outdat.dat with the only difference that complex amplitudes are stored instead of the significant wave height at all reference grid points. The structure of specamp.dat is as follows.

- 1. nr, mr
- 2. yr(jr), jr=1, nr

Then, for each x-position ir in the reference grid, the program stores the following information.

- 3. xr(ir), psibar(ir)
- 4. d(ir,jr), jr=1, nr
- 5. a(ir,jr), jr=1, nr

The set of output (3-5) is repeated for each x-position ir. Here, psibar(ir) is the value of reference phase function with no random phase included and a is the complex amplitude of the chosen wave component.

A listing of the program code is given in Appendix A. Also, an example of a plotting program using data stored in *outdat.dat* is included in Appendix B to provide guidance on reading the stored output file. An *fweb* documentation of the preprocessor **SPECGEN** is given in Appendix C.

3 EXAMPLE CALCULATIONS

This chapter presents calculations performed using the spectral combined refraction-diffraction model REF/DIF S.

Each section of this chapter describes in full the model's application to a specific problem. Following a description of the problem, defining parameters needed to create the input files are stated. Only part of the input file *indat.dat* will be presented for each example case. The rather long part where the user inputs periods and tides for each frequency as well as amplitudes and directions of all wave components for that frequency is not presented. Thus the wave parameters are omitted but the first and last parts of the data file are presented. The omitted lines correspond to items (10–12) described in Section 2.7 of this manual. Reference will be made to preprocessing programs that were used to create the input data.

The data files are then used to run the program REF/DIF S with no job-specific modifications to the program involved. Program output is then presented in such a way as to adequately indicate the results and compared to data at hand.

The output for the various examples has been presented using some plotting programs which are external to the main body of the supplied program REF/DIF S. A specialized program has been included in Appendix B in order to provide some guidance in reading the data files generated by the main program. However, plotting routines are likely to vary from one computer system to another. The extra programs are therefore likely to be extensively machine-specific to the systems on which the computations were performed.

Four example cases will be presented. Results of all four examples will be compared to available laboratory data. The results of the first example case will be compared to experiments described by Mase and Kirby (1992). The rest of the examples will be numerical reproductions of experimental test cases chosen out of a series of tests performed by Vincent and Briggs (1989).

Only two separate depth grids have to be constructed for the simulations of these four example cases. The first domain is a channel with a minimal width, a short region of constant depth and a beach with a 1:20 slope, corresponding to the setup used in the experiments described by Mase and Kirby (1992). The second is a two-dimensional domain with an elliptical shoal over an otherwise constant depth bottom used by Vincent and Briggs during their experiments. Both depth grids may be generated using the preprocessor **DATGEN**. Depth grids are stored in the file refdat.dat.

Section 3.1 provides a simulation of an experiment performed by Mase and Kirby (1992) with waves shoaling and breaking on a plane beach. This experiment only involves a unidirectional frequency spectrum. Section 3.2 presents simulations of multidirectional waves over a submerged shoal. The example cases in this section involve breaking and non-breaking waves, as well as broad and narrow directional distributions. Simulations of examples involving wide directional distributions make extensive use of the wide angle capabilities of **REF/DIF S**.

3.1 Unidirectional Waves Shoaling on a Plane Beach

3.1.1 Experimental Domain

Mase and Kirby (1992) conducted wave flume experiments using a Pierson-Moskowitz spectrum without directional spreading. The waves were generated in constant depth of 47 cm and shoaled and dissipated on a 1:20 beach. Wave gages were placed at several locations on the beach. Figure 12 shows the setup of the experiment. The depth grid for this case can be created using **DATGEN** using the information about the domain in this section. In the simulations of this experiment a constant depth region of 2 m length was used.

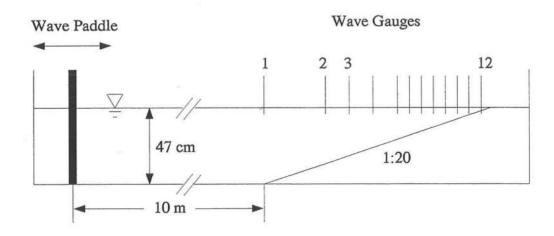


Figure 12: Experimental Setup (from Mase and Kirby (1992))

The comparisons between this data set and the model will be made using Case 1 in Mase and Kirby (1992). The available data are in terms of time series at the wave gage locations. The time series are used to compute the significant wave height for comparison with model results.

3.1.2 Model Comparison

Measured data at the offshore gage is used to create a smoothed incident frequency spectrum for the model input. The incident spectrum is divided into 50 discrete wave components

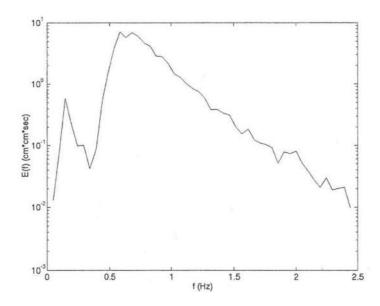


Figure 13: Incident Frequency Spectrum for Unidirectional Waves on a Beach

with constant bin width Δf . Figure 13 shows the incident frequency spectrum. Thus, the measured spectrum is used in the input rather than the target spectrum. The resulting indat.dat file is provided to the user with this manual.

The reference grid consists of 228×5 grid points with dx = 0.05 m and dy = 0.1 m. A depth tolerance value of dt=0.05 m was used. Also, MKS units were used (iu=1) and the linear model was applied (ntype=1). No currents (icur=0) or tides (tide=0) were considered. The program was run using closed lateral boundaries (ibc=0) and no subdivisions were specified by the user (ispace=0). Also no damping (iff(1)=iff(2)=iff(3)=0) and no user-specified subgrids (isp=0) were considered. The option to examine one of the wave components more closely was turned off (iopt=0).

The input data file *indat.dat* is partially given below.

```
1 2 3
228 5
1 0 0 0
0.05 0.1 0.05
0 1
0 0 0
0
1 1
1 50
```

The omitted part mentioned above is inserted here. The last two lines of the file *indat.dat* are as follows.

0.768317 0

Results for this case are presented in Figure 14. It can be seen that the model shoals the waves up to the breaking point. The location of the breaking point and the wave height at that location agree well with the data. The wave height decay after breaking starts is also very well predicted. The data points close to the shore show the presence of setup, which the model cannot predict.

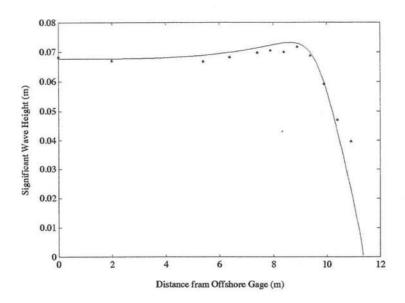


Figure 14: Waves Shoaling on a Beach: Significant Wave Height for Data (*) and Model (—)

3.2 Multidirectional Waves over a Submerged Shoal

3.2.1 Experimental Domain

Vincent and Briggs (1989) conducted their experiments in a wave tank that was 35 m (114 ft) wide and 29 m (96 ft) long. The waves were generated by the directional spectral wave generator, which was 27.43 m (90 ft) long. The center of an elliptical shoal was located at x = 6.10 m and y = 13.72 m. The elliptical shoal had a major radius of 3.96 m, a minor radius of 3.05 m and a maximum height of 30.48 cm at the center. Maximum water depth was 47 cm. Expressions for the shoal perimeter and the elevation of each point on the shoal are given by Vincent and Briggs (1989). Reference should be made to that document for further information about the domain.

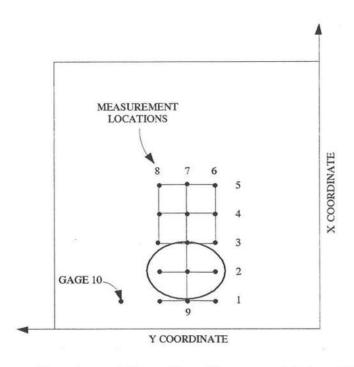


Figure 15: Experimental Setup (from Vincent and Briggs (1989))

Figure 15 shows the experimental setup in the basin and the measurement transects 1–9. Data collection for the cases of interest to this study were performed only along Transect 4.

In all the cases a TMA spectrum (Hughes, 1984) in conjunction with a wrapped normal directional spreading function (Borgman, 1985) was used to establish the target spectrum. The TMA spectrum is given by the energy density E(f) for frequency f

$$E(f) = \frac{\alpha g^2}{(2\pi)^4 f^5} \exp\left\{-1.25 \left(\frac{f_m}{f}\right)^4 + (\ln \gamma) \exp\left[\frac{-(f - f_m)^2}{2\sigma^2 f_m^2}\right]\right\} \phi(f, h)$$
(41)

where α is the Philips' constant, f_m is the peak frequency and γ is the peak enhancement factor. Furthermore, σ is the shape parameter defined by

$$\sigma = \begin{cases} \sigma_a = 0.07 & \text{if } f < f_m \\ \sigma_b = 0.09 & \text{if } f \ge f_m \end{cases}$$

$$\tag{42}$$

The factor $\phi(f, h)$ incorporates the effect of the depth h and is computed following Hughes (1984) by

$$\phi = \begin{cases} 0.5 (\omega_h)^2 & \text{if } \omega_h < 1\\ 1 - 0.5 (2 - \omega_h)^2 & \text{if } 1 \le \omega_h \le 2\\ 1 & \text{if } \omega_h > 2 \end{cases}$$
(43)

where

$$\omega_h = 2\pi f \sqrt{\frac{h}{g}} \tag{44}$$

Narrow or broad frequency spectra used in the experiment can be obtained by assigning the values 20 and 2 to the parameter γ , respectively. Representative plots of the narrow and broad frequency spectra are given in Figure 16.

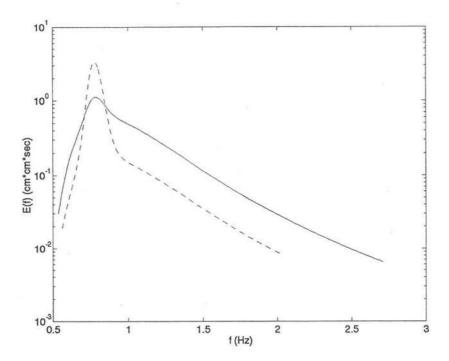


Figure 16: Narrow (---) and Broad (----) Frequency Spectra with $H_{1/3}=0.0254$

The directional spreading function (Borgman, 1985) is given by

$$D(\theta) = \frac{1}{2\pi} + \frac{1}{\pi} \sum_{i=1}^{J} \exp\left[-\frac{(j\sigma_m)^2}{2}\right] \cos j \left(\theta - \theta_m\right)$$
 (45)

where

 $\theta_m = \text{mean wave direction} = 0^{\circ}.$

J = number of terms in the series (chosen to be 20 in the numerical calculations).

The parameter σ_m is chosen to be 10° or 30° giving a narrow or broad directional spreading, respectively. These spreading characteristics were used in the experiment. Representative plots of the narrow and broad directional spreading functions are given in Figure 17.

The two-dimensional spectrum is then given by the product

$$S(f,\theta) = E(f)D(\theta) \tag{46}$$

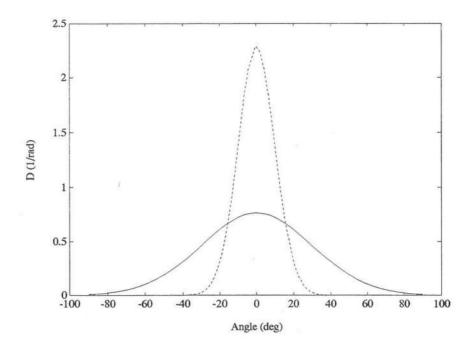


Figure 17: Narrow (---) and Broad (----) Directional Spreading Functions

The two-dimensional spectrum created by these two functions has to be divided up into discrete components to create input for REF/DIF S, which can be performed using SPECGEN. Section 2.6 and Appendix C give more information about this preprocessor.

Also, the reference grid was established and stored in *refdat.dat* using the preprocessor **DATGEN**. No additional information is necessary for this process.

3.2.2 Model Comparison

Three cases out of this data set will be used for comparisons with the model. The first two cases do not involve breaking waves, whereas the last case involves extensive breaking of the waves over the submerged shoal. Table 1 summarizes the wave parameters for these cases.

In order to recreate the results for these three example cases, the user has to create a specgen.dat input file for the preprocessor SPECGEN. The structure of this file is explained in Appendix C. The necessary information about the test cases can be found in this section as well as the previous section. The preprocessor SPECGEN will create an indat.dat file, that can be used to run REF/DIF S. The indat.dat files for the examples are included below, reference can be made to those listings.

The bottom bathymetry input file refdat.dat can be created using **DATGEN**. Note that

Table 1: Test Conditions for Waves over a Submerged Shoal (Vincent and Briggs, 1989)

Test Type	Case ID	T_p (sec)	$H_{1/3}$ (cm)	α	γ	σ_m
$Non ext{-}breaking$	N4	1.3	2.54	0.00047	20	10
	B4	1.3	2.54	0.00047	20	30
Breaking	B5	1.3	19.00	0.08650	2	30

DATGEN will give the user the option to create an *indat.dat* manually. The user must skip that option as the necessary *indat.dat* is created by **SPECGEN** and would otherwise be overwritten.

Non-Breaking Waves

These cases do not involve breaking waves and are simulated to show that the breaking term that is always in effect does not induce an unrealistic wave height reduction.

Case N4

The first case is denoted as Case N4 by Vincent and Briggs (1989) and involves a narrow frequency spectrum used in conjunction with a narrow directional spreading. Distinct values are given to the free parameters to define the spectrum.

$$H_{1/3} = 0.0254m$$
 $f_m = \frac{1}{1.3sec} = 0.769231Hz$
 $f_{max} = 5Hz$
 $\gamma = 20$
 $\theta_m = 0^{\circ}$
 $\sigma_m = 10^{\circ}$

SPECGEN was used to generate the input data file *indat.dat*. The spectrum was divided into 10 frequency components and 20 directional components, thus a total of 200 waves were present.

The model was run using grid dimensions 251×275 with a resolution of dx = dy = 0.1 m. A depth tolerance value of dt=0.05 m was used. MKS units were used (iu=1) and

the Stokes-to-Hedges nonlinear model was applied to test the behaviour of the modified nonlinear terms in the absense of breaking waves (ntype=1). No currents (icur=0) or tides (tide=0) were considered. The program was run using partially transmitting boundaries (ibc=1) and no subdivisions were specified by the user (ispace=0). Also no damping (iff(1)=iff(2)=iff(3)=0) and no user-specified subgrids (isp=0) were considered. The option to examine one of the wave components more closely was turned off (iopt=0).

Part of the input file *indat.dat* is given below. Note that the first seven lines are identical with the first seven lines of the input file *specgen.dat* for the preprocessor **SPECGEN**, which is documented in Appendix C. Also, see Section 2.6.

```
1 2 3

251 275

1 1 0 1

1.00000E-01 1.00000E-01 5.00000E-02

0 1

0 0 0

0

1 1

1 10
```

The omitted part consisting of the amplitudes and directions of the individual components is inserted here. The last two lines of the file *indat.dat* are as follows.

```
0.769231
```

The output is presented in two ways. Figure 18 shows a contour plot of the normalized significant wave height. Figure 19 shows a transect at x = 12.2 m and the comparison to available data at that transect. It can be seen that the general shape of the curve in Figure 19 is well modeled. The fact that the peak of the wave height for this non-breaking wave case is well represented shows that the coefficient of the breaking term stays small and does not dissipate significant energy from the system.

Case B4

In case B4 a spectrum with narrow frequency and broad directional spreading is generated. The free parameters have the following values.

$$H_{1/3} = 0.0254m$$

$$f_m = \frac{1}{1.3sec} = 0.769231Hz$$

$$f_{max} = 5Hz$$

$$\gamma = 20$$

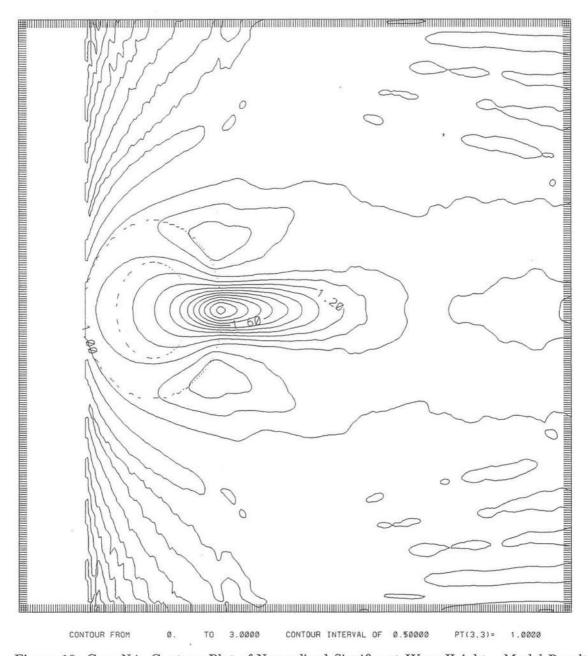


Figure 18: Case N4: Contour Plot of Normalized Significant Wave Height – Model Results

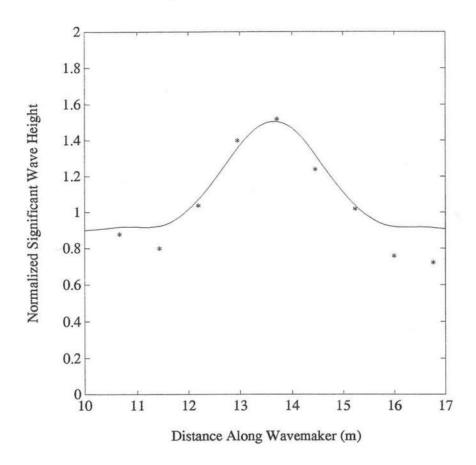


Figure 19: Case N4: Normalized Significant Wave Height for Data (*) and Model Results (\longrightarrow) at Transect at x = 12.2 m

$$\theta_m = 0^{\circ}$$
 $\sigma_m = 30^{\circ}$

SPECGEN was used to generate the input data file *indat.dat*. Again the spectrum was divided into 10 frequency components and 20 directional components, thus a total of 200 waves were present.

The parameters defining the domain and controling the program options are identical to those of Case N4 with the exception of *ntype*. In this case the linear model is used as the satisfactory behavior of the Stokes-to-Hedges nonlinear model was shown for Case N4.

The input data file *indat.dat* is given partially below. Again note that the first seven lines of this file correspond to the first seven lines of the input file *specgen.dat* for **SPECGEN**, which is documented in Appendix C. Also, see Section 2.6.

1 2 3

```
251 275

1 0 0 1

1.00000E-01 1.00000E-01 5.00000E-02

0 1

0 0 0

1 1

1 10
```

The omitted part mentioned above is inserted here. The last two lines of the file *indat.dat* are as follows.

Figure 20 shows a contour plot of the normalized significant wave height. Figure 21 shows a transect at x = 12.2m and the comparison to available data at that transect. These results also confirm that the breaking term in the model remains sufficiently small for non-dissipative cases.

Breaking Waves

Case B5

In this example Case B5 of the Vincent and Briggs experiments is modeled. The parameters used as input to the preprocessor SPECGEN to create the input file for REF/DIF S are as follows.

$$H_{1/3} = 0.19m$$
 $f_m = \frac{1}{1.3sec} = 0.769231Hz$
 $f_{max} = 5Hz$
 $\gamma = 2$
 $\theta_m = 0^{\circ}$
 $\sigma_m = 30^{\circ}$

Thus, a random sea with a broad frequency and a broad directional spectrum is generated. The spectrum is again divided into 10 frequency and 20 directional components, yielding 200 wave components. The Stokes-to-Hedges nonlinear model is used as strong nonlinear effects occur in this case. Again no bottom friction, no tides or currents were considered. The reference grid is constructed with the same resolution as in the first two cases simulated.

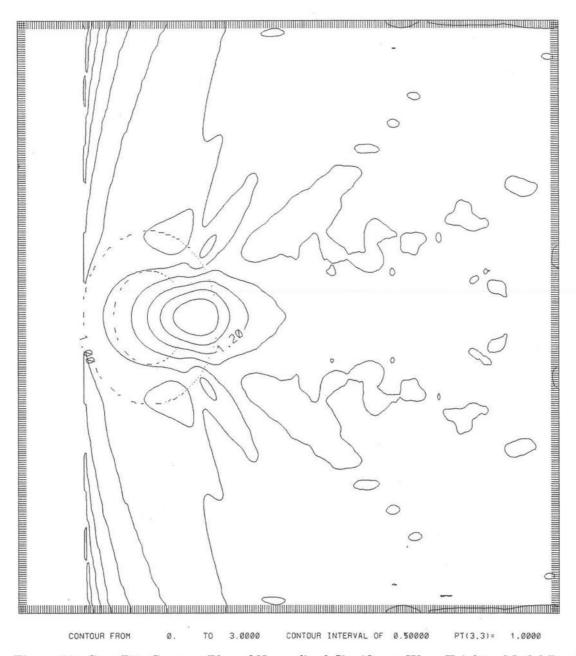


Figure 20: Case B4: Contour Plot of Normalized Significant Wave Height - Model Results

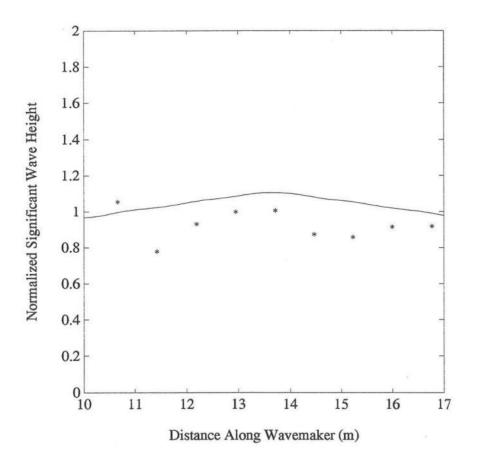


Figure 21: Case B4: Normalized Significant Wave Height for Data (*) and Model Results (—) at Transect at x = 12.2 m

and all other model controlling parameters are the same as well. The reader is referred to the discussion of Cases N4 and B4 for that information.

Part of the input file *indat.dat* is listed below. Note that the first seven lines are identical to those of the input file *specgen.dat* for the program **SPECGEN**, which is documented in Appendix C. Also, see Section 2.6.

```
1 2 3

251 275

1 1 0 1

1.00000E-01 1.00000E-01 5.00000E-02

0 1

0 0 0

0

1 1

1 10
```

The last two lines of the input file are as follows.

0.769231

Results are shown in Figure 22 in terms of a contour plot of the normalized significant wave height and in Figure 23 which shows a comparison between the data and the results of the model for the normalized significant wave height at the transect at x = 12.2 m.

The results show that the model is partially successful in predicting the defocusing effect behind the shoal that arises because of nonlinear effects. This effect can not be predicted using the linear model. It is seen though that the data shows a tendency to recover the initial significant wave height at y = 10 m and y = 17 m, the model does not predict this behavior. The discrepency between the data and the model results is believed to arise because the breaking model currently used in **REF/DIF** S does not take directional effects into account.

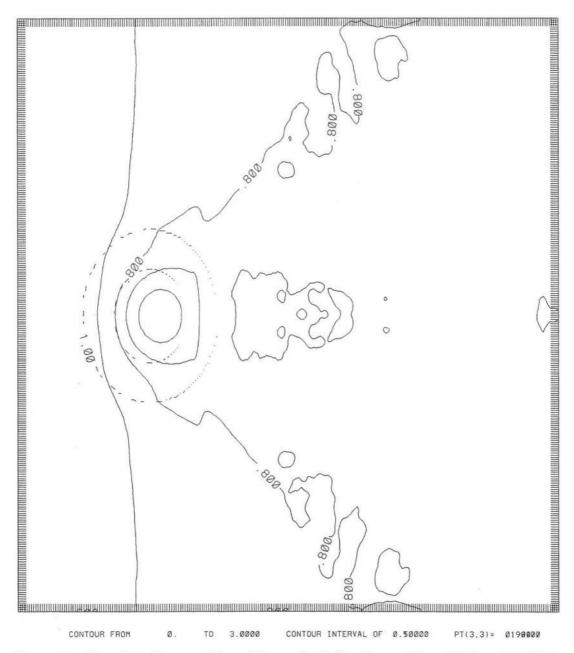


Figure 22: Case B5: Contour Plot of Normalized Significant Wave Height – Model Results

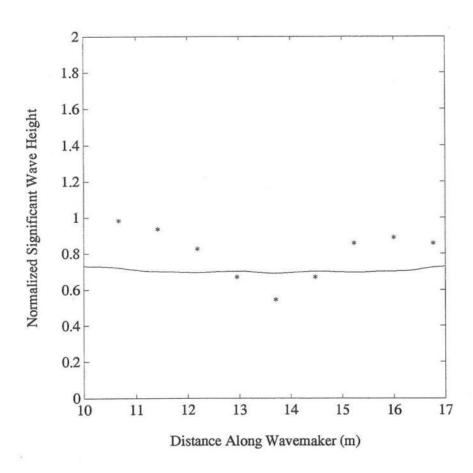


Figure 23: Case B5: Normalized Significant Wave Height for Data (*) and Model Results (—) at Transect at x = 12.2 m

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5 Appendix A: REF/DIF S Program Listing

The following section contains the listing for REF/DIF S. An attempt has been made to make the code readible by using comments at appropriate locations in the program. The anotations may be improved significantly using the *fweb* documentation program (Krommes, 1992).

The main program REF/DIF S and the eleven subroutines are listed here in the following order:

- 1. refdifs
- 2. inref
- 3. inwave
- 4. model
- 5. grid
- 6. con
- 7. fdcalc
- 8. rbcon
- 9. vwnum
- 10. vtrida
- 11. diss

```
C*
     program: ref/dif s
C*
     version: 1.0
C*
C*
     program to calculate the forward scattered wave field in
C*
C*
     regions with slowly varying depth and current, including the
     effects of refraction and diffraction. program is based on
C*
     parabolic equation method.
     parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
    1nnd=20,nnii=200)
     common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
    1dconv2(2)
     common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
    1ioutput
     common/ref3/dxr,dyr,xr(ixr),yr(iyr),x(ix),y(iy)
     common/ref4/isd(ixr,iyr)
     common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
     common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
     common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
     common/nlin/an,anl,ntype
     common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
    1, nii, fpeak
     common/wav2/amp(ncomp,nnd),dir(ncomp,nnd),tide
     common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
     real k, kb
     complex a, w, az
     dconv(1)=1.
     dconv(2)=0.30488
     dconv2(1)=1.
     dconv2(2)=14.594
C*----
     specify unit number for line printer output. standard choice
C*
C*
     would be *; some systems may vary.
     this line must be set during installation.
C*----
     print headers on output
     write(*,100)
     read control parameters and reference grid data
     call inref
C*-----
     read control parameters and initializing wave data
     call inwave
                ______
     pass program control to subroutine model
     for each frequency component specified in inwave, model
C*
```

```
executes the model throughout the entire grid and then
C*
     reinitializes the model for the next frequency
C*
     all done
     close output data files if open and close statements are
C*
C*
     being used.
C*----
     do 1 i=1,3
     close(iun(i))
     continue
     close(22)
     if(iopt.eq.1) close(36)
     stop
 100 format(/////20x, 'Refraction-Diffraction Model for'/
    120x, 'Spectral Wave Conditions'///
    120x, 'REF/DIF S Version 1.1'///
    120x, 'Center for Applied Coastal Research'/
    120x, 'Department of Civil Engineering'/
    120x, 'University of Delaware'/
    120x, 'Newark, Delaware 19716'///
    105x,'James T. Kirby and H.Tuba Ozkan, July 1992, February 1994')
     end
```

```
C*
      subroutine inref
C*
      subroutine reads in and checks dimensions and values for
c*
      large scale reference grid. wave parameters for the
C*
      particular run are read in later by subroutine inwave.
C*
      the following unit (device) numbers are assumed:
C*
C*
            unit
                       contents
C*
C*
                       input reference grid values of d, u, and v
C*
             iun(1)
                       input user specified subgrid divisions
             iun(2)
c*
                       output results at reference grid locations to
             iun(3)
C*
                       disk file
C*
                       assumed unit for line printer output
C*
                       (must be specified in program, machine
C*
                        specific. * is sun workstation version)
c*
                       dimension information and run instructions
             iun(5)
c*
                       (must be specified in program, machine
C*
                        specific)
C*
C*
      variable definitions
c*
C*
          mr,nr - reference grid dimensions
C*
          dxr, dyr-
                     grid spacing for reference grid
C*
                    physical unit descriptor (1=mks, 2=english)
C*
                     default value is 1, mks units
C*
                  - depth tolerence value (to check for anomalous
C*
          dt
                     depth values)
c*
                     switch to control grid subdivision.
c*
                     =0, program attempts its own subdivisions
C*
                     =1, user specifies subdivisions
C*
                    y direction subdivision (ispace=0 or 1)
C*
          nd
C*
                     (must be .lt. iy/nr-1)
          md(mr-1) - x direction subdivisions (if ispace=1)
C*
C*
                     (must be .le. ix-1)
                     nonlinearity control parameter
C*
          ntype -
                     =0, linear model
C*
                     =1, stokes matched to hedges in shallow water
C*
                     =2, stokes throughout
c*
                    switch to tell program if current data is to
C*
          icur
                     be used and read on input
c*
C*
                     =0, no input current data
                     =1, input current data to be read
C*
                     program defaults to icur=0
C*
          ibc
                     boundary condition switch
                     =0, use closed lateral boundaries
C*
C*
                     =1, use open lateral conditions
                     program defaults to ibc=0
C*
C*
          dr

    depths at grid points
```

```
>0, submerged areas
C*
                <0, elevation above surface datum
C*
              - x velocities at grid points
C*
        117
                (only entered if icur=1)
             - y velocities at grid points
        vr
C*
                (only entered if icur=1)
C*
C*
    subroutine is called from main and returns control to
C*
    calling location
                     _____
    parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
    1nnd=20, nnii=200)
    common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
    1dconv2(2)
    common/ref2/dr(ixr,iyr), ur(ixr,iyr), vr(ixr,iyr),iun(8),iinput,
    1ioutput
    common/ref3/dxr,dyr,xr(ixr),yr(iyr),x(ix),y(iy)
     common/ref4/isd(ixr,iyr)
    common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
    common/nlin/an,anl,ntype
    constants
    define device number iun(5) for reference grid input
    this line must be set during program installation, prior to
C*
    compilation. also construct an open statement if needed.
    open(unit=iun(5),file='indat.dat',status='old')
C*-----
    remaining unit numbers in i4 format
C*----
    read(iun(5),*)(iun(i),i=1,3)
    open(unit=iun(1),file='refdat.dat',status='old')
     open(unit=iun(2),file='subdat.dat')
    open(unit=iun(3),file='outdat.dat')
    open(unit=22,file='outrad.dat')
C*----
    headers on output file *
C*----
    write(*,106)
C*-----
    read control data from unit iun(5)
    read(iun(5),*) mr,nr
    read(iun(5),*) iu, ntype, icur, ibc
    read(iun(5),*) dxr, dyr, dt
    read(iun(5),*) ispace, nd
     if(ispace.eq.1) then
    read(iun(5),*) (md(i),i=1,mr-1)
     else
     endif
```

```
write(*,107) mr,nr,dxr,dyr
     if(iu.eq.1) write(*,114) iu
     if(iu.eq.2) write(*,115) iu
     if(icur.eq.0) write(*,200)
     if(icur.eq.1) write(*,201)
     if(ibc.eq.0) write(*,202)
     if(ibc.eq.1) write(*,203)
     if(ispace.eq.0)write(*,108)
     if(ispace.eq.1)write(*,109)
     write(*,119) nd
     if(ntype.eq.0) write(*,110)
     if(ntype.eq.1) write(*,111)
     if(ntype.eq.2) write(*,112)
C*----
     check input from unit iun(5)
     if((mr.gt.ixr).or.(nr.gt.iyr)) then
       stop 'dimensions for reference grid too large, stopping'
     else
     end if
     if((iu.ne.1).and.(iu.ne.2)) iu=1
     dt=dt*dconv(iu)
     dxr=dxr*dconv(iu)
     dyr=dyr*dconv(iu)
     if(dt.eq.0.) dt=2.
     if(nd.gt.(ifix(float(iy-1)/float(nr-1)))) then
      write(*,102) nd
       stop
     else
     endif
     if(ispace.eq.1) then
      test=0.
      do 1 i=1,mr-1
      if(md(i).gt.(ix-1)) then
        write(*,103) md(i),i
        test=1.
       else
       endif
1
       continue
     if(test.eq.1.) stop
     else
     end if
C*----
     read depth grid and velocities from unit iun(1)
C*----
     do 2 i=1,mr
     read(iun(1),101)(dr(i,j),j=1,nr)
2
     continue
     if(icur.eq.1)then
     do 3 i=1,mr
     read(iun(1),101)(ur(i,j),j=1,nr)
3
     continue
     do 4 i=1.mr
     read(iun(1),101)(vr(i,j),j=1,nr)
```

```
continue
    endif
            ______
    convert depth and currents
C*----
    do 5 i=1,mr
    do 5 j=1,nr
    dr(i,j)=dr(i,j)*dconv(iu)
5
    continue
    if (icur.eq.1) then
    do 55 i=1,mr
    do 55 j=1,nr
    ur(i,j)=ur(i,j)*dconv(iu)
    vr(i,j)=vr(i,j)*dconv(iu)
55
    continue
    endif
    check for large depth changes and large currents
    in reference grid data
C*----
    do 6 i=2,mr-1
    do 6 j=2,nr-1
     dcheck=(dr(i+1,j)+dr(i-1,j)+dr(i,j-1)+dr(i,j+1))/4.
    if(abs(dcheck-dr(i,j)).gt.dt)
    1 write(*,104) dr(i,j), i,j,dt
    continue
6
    if(icur.eq.1) then
    do 7 i=1,mr
     do 7 j=1,nr
     if(dr(i,j).le.0.0) go to 7
     fr=(ur(i,j)*ur(i,j)+vr(i,j)*vr(i,j))/(g*dr(i,j))
     if(fr.gt.1.) write(*,105) i,j,fr
    continue
7
     endif
C*----
  establish coordinates for reference grid
    do 8 ir=1,mr
    xr(ir)=float(ir-1)*dxr
8
     continue
     do 9 jr=1,nr
     yr(jr)=float(jr-1)*dyr
     continue
C*----
    establish y coordinates for interpolated grid
    n=nd*(nr-1)+1
     dy=dyr/float(nd)
     do 10 j=1,n
     y(j)=float(j-1)*dy
     continue
              _____
     write grid information on output unit iun(3)
```

```
write(iun(3),*) nr,mr
     write(iun(3),*) (yr(jr)/dconv(iu),jr=1,nr)
     read in changes for friction values
c*
         if(1)=1, turbulent boundary layer damping everywhere
c*
         if(2)=1, porous bottom damping everywhere
C*
         if(3)=1, laminar boundary layer damping everywhere
     read(iun(5),*) (if(i),i=1,3)
     do 11 i=1,3
     if((if(i).ne.0).and.(if(i).ne.1)) if(i)=0
11
     continue
     write(*,116)(if(i),i=1,3)
     specify whether or not user specified subgrids are to be
     read in during model operation.
C*
c*
     isp=0, no subgrids specified
     isp=1, subgrids to be read in later from unit iun(2)
C*----
     read(iun(5),*) isp
      if(isp.eq.0) write(*,117)
      if(isp.eq.1) write(*,118)
      if((isp.eq.1).and.(ispace.eq.0))write(*,113)
      if(isp.eq.0)then
      do 14 ir=1,mr
      do 14 jr=1,nr
      isd(ir,jr)=0
      continue
 14
      else
      do 15 ir=1,mr-1
      read(iun(2),100)(isd(ir,jr),jr=1,nr-1)
 15
     continue
      end if
     input done, return to main program
     return
 100 format(15i4)
 101 format(16f8.4)
 102 format('y-direction subdivision nd=',i4,'too fine.'/
     1' maximum number of y grid points will be exceeded.'/
     1' execution terminating.')
 103 format('x-direction subdivision md=',i4,
     1'too fine on grid block',2x,i3/' execution terminating')
 104 format(' depth', 2x, f7.2, '(m) at reference grid location',
     12(2x,i3)/' differs from the average of its neighbors by',
     1' more than', 2x, f7.2, '(m).'/' execution continuing')
 105 format(' ambient current at reference grid location'
     1,2(2x,i3),' is supercritical with froude number =',f7.4/
     1' execution continuing')
 106 format('0'///20x,'input section, reference grid values'///)
 107 format(' reference grid dimensions mr=',i3/
                                         nr=',i3///
     1
```

```
' reference grid spacings
                                         dxr=',f8.4/
                                         dyr=',f8.4)
108 format(' '/' ispace = 0 chosen, program will attempt its own ',
    1'reference grid subdivisions')
109 format(' '/' ispace =1 chosen, subdivision spacings will be',
    1' input as data')
110 format(' '/' ntype = 0, linear model')
111 format(' '/' ntype = 1, stokes model matched to hedges model')
112 format(' '/' ntype = 2, stokes model')
113 format(' warning: input specifies that user will be supplying',
    1' specified subgrids (isp=1),'/
    1' while program has been told to generate its own subgrid',
    1' spacings (ispace=0).'/
    1' possible incompatibility in any or all subgrid blocks')
114 format(' '/' physical unit switch iu=',i1,
    1', input in mks units')
115 format(' '/' physical unit switch iu=',i1,
    1', input in english units')
116 format(' '//' switches for dissipation terms'//
    1' ', i1, '
               turbulent boundary layer'/
    1' ',i1,' porous bottom'/
    1' ',i1,' laminar boundary layer')
117 format(' '/' isp=0, no user defined subgrids')
118 format(' '/' isp=1, user defined subgrids to be read')
119 format(' '/' y-direction subdivision according to nd=',i3)
200 format(' '/' icur=0, no current values read from input files')
201 format(' '/' icur=1, current values read from data files')
202 format(' '/' ibc=0, closed (reflective) lateral boundaries')
203 format(' '/' ibc=1, open lateral boundaries')
```

```
_____
C*
      subroutine inwave
C*
     read in wave parameters.
C*
      variable definitions
C*
C*
          iinput - determine method of specifying the first
C*
                     row of computational values
C*
                     =1, input values from indat.dat according to
C*
                     value of iwave
C*
         =2, input complex a values from file wave.dat
C*
C*
                     determine whether last row of complex amplitudes
C*
          ioutput -
                     are stored in separate file owave.dat
C*
                     =1, extra data not stored
C*
                      =2, extra data stored in file owave.dat
C*
C*
      if iinput=1:
C*
C*
C*
                     input wave type
                      =1, input discrete wave amplitudes and
C*
                          directions
C*
                      =2, read in dominant direction, total average
C*
                          energy density, and spreading factor
C*
                     number of frequency components to be used
                      (separate model run for each component)
C*
                     wave frequency for each of nfreqs runs
C*
          freqs
                  - tidal offset for each of nfreqs runs
          tide
C*
c*
      if iwave = 1
C*
C*
                  - number of discrete wave components at
C*
                     each of nfreqs runs
C*
          amp
                  - initial amplitude of each component
C*
                  - direction of each discrete component in
          dir
C*
C*
                     + or - degrees from the x-direction
C*
      if iinput=2:
C*
C*
C*
                  - wave frequency for one run
         freqs
                  - tidal offset for one run
C*
         tide
      parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20, nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
     1dconv2(2)
      common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
     1ioutput
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1,nii,fpeak
```

```
common/wav2/amp(ncomp,nnd),dir(ncomp,nnd),tide
     pi=3.1415927
     read value of iinput, ioutput
     read(iun(5),*) iinput, ioutput
     if(iinput.ne.1) iinput=1
     if(ioutput.ne.1) ioutput=1
     if(iinput.eq.1)then
     write(*,*)'iinput = 1, program specifies initial row of a'
     read iwave, nfreqs
     write(*,102)
     read(iun(5),*) iwave, nfreqs
     if(iwave.ne.1) iwave=1
     write(*,103)
     if(nfreqs.gt.ncomp) then
     write(*,104)
     stop
     endif
     write(*,105) nfreqs
     for each frequency, enter the wave period and tidal offset
     do 3 ifreq=1,nfreqs
     read(iun(5),*) freqs(ifreq), tide
     write(*,107) ifreq,freqs(ifreq),tide
     freqs(ifreq)=2.*pi/freqs(ifreq)
     tide=tide*dconv(iu)
     if iwave = 1, read the number of discrete components
C*----
     if(iwave.eq.1) then
     read(iun(5),*) nwavs
     if(nwavs.gt.nnd) then
     write(*,109)
     stop
     endif
     do 1 iwavs=1,nwavs
     read(iun(5),*) amp(ifreq,iwavs), dir(ifreq,iwavs)
     write(*,106) iwavs,amp(ifreq,iwavs),dir(ifreq,iwavs)
     dir(ifreq,iwavs)=dir(ifreq,iwavs)*pi/180.
     amp(ifreq,iwavs)=amp(ifreq,iwavs)*dconv(iu)
     continue
1
     else
     end if
     continue
3
     endif
     return
100 format(15i4)
101 format(16f8.4)
102 format('1'///20x,' input section, wave data values'///)
103 format(' '///' iwave=1, discrete wave amps and directions')
```

- 104 format('Too many frequency components; stopping')
- 106 format(' '/' wave component ',i2,', amplitude =',f8.4, 1', direction=',f8.4)
- 107 format(' '//' frequency component ',i2//
 1' wave period=',f8.4,'sec., tidal offset=',f8.4)
- 108 format(' '/' total variance density =',f8.4,', spreading factor 1n=',i2,' seed number =',i5)
- 109 format('Too many directional components; stopping')
 end

```
C*
      subroutine model
C*
      this subroutine is the control level for the actual wave
C*
      model. data read in during inref and inwave is conditioned
C*
      and passed to the wave model. this routine is executed
C*
      once for each frequency component specified in inwave.
C*
C*
      the wave model is split in three parts which are run
C*
      sequentially for each reference grid row.
C*
c*
                - subroutine performs the interpolation of
c*
         grid
                   depth and current values.
C*
C*
                - calculate the constants needed by the finite
         con
                   difference scheme.
c*
         fdcalc - perform the finite difference calculations.
c*
      parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20, nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
     1dconv2(2)
      common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
     1ioutput
      common/ref3/dxr,dyr,xr(ixr),yr(iyr),x(ix),y(iy)
      common/ref4/isd(ixr,iyr)
      common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
      common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
      common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
      common/nlin/an,anl,ntype
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1, nii, fpeak
      common/wav2/amp(ncomp,nnd),dir(ncomp,nnd),tide
      common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
      dimension s(iy),th(iy,nnii),sxy(iy),sxx(iy),syy(iy)
      dimension npts(ncomp), sumk(ncomp), xd(iyr), xu(iyr), xk(iyr, ncomp)
      complex a, w, az
      real k,kb
      g=9.80621
      rho=1000.
      pi=3.1415927
      eps=1.0e-05
      nii=nwavs*nfregs
      if(nii.eq.1) then
      fpeak=freqs(1)/(2.*pi)
      iopt=0
      else
      read(iun(5),*) fpeak
      read(iun(5),*) iopt
      if(iopt.eq.1) then
      read(iun(5),*) ifcomp,idcomp
```

```
icomp=nwavs*(ifcomp-1)+idcomp
     open(unit=36,file='specamp.dat')
     write(36,*) nr,mr
     write(36,*)(yr(jr)/dconv(iu),jr=1,nr)
     else
     iopt=0
     endif
     endif
     specify initial nonlinear parameters once
C*----
     if(ntype.eq.0) an=0.
     if(ntype.ne.0) an=1.
     if(ntype.ne.2) anl=0.
     if(ntype.eq.2) anl=1.
    Compute kb on first row, for use in specifying initial condition
     do 32 jr=1,nr
     xd(jr)=dr(1,jr)
     xu(jr)=ur(1,jr)
 32 continue
     call vwnum(xd,xu,freqs,xk,eps,nfreqs,nr)
     do 35 ifreq=1,nfreqs
     npts(ifreq)=0.
     sumk(ifreq)=0.
 35 continue
     do 36 ifreq=1,nfreqs
     do 37 jr=1,nr
     if(dr(1,jr).gt.0.05) then
     sumk(ifreq)=sumk(ifreq)+xk(jr,ifreq)
     npts(ifreq)=npts(ifreq)+1
     endif
  37 continue
     if(npts(ifreq).eq.0) then
     kb(1,ifreq)=xk(1,ifreq)
     kb(1,ifreq)=sumk(ifreq)/float(npts(ifreq))
     endif
  36 continue
C*-----
     Set-up initial condition
     do 2 ii=1,nii
     do 3 j=1,n
     do 4 i=1,2
     a(i,j,ii)=cmplx(0.,0.)
   4 continue
   3 continue
   2 continue
     do 5 j=1,n
     do 6 ifreq=1,nfreqs
     do 7 iwavs=1,nwavs
     ii=nwavs*(ifreq-1)+iwavs
```

```
a(1,j,ii)=amp(ifreq,iwavs)*cexp(cmplx(0.,kb(1,ifreq)
    1*sin(dir(ifreq,iwavs))*y(j)))
    istore(ii)=ifreq
    th(j,ii)=dir(ifreq,iwavs)
   7 continue
   6 continue
   5 continue
   compute significant waveheight hs for initial condition
    if(nii.eq.1) then
    do 31 j=1,n
    h13(j)=2.*cabs(a(1,j,1))
 31 continue
    else
    do 13 j=1,n
    s(j)=0.
    do 14 ii=1,nii
    s(j)=s(j)+((cabs(a(1,j,ii)))**2.)
  14 continue
    h13(j)=sqrt(8.*s(j))
  13 continue
    endif
           ______
    initial condition set up
    now execute model for each grid block
C*----
    do 20 ii=1,nii
    psibar(ii)=0.
  20 continue
    do 100 ir=1,(mr-1)
    establish interpolated grid block for segment ir
C*----
    call grid(ir)
    write initial values on iun(3)
    if(ir.eq.1) then
    write(iun(3),*) x(1)/dconv(iu)
    write(iun(3),*)(d(1,j)/dconv(iu),j=1,n,nd)
    write(iun(3),*)(h13(j)/dconv(iu),j=1,n,nd)
    if(iopt.eq.1) then
    write(36,*) x(1)/dconv(iu),psibar(icomp)
    write(36,*)(d(1,j)/dconv(iu),j=1,n,nd)
    write(36,*)(a(1,j,icomp)/dconv(iu),j=1,n,nd)
    endif
    endif
    do 11 ii=1,nii
    do 12 j=1,n
    az(j,ii)=a(1,j,ii)
  12 continue
  11 continue
```

```
calculate constants for first row
    if(ir.eq.1) then
    call con(ir,1,1)
    endif
C*----
    Line printer output for first row
    if(ir.eq.1) then
    do 21 j=1,n
    sxx(j)=0.
    syy(j)=0.
    sxy(j)=0.
  21 continue
    do 25 j=1,n
    do 22 ii=1,nii
    sxx(j)=sxx(j)+(cabs(a(1,j,ii))**2)
    1*(q(1,j,istore(ii))*(1+(cos(th(j,ii))**2))-0.5)
    syy(j)=syy(j)+(cabs(a(1,j,ii))**2)
    1*(q(1,j,istore(ii))*(1+(sin(th(j,ii))**2))-0.5)
    sxy(j)=sxy(j)+q(1,j,istore(ii))
    1*(cabs(a(1,j,ii))**2)*sin(th(j,ii))
  22 continue
    sxx(j)=sxx(j)*rho*g/8.
    syy(j)=syy(j)*rho*g/8
    sxy(j)=sxy(j)*rho*g/16.
  25 continue
                _____
    Line printer output
    mm1=m-1
    write(*,205) ir,mm1
    write(*,202) x(1)/dconv(iu)
    write(*,206)
    write(*,203) (h13(j)/dconv(iu),j=1,n,nd)
    Output of radiation stresses
    write(22,*) nr,mr
    write(22,*) (sxx(j)/dconv2(iu), j=1,n,nd)
    write(22,*) (sxy(j)/dconv2(iu), j=1,n,nd)
    write(22,*) (syy(j)/dconv2(iu),j=1,n,nd)
    endif
C*----
    for next row, evaluate constants
    do 9 icount=2,m
    call con(ir,icount,2)
C*----
    perform finite differencing calculations
C*----
    call fdcalc(ir,icount)
C*----
    roll back values to row 1
```

```
call rbcon
   9 continue
    write out reference grid data on disk file iun(3)
C*----
     write(iun(3),*) x(m)/dconv(iu)
     write(iun(3),*)(d(m,j)/dconv(iu),j=1,n,nd)
     write(iun(3),*)(h13(j)/dconv(iu),j=1,n,nd)
     if(iopt.eq.1) then
     write(36,*) x(m)/dconv(iu),psibar(icomp)
     write(36,*)(d(m,j)/dconv(iu),j=1,n,nd)
     write(36,*)(a(1,j,icomp)/dconv(iu),j=1,n,nd)
     endif
 134 continue
 100 continue
202 format(' x=',f10.2)
203 format(' ',12(f10.4))
205 format(' grid row ir=',i3,', ',i3,' x-direction subdivisions',
    1' used')
206 format(' Significant Wave Height')
     return
     end
```

```
C*
      subroutine grid(ir)
C*
     interpolate the depth and current grids for reference
C*
     grid block ir.
     parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
    1nnd=20, nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
    1dconv2(2)
      common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
    1ioutput
      common/ref3/dxr,dyr,xr(ixr),yr(iyr),x(ix),y(iy)
      common/ref4/isd(ixr,iyr)
      common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
      common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1, nii, fpeak
      common/wav2/amp(ncomp,nnd),dir(ncomp,nnd),tide
      dimension dref(iy),npts(ncomp),sumk(ncomp),xu(iy),xk(iy,ncomp)
      real k,kb
      complex w
     pi=3.1415927
      eps=1.0e-05
     perform y-interpolation on reference grid
      interpolate first row
C*----
      do 10 j=1,n,nd
      d(1,j)=dr(ir,((j-1)/nd+1))
      u(1,j)=ur(ir,((j-1)/nd+1))
      v(1,j)=vr(ir,((j-1)/nd+1))
   10 continue
      if(nd.gt.1) then
      do 12 jj=2,nr
      do 11 j=1,(nd-1)
      jjj=nd*(jj-2)+(j+1)
      d(1,jjj)=(dr(ir,jj)-dr(ir,jj-1))*y(jjj)/dyr
     1+(yr(jj)*dr(ir,jj-1)-yr(jj-1)*dr(ir,jj))/dyr
      u(1,jjj)=(ur(ir,jj)-ur(ir,jj-1))*y(jjj)/dyr
     1+(yr(jj)*ur(ir,jj-1)-yr(jj-1)*ur(ir,jj))/dyr
      v(1,jjj)=(vr(ir,jj)-vr(ir,jj-1))*y(jjj)/dyr
     1+(yr(jj)*vr(ir,jj-1)-yr(jj-1)*vr(ir,jj))/dyr
   11 continue
   12 continue
      set number of x points and define x values
```

```
if(ispace.eq.0) then
  ispace=0, program sets subdivisions
  do 13 j=1,n
   dref(j)=d(1,j)+tide
   if(dref(j).lt.0.001) dref(j)=0.001
13 continue
  do 27 j=1,n
  xu(j)=u(1,j)
27 continue
   call vwnum(dref,xu,freqs,xk,eps,nfreqs,n)
   do 29 ifreq=1,nfreqs
   do 28 j=1,n
  k(1,j,ifreq)=xk(j,ifreq)
28 continue
29 continue
   do 14 ifreq=1,nfreqs
   npts(ifreq)=0
   sumk(ifreq)=0.
14 continue
   do 15 ifreq=1,nfreqs
   do 16 j=1,n
   if(d(1,j).gt.0.05) then
     sumk(ifreq)=sumk(ifreq)+k(1,j,ifreq)
     npts(ifreq)=npts(ifreq)+1
   else
   endif
16 continue
15 continue
   do 17 ifreq=1,nfreqs
   if(npts(ifreq).eq.0) then
   kb(1,ifreq)=k(1,1,ifreq)
   kb(1,ifreq)=sumk(ifreq)/float(npts(ifreq))
   endif
17 continue
   find maximum wavenumber
   big=kb(1,1)
   if(nfreqs.eq.1) go to 40
   do 18 ifreq=2,nfreqs
   if(kb(1,ifreq).gt.big) big=kb(1,ifreq)
18 continue
   compute spacing
40 alw=2.*pi/big
   anw=dxr/alw
   np=ifix(5.*anw)
   if(np.lt.1) np=1
   md(ir)=min((ix-1),np)
   if(np.gt.(ix-1)) write(*,100) np,ir
```

```
else
     ispace=1, user specified subdivision
     endif
     m=md(ir)+1
     dx=dxr/float(md(ir))
     do 19 i=1,m
     x(i)=xr(ir)+float(i-1)*dx
  19 continue
     interpolate values on m row
C*----
     do 20 j=1,n,nd
     d(m,j)=dr(ir+1,((j-1)/nd+1))
     u(m,j)=ur(ir+1,((j-1)/nd+1))
     v(m,j)=vr(ir+1,((j-1)/nd+1))
  20 continue
     if(nd.gt.1) then
     do 21 jj=2,nr
     do 22 j=1,(nd-1)
     jjj=nd*(jj-2)+(j+1)
     d(m,jjj)=(dr(ir+1,jj)-dr(ir+1,jj-1))*y(jjj)/dyr
    1+(yr(jj)*dr(ir+1,jj-1)-yr(jj-1)*dr(ir+1,jj))/dyr
     u(m,jjj)=(ur(ir+1,jj)-ur(ir+1,jj-1))*y(jjj)/dyr
    1+(yr(jj)*ur(ir+1,jj-1)-yr(jj-1)*ur(ir+1,jj))/dyr
     v(m,jjj)=(vr(ir+1,jj)-vr(ir+1,jj-1))*y(jjj)/dyr
    1+(yr(jj)*vr(ir+1,jj-1)-yr(jj-1)*vr(ir+1,jj))/dyr
  22 continue
  21 continue
     endif
     interpolate values in x-direction
     do 23 i=2,m-1
     do 24 j=1,n
     d(i,j)=(d(m,j)-d(1,j))*x(i)/dxr+(x(m)*d(1,j)-x(1)*d(m,j))/dxr
     u(i,j)=(u(m,j)-u(1,j))*x(i)/dxr+(x(m)*u(1,j)-x(1)*u(m,j))/dxr
     v(i,j)=(v(m,j)-v(1,j))*x(i)/dxr+(x(m)*v(1,j)-x(1)*v(m,j))/dxr
   24 continue
   23 continue
   add in user specified grid subdivisions (read from unit iun(2))
C*----
     do 30 jr=1,nr-1
     if(isd(ir,jr).eq.1) then
     js=nd*jr+(1-nd)
     jf=js+nd
     read(iun(2),101)((d(i,j),j=js,jf),i=1,m)
     if(icur.eq.1)then
     read(iun(2),101)((u(i,j),j=js,jf),i=1,m)
     read(iun(2),101)((v(i,j),j=js,jf),i=1,m)
     endif
     do 31 i=1,m
```

```
do 32 j=js,jf
     d(i,j)=d(i,j)*dconv(iu)
     u(i,j)=u(i,j)*dconv(iu)
     v(i,j)=v(i,j)*dconv(iu)
32
     continue
     continue
31
     else
     end if
30
     continue
     add tidal offset to all rows and establish thin film
     do 33 i=1,m
     do 34 j=1,n
     d(i,j)=d(i,j)+tide
     if(d(i,j).lt.0.001) d(i,j)=0.001
34
     continue
     continue
     interpolation complete, return to model
C*----
     return
100 format(' model tried to put more spaces (md=',i4,
    1') than allowed in grid block ',i3)
101 format(16f8.4)
     end
```

```
C*
      subroutine con(ir,icount,ij)
C*
      subroutine calculates constants for row ij in reference grid
C*
      block ir.
      parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20,nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
     1dconv2(2)
      common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
      common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
      common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1, nii, fpeak
      common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
      dimension akd(2,iy,ncomp),xd(iy),xu(iy),xk(iy,ncomp)
     1,npts(ncomp),sumk(ncomp)
      complex a, w, az
      real k,kb
      eps=1.0e-05
      calculate constants
      do 1 j=1,n
      xd(j)=d(icount,j)
      xu(j)=u(icount,j)
    1 continue
      call vwnum(xd,xu,freqs,xk,eps,nfreqs,n)
      do 5 ifreq=1,nfreqs
      do 2 j=1,n
      k(ij,j,ifreq)=xk(j,ifreq)
    2 continue
    5 continue
      do 3 j=1,n
      do 4 ifreq=1,nfreqs
      sig(ij,j,ifreq)=freqs(ifreq)-k(ij,j,ifreq)*u(icount,j)
      akd(ij,j,ifreq)=k(ij,j,ifreq)*d(icount,j)
      q(ij,j,ifreq)=(1.+akd(ij,j,ifreq)/
     1(sinh(akd(ij,j,ifreq))*cosh(akd(ij,j,ifreq))))/2.
      p(ij,j,ifreq)=q(ij,j,ifreq)*9.80621*tanh(akd(ij,j,ifreq))
     1/k(ij,j,ifreq)
      dd(ij,j,ifreq)=(cosh(4.*akd(ij,j,ifreq))+8.-2.*
     1(tanh(akd(ij,j,ifreq))**2))/(8.*(sinh(akd(ij,j,ifreq))**4.))
    4 continue
    3 continue
      calculate the dissipation term w
      call diss(ir,icount,ij)
```

```
calculate the mean kb on each row
  do 6 ifreq=1,nfreqs
 npts(ifreq)=0
 sumk(ifreq)=0.
6 continue
 do 7 ifreq=1,nfreqs
  do 8 j=1,n
  if(d(icount,j).gt.0.05) then
  sumk(ifreq)=sumk(ifreq)+k(ij,j,ifreq)
 npts(ifreq)=npts(ifreq)+1
 else
 endif
8 continue
  if(npts(ifreq).eq.0) then
 kb(ij,ifreq)=k(ij,1,ifreq)
 kb(ij,ifreq)=sumk(ifreq)/float(npts(ifreq))
  endif
7 continue
 return
  end
```

```
subroutine fdcalc(ir,icount)
C*
      perform the crank-nicolson finite-difference calculations
c*
      on grid block ir. method is the implicit-implicit iteration
C*
      used by kirby and dalrymple(1983)
      parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20,nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
     1dconv2(2)
      common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
     1ioutput
      common/ref3/dxr,dyr,xr(ixr),yr(iyr),x(ix),y(iy)
      common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
      common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
      common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1, nii, fpeak
      common/wav2/amp(ncomp,nnd),dir(ncomp,nnd),tide
      common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
      common/nlin/an,anl,ntype
      real k,kb,ksth1,ksth2
      complex a, w, az
      complex first, second, third, fourth, fifth, sixth
      complex c1,c2,c3,cp1,cp2,cp3,ci,damp
      complex ac(iy,nnii),bc(iy,nnii),cc(iy,nnii),rhs(iy,nnii)
     1,sol(iy,nnii),rhso(iy,nnii)
      dimension ucp1(iy),uc(iy),vc(iy),vcp1(iy)
     1,dc(iy),dcp1(iy),ksth1(nnii),ksth2(nnii),s(iy)
      dimension thet(iy,nnii),sxy(iy),sxx(iy),syy(iy),urs(iy,nnii)
      dimension beta(iv)
      statement functions for crank-nicolson coefficients
      cg(i,j,ifreq)=sqrt(p(i,j,ifreq)*q(i,j,ifreq))
      pv(i,j,ifreq)=p(i,j,ifreq)-vc(j)*vc(j)
      pvp1(i,j,ifreq)=p(i+1,j,ifreq)-vcp1(j)*vcp1(j)
      bet(i,j,ifreq)=-4.*(k(i+1,j,ifreq)-k(i,j,ifreq))
     1/(dx*((k(i+1,j,ifreq)+k(i,j,ifreq))**2))
     1-4.*(k(i+1,j,ifreq)*(p(i+1,j,ifreq)-ucp1(j)**2)
     1-k(i,j,ifreq)*(p(i,j,ifreq)-uc(j)**2))
     1/(dx*((k(i+1,j,ifreq)+k(i,j,ifreq))**2.)
     1*(p(i+1,j,ifreq)+p(i,j,ifreq)-(ucp1(j)**2.+uc(j)**2)))
      dv(i,j,ifreq)=(cg(i+1,j,ifreq)+ucp1(j))/sig(i+1,j,ifreq)
     1-(cg(i,j,ifreq)+uc(j))/sig(i,j,ifreq)
     1+(-delta1)*dx*((vcp1(j+1)/sig(i+1,j+1,ifreq))
     1+(vc(j+1)/sig(i,j+1,ifreq))
```

```
1-(vcp1(j-1)/sig(i+1,j-1,ifreq))
1-(vc(j-1)/sig(i,j-1,ifreq)))/(2.*dy)
 damp(i,j,ifreq)=2.*ci*cdamp*((cg(i+1,j,ifreq)+ucp1(j))
1+(cg(i,j,ifreq)+uc(j)))/(dy*dy*(k(i+1,j,ifreq)**2
1+k(i,j,ifreq)**2))
 deltap(i,j,ifreq)=a1-b1*kb(i,ifreq)/k(i,j,ifreq)
 hrms(j)=h13(j)/(sqrt(2.))
 alpha1(j)=(0.75*sqrt(pi)*(b**3)*fpeak/(gam**4))
1*(hrms(j)/dcp1(j))**5
 alpha2(j)=0.75*sqrt(pi)*(b**3)
1*fpeak*(hrms(j)**5)/((gam**4)*(dc(j)**5))
 first(i,j,ifreq)=(cg(i+1,j,ifreq)+ucp1(j))
1*cmplx(1.,dx*(kb(i+1,ifreq)-a0*k(i+1,j,ifreq)))
1+cmplx(1.,0.)*(cg(i,j,ifreq)+uc(j)
1+dv(i,j,ifreq)*(sig(i+1,j,ifreq)+sig(i,j,ifreq))/4.)
1+2.*freqs(ifreq)*cmplx(0.,1.)*(-b1)*bet(i,j,ifreq)
1*(ucp1(j)+uc(j))/sig(i+1,j,ifreq)
1+4.*freqs(ifreq)*(-b1)*cmplx(0.,1.)*
1(3.*(ucp1(j)-uc(j))/dx+(vcp1(j+1)+vc(j+1)-vcp1(j-1)-vc(j-1))
1/(4.*dy))/(sig(i+1,j,ifreq)*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1+cmplx(-2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1+b1*bet(i,j,ifreq)*dx/(2.*dy*dy),(-deltap(i,j,ifreq))*dx
1/(2.*dy*dy))*(pvp1(i,j+1,ifreq)+2.*pvp1(i,j,ifreq)
1+pvp1(i,j-1,ifreq))/sig(i+1,j,ifreq)
 cp1(i,j,ifreq)=first(i,j,ifreq)
1-cmplx(1.,0.)*freqs(ifreq)*delta2*(3.*ucp1(j)+uc(j))
1/(2.*sig(i+1,j,ifreq))+ci*freqs(ifreq)*(a0-1.)*k(i+1,j,ifreq)
1*ucp1(j)*dx/sig(i+1,j,ifreq)+2.*ifilt*damp(i,j,ifreq)
1+cmplx(1.,0.)*2.*beta(j)*alpha1(j)*dx
 second(i,j,ifreq)=cmplx((-delta1)*dx*(vcp1(j)+vc(j))/(2.*dy)
1+b1*u2*bet(i,j,ifreq)*(ucp1(j)*vcp1(j)+uc(j)*vc(j))
1/(dy*sig(i+1,j+1,ifreq)),(-delta1*u2)*(ucp1(j+1)*vcp1(j+1)
1+uc(j+1)*vc(j+1)+2.*ucp1(j)*vcp1(j))/(2.*dy*sig(i+1,j+1,ifreq))
1+dx*(-b1)*bet(i,j,ifreq)*(sig(i+1,j,ifreq)*vcp1(j)
1+sig(i,j,ifreq)*vc(j))/(2.*dy*sig(i+1,j+1,ifreq)))
1+cmplx(2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1+(-b1)*bet(i,j,ifreq)*dx/(2.*dy*dy),-(-deltap(i,j,ifreq)*dx)
1/(2.*dy*dy))*(pvp1(i,j+1,ifreq)+pvp1(i,j,ifreq))
1/sig(i+1,j+1,ifreq)
1+4.*cmplx(0.,1.)*(-b1)*sig(i+1,j,ifreq)*vcp1(j)
1/(dy*sig(i+1,j+1,ifreq)*(k(i+1,j,ifreq)+k(i,j,ifreq)))
 cp2(i,j,ifreq)=second(i,j,ifreq)
1-ifilt*damp(i,j,ifreq)
 third(i,j,ifreq)=cmplx(-(-delta1)*dx*(vcp1(j)+vc(j))/(2.*dy)
1+(-b1)*u2*bet(i,j,ifreq)*(ucp1(j)*vcp1(j)+uc(j)*vc(j))
1/(dy*sig(i+1,j-1,ifreq)),-(-delta1)*u2*(ucp1(j-1)*vcp1(j-1)
```

```
1+uc(j-1)*vc(j-1)+2.*ucp1(j)*vcp1(j))
1/(2.*dy*sig(i+1,j-1,ifreq))
1-dx*(-b1)*bet(i,j,ifreq)*(sig(i+1,j,ifreq)*vcp1(j)
1+sig(i,j,ifreq)*vc(j))/(2.*dy*sig(i+1,j-1,ifreq)))
1+cmplx(2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1-b1*bet(i,j,ifreq)*dx/(2.*dy*dy),-(-deltap(i,j,ifreq)*dx)
1/(2.*dy*dy))*(pvp1(i,j,ifreq)+pvp1(i,j-1,ifreq))
1/sig(i+1,j-1,ifreq)-4.*cmplx(0.,1.)*(-b1)*sig(i+1,j,ifreq)
1*vcp1(j)/(dy*sig(i+1,j-1,ifreq)*(k(i+1,j,ifreq)+k(i,j,ifreq)))
 cp3(i,j,ifreq)=third(i,j,ifreq)
1-ifilt*damp(i,j,ifreq)
 fourth(i,j,ifreq)=cmplx(cg(i+1,j,ifreq)+ucp1(j)-
1dv(i,j,ifreq)*(sig(i+1,j,ifreq)+sig(i,j,ifreq))/4.,0.)
1+cmplx(1.,-dx*(kb(i,ifreq)-a0*k(i,j,ifreq)))
1*(cg(i,j,ifreq)+uc(j))+2.*cmplx(0.,1.)*freqs(ifreq)*(-b1)
1*bet(i,j,ifreq)*(ucp1(j)+uc(j))/sig(i,j,ifreq)
1+4.*cmplx(0.,1.)*freqs(ifreq)*(-b1)*(3.*(ucp1(j)-uc(j))/dx
1+(vcp1(j+1)+vc(j+1)-vcp1(j-1)-vc(j-1))
1/(4.*dy))/(sig(i,j,ifreq)*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1+cmplx(-2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))
1+(-b1)*bet(i,j,ifreq)*dx/(2.*dy*dy),-(-deltap(i,j,ifreq)*dx)
1/(2.*dy*dy))*(pv(i,j+1,ifreq)
1+2.*pv(i,j,ifreq)+pv(i,j-1,ifreq))/sig(i,j,ifreq)
 c1(i,j,ifreq)=fourth(i,j,ifreq)
1-cmplx(1.,0.)*freqs(ifreq)*delta2*(3.*ucp1(j)+uc(j))
1/(2.*sig(i,j,ifreq))-ci*freqs(ifreq)*(a0-1.)*k(i,j,ifreq)
1*uc(j)*dx/sig(i,j,ifreq)+2.*ifilt*damp(i,j,ifreq)
1-cmplx(1.,0.)*2.*(1-beta(j))*alpha1(j)*dx
 fifth(i,j,ifreq)=cmplx(-(-delta1)*dx*(vcp1(j)+vc(j))/(2.*dy)
1+b1*u2*bet(i,j,ifreq)*(ucp1(j)*vcp1(j)+uc(j)*vc(j))/
1(dy*sig(i,j+1,ifreq)),(-delta1)*u2*(ucp1(j+1)*vcp1(j+1)
1+uc(j+1)*vc(j+1)+2.*uc(j)*vc(j))/(2.*dy*sig(i,j+1,ifreq))
1+4.*(-b1)*sig(i,j,ifreq)*vc(j)/(dy*(k(i+1,j,ifreq)
1+k(i,j,ifreq))*sig(i,j+1,ifreq))
1-dx*(-b1)*bet(i,j,ifreq)*(sig(i+1,j,ifreq)*vcp1(j)
1+sig(i,j,ifreq)*vc(j))/(2.*dy*sig(i,j+1,ifreq)))
1+cmplx(2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))+b1
1*bet(i,j,ifreq)*dx/(2.*dy*dy),(-deltap(i,j,ifreq))*dx
1/(2.*dy*dy))
1*(pv(i,j+1,ifreq)+pv(i,j,ifreq))/sig(i,j+1,ifreq)
 c2(i,j,ifreq)=fifth(i,j,ifreq)
1-ifilt*damp(i,j,ifreq)
 sixth(i,j,ifreq)=cmplx((-delta1)*dx*(vcp1(j)+vc(j))/(2.*dy)
1+(-b1)*u2*bet(i,j,ifreq)*(ucp1(j)*vcp1(j)+uc(j)*vc(j))
1/(dy*sig(i,j-1,ifreq)),-(-delta1)*u2
1*(ucp1(j-1)*vcp1(j-1)+uc(j-1)*vc(j-1)+2.*uc(j)*vc(j))
1/(2.*dy*sig(i,j-1,ifreq))-4.*(-b1)*sig(i,j,ifreq)
1*vc(j)/(dv*(k(i+1,j,ifreq)+k(i,j,ifreq))*sig(i,j-1,ifreq))
1+dx*(-b1)*bet(i,j,ifreq)*(sig(i+1,j,ifreq)*vcp1(j)
1+sig(i,j,ifreq)*vc(j))/(2.*dy*sig(i,j-1,ifreq)))
1+cmplx(2.*(-b1)/(dy*dy*(k(i+1,j,ifreq)+k(i,j,ifreq)))-(-b1)
```

```
1*bet(i,j,ifreq)*dx/(2.*dy*dy),(-deltap(i,j,ifreq))*dx
1/(2.*dy*dy))*(pv(i,j,ifreq)+pv(i,j-1,ifreq))/sig(i,j-1,ifreq)
 c3(i,j,ifreq)=sixth(i,j,ifreq)
1-ifilt*damp(i,j,ifreq)
 f1(i,j,ifreq)=tanh(k(i,j,ifreq)*dc(j))**5.
 f1p1(i,j,ifreq)=tanh(k(i+1,j,ifreq)*dcp1(j))**5.
 f2(i,j,ifreq)=(k(i,j,ifreq)*dc(j)
1/sinh(k(i,j,ifreq)*dc(j)))**4.
 f2p1(i,j,ifreq)=(k(i+1,j,ifreq)*dcp1(j)
1/sinh(k(i+1,j,ifreq)*dcp1(j)))**4.
Booij coefficients.
 a0=1.
 a1 = -0.75
 b1 = -0.25
 do 1 j=1,n
 dc(j)=d(icount-1,j)
 uc(j)=u(icount-1,j)
 vc(j)=v(icount-1,j)
 dcp1(j)=d(icount,j)
 ucp1(j)=u(icount,j)
 vcp1(j)=v(icount,j)
 beta(j)=0.5+0.5*((0.001/dc(j))**3)
1 continue
 u2=1.0
 b=1.
 gam=0.6
 g=9.80621
 pi=3.1415927
 rho=1000.
 ci=cmplx(0.,1.)
 cdamp=0.025
 ifilt=1
 delta1=a1-b1
 delta2=1+2.*a1-2.*b1
 nii=nfreqs*nwavs
 it=0
 setup r.h.s.
 do 2 ii=1,nii
 rhs(1,ii)=cmplx(0.,0.)
 do 3 j=2, (n-1)
 rhso(j,ii)=c1(1,j,istore(ii))*a(1,j,ii)+c2(1,j,istore(ii))
 1*a(1,j+1,ii)+c3(1,j,istore(ii))*a(1,j-1,ii)
 rhs(j,ii)=rhso(j,ii)
 1-dx*w(1,j,ii)*a(1,j,ii)/2.-dx*cmplx(0.,1.)*an*anl
 1*sig(1,j,istore(ii))*k(1,j,istore(ii))*k(1,j,istore(ii))
```

```
1*dd(1,j,istore(ii))*(cabs(a(1,j,ii))**2.)
  1*a(1,j,ii)/2.-dx*cmplx(0.,1.)*an*(1.-anl)
  1*sig(1,j,istore(ii))*((1.+f1(1,j,istore(ii))*k(1,j,istore(ii))
  1*k(1,j,istore(ii))*(cabs(a(1,j,ii))**2.)
  1*dd(1,j,istore(ii)))*tanh(k(1,j,istore(ii))*dc(j)
  1+f2(1,j,istore(ii))*k(1,j,istore(ii))*0.5*h13(j))
  1/\tanh(k(1,j,istore(ii))*dc(j))-1.)*a(1,j,ii)/2.
 3 continue
   rhs(n,ii)=cmplx(0.,0.)
 2 continue
   return here for iterations
20 it=it+1
   if(it.eq.1)iii=1
   if(it.eq.2)iii=2
    establish boundary conditions
   if(ibc.eq.1)then
    do 4 ii=1,nii
    ksth1(ii)=real((2.*(a(1,2,ii)-a(1,1,ii)))/((a(1,2,ii)
   1+a(1,1,ii))*dy))*cmplx(0.,-1.))
    ksth2(ii)=real((2.*(a(1,n,ii)-a(1,n-1,ii))/((a(1,n,ii)
   1+a(1,n-1,ii))*dy))*cmplx(0.,-1.))
    bc(1,ii)=cmplx(1.,ksth1(ii)*dy/2.)
    cc(1,ii) = -cmplx(1.,-ksth1(ii)*dy/2.)
    bc(n,ii) = -cmplx(1.,-ksth2(ii)*dy/2.)
    ac(n,ii)=cmplx(1.,ksth2(ii)*dy/2.)
 4 continue
    else
    do 5 ii=1,nii
    bc(1,ii) = cmplx(1.,0.)
    cc(1,ii) = -bc(1,ii)
    bc(n,ii)=cmplx(1.,0.)
    ac(n,ii) = -bc(n,ii)
  5 continue
    endif
    coefficients for forward row
    do 6 ii=1,nii
    do 7 j=2, (n-1)
    ac(j,ii)=cp3(1,j,istore(ii))
    bc(j,ii)=cp1(1,j,istore(ii))+(dx/2.)*(w(2,j,ii))
   1+cmplx(0.,an*anl)*sig(2,j,istore(ii))*k(2,j,istore(ii))
   1*k(2,j,istore(ii))*dd(2,j,istore(ii))*(cabs(a(iii,j,ii))**2.)
   1*(dx/2.)+cmplx(0.,an*(1.-anl))
   1*sig(2,j,istore(ii))*(dx/2.)*((1.+f1p1(1,j,istore(ii))*
   1k(2,j,istore(ii))*k(2,j,istore(ii))*(cabs(a(iii,j,ii))**2.)
   1*dd(2,j,istore(ii)))*tanh(k(2,j,istore(ii))
   1*dcp1(j)+f2p1(1,j,istore(ii))*k(2,j,istore(ii))
   1*0.5*h13(j))/tanh(k(2,j,istore(ii))*dcp1(j))-1.)
    cc(j,ii)=cp2(1,j,istore(ii))
```

```
7 continue
  6 continue
    update solution one step
    call vtrida(1,n,ac,bc,cc,rhs,sol,nii)
    do 8 ii=1,nii
    do 9 j=1,n
    a(2,j,ii)=sol(j,ii)
    sol(j,ii)=cmplx(0.,0.)
  9 continue
  8 continue
    if(it.eq.1) go to 20
                           ____
    Check Ursell parameter for Stokes solution
    if (ntype.eq.2) then
    do 23 j=1,n
    do 24 ii=1,nii
    urs(j,ii)=(cabs(a(2,j,ii))/dcp1(j))
   1/((k(2,j,istore(ii))*dcp1(j))**2)
    if(urs(j,ii).gt.0.5) write(*,204)urs(j,ii),icount,j,ii
24 continue
23 continue
     endif
    calculate reference phase function
    do 10 ifreq=1,nfreqs
    psibar(ifreq)=psibar(ifreq)+(kb(2,ifreq)+kb(1,ifreq))*dx/2.
  10 continue
     if(icount.eq.m) then
     open(111,file='alf.dat')
    write(111,201) ((alpha1(j)), j=1,n)
201 format(100e10.3)
    endif
    calculate significant waveheight h13
    if(nii.eq.1) then
    do 33 j=1,n
    h13(j)=2.*cabs(a(2,j,1))
  33 continue
     else
    do 25 j=1,n
     s(j)=0.
     do 26 ii=1,nii
     s(j)=s(j)+(cabs(a(2,j,ii))**2)
  26 continue
    h13(j)=sqrt(8.*s(j))
  25 continue
     endif
     if (icount.eq.m) then
```

```
c Calculate wave angles at reference grid rows. Note: angles are not well
c defined in a directional, multicomponent sea, or where waves become short
c crested. This routine was heavily modified by Raul Medina, University of
c Cantabria.
      do 16 ii=1,nii
      do 15 j=1,n
      if(a(2,j,ii).eq.(0.,0.)) then
          akx2=0.
      else
          akx2=aimag(clog(a(2,j,ii)))
      endif
      if(a(1,j,ii).eq.(0.,0.)) then
          akx1=0.
      else
          akx1=aimag(clog(a(1,j,ii)))
      endif
      if(abs(akx2-akx1).gt.pi)then
      akx=sign((2.*pi-(abs(akx1)+abs(akx2)))/dx,akx1)
      akx=(akx2-akx1)/dx
      endif
      if(j.ne.n) then
          if(a(2,j+1,ii).eq.(0.,0.)) then
              aky2=0.
          else
              aky2=aimag(clog(a(2,j+1,ii)))
          if(a(2,j,ii).eq.(0.,0.)) then
              aky1=0.
          else
              aky1=aimag(clog(a(2,j,ii)))
          endif
      else
          if(a(2,j,ii).eq.(0.,0.)) then
              aky2=0.
          else
              aky2=aimag(clog(a(2,j,ii)))
          if(a(2,j-1,ii).eq.(0.,0.))then
              aky1=0.
          else
              aky1=aimag(clog(a(2,j-1,ii)))
          endif
      endif
      if(abs(aky2-aky1).gt.pi)then
          aky=sign((2.*pi-(abs(aky1)+abs(aky2)))/dy,aky1)
      else
          aky=(aky2-aky1)/dy
      thet(j,ii)=atan2(aky,(akx+kb(2,istore(ii))))
 15
      continue
 16
      continue
```

```
Estimation of radiation stresses
     do 17 j=1,n
     sxx(j)=0.
     syy(j)=0.
     sxy(j)=0.
 17 continue
     do 22 j=1,n
     do 18 ii=1,nii
     sxx(j)=sxx(j)+(cabs(a(2,j,ii))**2)
    1*(q(2,j,istore(ii))*(1+(cos(thet(j,ii))**2))-0.5)
     syy(j)=syy(j)+(cabs(a(2,j,ii))**2)
    1*(q(2,j,istore(ii))*(1+(sin(thet(j,ii))**2))-0.5)
     sxy(j)=sxy(j)+q(2,j,istore(ii))
    1*(cabs(a(2,j,ii))**2)*sin(thet(j,ii))
  18 continue
     sxx(j)=sxx(j)*rho*g/2.
     syy(j)=syy(j)*rho*g/2.
     sxy(j)=sxy(j)*rho*g/4.
  22 continue
    Line printer output
     mm1=m-1
     write(*,205) (ir+1),mm1
     write(*,202) x(m)/dconv(iu)
     write(*,206)
     write(*,203) (h13(j)/dconv(iu),j=1,n,nd)
     Output of radiation stresses
     write(22,*) (sxx(j)/dconv2(iu), j=1,n,nd)
     write(22,*) (sxy(j)/dconv2(iu), j=1,n,nd)
     write(22,*) (syy(j)/dconv2(iu),j=1,n,nd)
     endif
    return control back to model
     return
202 format(' x=',f10.2)
203 format(',12(f10.4))
204 format(' '//' Warning: Ursell number =',f10.4,' encountered at','g
    1rid location', i6,',', i6/' during computation of wave component', i3,
    1'should be using Stokes-Hedges model
    1(ntype=1) due to shallow', 'water')
205 format(' grid row ir=',i3,', ',i3,' x-direction subdivisions',
    1' used')
206 format(' Significant Wave Height')
     end
```

```
subroutine rbcon
C*
      rolls back constants and solution to row 1
C*
      parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20, nnii=200)
      common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
     1dconv2(2)
      common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
      common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
      common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
      common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
     1, nii, fpeak
      common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
      complex a, w, az
      real k,kb
      do 1 j=1,n
      do 2 ifreq=1,nfreqs
      k(1,j,ifreq)=k(2,j,ifreq)
      sig(1,j,ifreq)=sig(2,j,ifreq)
      q(1,j,ifreq)=q(2,j,ifreq)
      p(1,j,ifreq)=p(2,j,ifreq)
      dd(1,j,ifreq)=dd(2,j,ifreq)
      w(1,j,ii)=w(2,j,ii)
      k(2,j,ifreq)=0.
      sig(2,j,ifreq)=0.
      q(2,j,ifreq)=0.
      p(2,j,ifreq)=0.
      dd(2,j,ifreq)=0.
      w(2,j,ii)=0.
    2 continue
    1 continue
      do 3 ifreq=1,nfreqs
      kb(1,ifreq)=kb(2,ifreq)
      kb(2,ifreq)=0.
    3 continue
      roll back solution
      do 4 j=1,n
      do 5 ii=1,nii
      a(1,j,ii)=a(2,j,ii)
      a(2,j,ii)=cmplx(0.,0.)
    5 continue
    4 continue
      return
      end
```

```
subroutine vwnum(dt,u,freqs,k,eps,nfreqs,n)
C*
     vectorized wavenumber calculations. variable definitions
C*
     same as in subroutine wnum.
C*
c*
     calculate wavenumber k according to the form
         s*s-2*s*k*u+k*k*u*u=g*k*tanh(k*d)
C*
C*
     where
C*
         d = local water depth
C*
         s = absolute frequency
         g = gravitational acceleration constant
         u = x-component of ambient current
         eps = tolerance for iteration convergence
C*
         i, j = indices in finite-difference grid
C*
C*
      solution by newton-raphson iteration using eckart's
C*
C*
      approximation as a seed value.
     parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
     1nnd=20, nnii=200)
      common/ref2/dr(ixr,iyr),ur(ixr,iyr),vr(ixr,iyr),iun(8),iinput,
C
     1ioutput
      dimension dt(iy),u(iy),freqs(ncomp),f(iy,ncomp),
     1fp(iy,ncomp)
      real k(iy,ncomp),kn(iy,ncomp)
      g = 9.806
     pi=3.1415927
                           _____
      calculate first guess
      do 1 ifreq=1,nfreqs
      do 2 j=1,n
      k(j,ifreq)=freqs(ifreq)*freqs(ifreq)/
     1(g*sqrt(tanh(freqs(ifreq)*freqs(ifreq)*dt(j)/g)))
    2 continue
    1 continue
      do 4 ifreq=1,nfreqs
      do 5 j=1,n
      do 3 i=1,40
      f(j,ifreq)=freqs(ifreq)**2-2.*freqs(ifreq)*k(j,ifreq)*u(j)+
     1(k(j,ifreq)*u(j))**2-g*k(j,ifreq)*tanh(k(j,ifreq)*dt(j))
      fp(j,ifreq)=-2.*freqs(ifreq)*u(j)+2.*k(j,ifreq)*(u(j)**2)-
     1g*tanh(k(j,ifreq)*dt(j))-g*k(j,ifreq)*dt(j)
     1/(cosh(k(j,ifreq)*dt(j))**2.)
```

```
kn(j,ifreq)=k(j,ifreq)-f(j,ifreq)/fp(j,ifreq)
if((abs(kn(j,ifreq)-k(j,ifreq))/kn(j,ifreq)).lt.eps) go to 10
k(j,ifreq)=kn(j,ifreq)
3 continue
  write(*,*)'Wavenumber failed to converge for frequency component'
1,ifreq,'on column',j
10 k(j,ifreq)=kn(j,ifreq)
5 continue
4 continue
return
end
```

```
C*
     subroutine vtrida(if,l,a,b,c,d,v,mm)
C*
C*
     vectorized tridiagonal matrix solution by double sweep algorithm.
     present subroutine adopted from the subroutine described in:
C*
C*
         carnahan, luther and wilkes, applied numerical
C*
         methods, wiley, 1969
C*
     modified to handle complex array coefficients and solution
C*
     values. input and output are
C*
C*
         a,b,c = coefficients of row in tridiagonal matrix
C*
                = right hand side vector of matrix equation
                = solution vector
C*
                = beginning and end indices of positions in the
C*
                  dimensioned range of the column vector to be
C*
                  considered.
C*
                  _____
     parameter(iy=275,nnii=200)
     complex a(iy,nnii),b(iy,nnii),c(iy,nnii),d(iy,nnii)
     1,v(iy,nnii),beta(iy,nnii),gamma(iy,nnii)
C*----
     compute intermediate vectors beta and gamma
     do 1 j=1,mm
     beta(if,j)=b(if,j)
     gamma(if,j)=d(if,j)/beta(if,j)
    1 continue
     ifp1=if+1
     do 2 i=ifp1,1
      do 3 j=1,mm
     beta(i,j)=b(i,j)-a(i,j)*c(i-1,j)/beta(i-1,j)
      gamma(i,j)=(d(i,j)-a(i,j)*gamma(i-1,j))/beta(i,j)
    3 continue
    2 continue
    compute solution vector v
     do 4 j=1,mm
      v(1,j)=gamma(1,j)
    4 continue
      last=l-if
      do 5 k=1,last
      i=1-k
      do 6 j=1,mm
      v(i,j)=gamma(i,j)-c(i,j)*v(i+1,j)/beta(i,j)
    6 continue
    5 continue
```

return end

```
C*
     subroutine diss(ir,icount,i)
C*
     subroutine calculates the dissipation at a single grid point
C*
     based on values of the switch iw at that point.
     parameter(ixr=251,iyr=275,ix=20,iy=275,ncomp=10,
    1nnd=20,nnii=200)
     common/ref1/mr,nr,ispace,nd,md(ixr),iu,dconv(2),if(3),icur,ibc,
    1dconv2(2)
     common/con1/q(2,iy,ncomp),p(2,iy,ncomp),sig(2,iy,ncomp)
     common/con2/k(2,iy,ncomp),kb(2,ncomp),w(2,iy,nnii),dd(2,iy,ncomp)
     common/block1/d(ix,iy),u(ix,iy),v(ix,iy),m,n,dx,dy
     common/wav1/iwave,nfreqs,freqs(ncomp),nwavs,istore(nnii)
    1, nii, fpeak
     common/comp/a(2,iy,nnii),psibar(ncomp),az(iy,nnii),h13(iy)
     real k,kb,nu,cp,kd(iy,nnii)
     complex w,a,az
     sq(i,j,ifreq)=sqrt(nu/(2.*sig(i,j,ifreq)))
     nu=1.3e-06
     cp=4.5e-11
     g=9.80621
     pi=3.1415927
        ______
     value of f here is value assuming tau=(f/8)*u**2
C*
     f=4.*fw; fw=wave friction factor
     f=0.01*4.0
     do 1 j=1,n
     do 2 ii=1,nii
     w(i,j,ii)=cmplx(0.,0.)
     kd(j,ii)=k(i,j,ii)*d(icount,j)
   if if(1) = 1, use turbulent boundary layer damping
     if(if(1).eq.1) w(i,j,ii)=2.*f*cabs(az(j,ii))
    1*sig(i,j,istore(ii))*k(i,j,istore(ii))/(sinh(2.*kd(j,ii))
    1*sinh(kd(j,ii))*3.*pi)
     if(if(2) = 1, add porous bottom damping
     if(if(2).eq.1) w(i,j,ii)=w(i,j,ii)+(g*k(i,j,istore(ii))
    1*cp/(nu*(cosh(kd(j,ii))**2)))*cmplx(1.,0.)
C*----
    if if(3) = 1, add boundary layer damping
     if(if(3).eq.1) w(i,j,ii)=w(i,j,ii)+2.*k(i,j,istore(ii))
     1*sig(i,j,istore(ii))*sq(i,j,istore(ii))
    1*(1.+(\cosh(kd(j,ii))**2))*cmplx(1.,-1.)/sinh(2.*kd(j,ii))
    2 continue
```

1 continue return end

6 Appendix B: PLOTAMP Program Listing

The following section contains the listing of the program **PLOTAMP** as an aid to read the stored output file *outdat.dat*. This program was used to create contour plots of normalized significant wave height in the example cases shown in this manual.

```
_____
     plotamp.f
C*
C*
     Plot significant wave height field constructed from output
C*
     in 'outdat.dat' from a REF/DIF S model.
     dimension wave (251,301), h(251,301), x(251), y(301)
     dimension hs(251,301)
     open(10,file='outdat.dat')
     read(10,*)n,m
     write(*,*)' grid dimensions are ',m,' by ',n
     read(10,*)(y(j),j=1,n)
     do 1 i=1,m
     read(10,*)x(i)
     read(10,*)(h(i,j),j=1,n)
     read(10,*)(hs(i,j),j=1,n)
     continue
     close(10)
     h0=hs(1,1)
     hh=h(1,1)
     construct normalized wave height field
     do 3 i=1,m
     do 2 j=1,n
     wave(i,j)=hs(i,j)/h0
     h(i,j)=h(i,j)/hh
2
     continue
     continue
     enter corner points for plotted grid
C*----
     write(*,*)' enter imin,imax,jmin,jmax'
     read(*,*)imin,imax,jmin,jmax
     mnew=(imax-imin)+1
     nnew=(jmax-jmin)+1
     do 5 i=1,mnew
     ii=(i+imin)-1
     do 4 j=1, nnew
     jj=(j+jmin)-1
     wave(i,j)=wave(ii,jj)
     h(i,j)=h(ii,jj)
 4
     continue
 5
     continue
     contour the normalized wave height field
C*----
     call gopks(6,idum)
     call gopwk(1,2,1)
     call gacwk(1)
     write(*,*)' enter contour interval for depth'
```

```
read(*,*)cont
call conrec(h,251,mnew,nnew,0.,3.,cont,0,-1,200)
write(*,*)' enter contour interval for amplitude'
read(*,*)cont
call conrec(wave,251,mnew,nnew,0.,3.,cont,0,-1,0)
call frame
call gdawk(1)
call gclwk(1)
call gclks
stop
end
```

7 Appendix C: fweb Documentation of the SPECGEN Program Listing

The following section contains a document which provides a heavily anotated program listing of SPECGEN, which is the preprocessing program used to create most of the input files to REF/DIF S for the example cases in this manual.

1. Data Generation for REF/DIF S: Introduction and Program Input.

This program generates the input file *indat.dat* for **REF/DIF** S Version 1.0, which models refraction and diffraction of two-dimensional spectra on arbitrary bathymetry.

The two-dimensional spectrum is obtained by generating a TMA spectrum which is used in conjunction with an angular spreading function. The resulting two-dimensional spectrum is divided into discrete components. The generated input file can be used with REF/DIF S as well as REF/DIF 1.

Some variable definitions are as follows:

f : Frequency array in (Hz)

sptma : TMA spectrum array in (cm²sec)

ff : Array containing cut-off values of frequency bins in (Hz)

thet : Directional array in (rad)

spd : Wrapped normal directional distribution array in (rad⁻¹)

tt : Array containing cut-off values for directional bins in (rad)

The input to SPECGEN is read in free format from the file specgen.dat which contains information about the domain and options desired to be used in REF/DIF S as well as information about the desired spectra. The first part of information is not used in the computational part of SPECGEN but is just past to the input file of the model REF/DIF S, whereas information about the spectra is used to generate and discretize the spectra. This results in discrete wave components to be input into REF/DIF S.

The input file specgen.dat is set up in the following manner:

The parameters in this first part are not related to the operation of SPECGEN, but will control the operation of the model REF/DIF S

1. iun(1), iun(2), iun(3) :

Logical device numbers of reference grid data, user-specified subdivision data and reference grid output.

2. mr, nr: Reference grid dimensions.

3. iu, ntype, icur, ibc

iu is the switch for physical units; iu = 1, MKS; iu = 2, English.

ntype is switch for nonlinearity; ntype = 0, linear model: ntype = 1, composite model; ntype = 2, Stokes wave model.

icur is switch for input current data. icur = 0, nocurrents input; icur = 1, currents input.

ibc is the boundary condition switch. ibc = 0, closed boundaries; ibc = 1, open boundaries.

- 4. dxr, dyr, dt : Reference grid x-spacing and y-spacing and depth tolerance value.
- 5. ispace, nd

ispace is switch controlling subdivisions. ispace = 0, program attempts its own x subdivisions; ispace = 1, user specifies x subdivisions.

nd is the number of y subdivisions.

If ispace = 1, the x subdivisions are inserted here.

$$(md(ir), ir = 1, mr - 1)$$

6. iff(1), iff(2), iff(3) :Dissipation switches.

iff(1) = 1, turn on turbulent boundary layer; iff(2) = 1, turn on porous bottom damping; iff(3) = 1, turn on laminar boundary layer.

No damping if all values are zero.

7. isp : Switch for user-specified subgrid specifications.

isp = 0, no subgrids to be read; isp = 1, subgrids will be read.

- 8. tide : Tidal offset.
- iopt : iopt specifies whether amplitudes values of a specific components are to be stored in file specamp.dat.

iopt = 1 turns the option on;

iopt = 0 skips the option. The default value is zero.

If iopt = 1 the desired wave component is specified here:

ifcomp, *idcomp* : The amplitude output is to be done for frequency component *ifcomp* and directional component *idcomp*.

This line concludes the first part of the input file. The following lines of input specify the input spectra.

10. depth

: Water depth in (m)

11. f_m

: Peak frequency of frequency spectrum in (Hz)

12. f_{max}

: Maximum frequency that is desired to be resolved in (Hz)

13. H_s

: Significant wave height in (m)

14. γ

: Peak enhancement factor

15. nebin

: Desired number of frequency bins

16. θ_m

: Mean direction in (°)

17. σ_m

: Directional spreading parameter in (°)

18. ndbin

: Desired number of directional bins

2. The Code.

program specgen

3. The definition and input sections of the code are as follows.

```
\langle \text{ Common Declarations } 15 \rangle
dimension f(200000), tma(200000), ff(200000), d(200000), tt(200000), thet(200000)
\langle \text{ Open Statements } 16 \rangle
call readin
```

4. Additional parameters are defined and conversions are performed.

```
depth = depth*100.
hs = hs*100.
pi = 4.*atan(1.)
g = 980.621
l = 20
deltaf = 0.025
fmin = deltaf
nn = nint(fmax/deltaf)
deltat = pi*4.5/180.
thetm = pi*thetm/180.
sigm = pi*sigm/180.
```

5. Construction of frequency spectrum array.

```
alpha = 1.
\mathbf{do} \ 1 \ i = 1, nn
f(i) = deltaf * float(i)
tma(i) = sptma(f(i))
```

1: continue

write (*, *) 'Frequency spectrum computed'

Computation of area under frequency spectrum and proper value of α .

33: continue

$$\mathbf{call}\ integ(1,1,tma,1,nn,deltaf,am)$$

Determination of cut-off points.

$$ij = 1$$

call $integ(2, 1, tma, ij, nn, deltaf, am)$
 $fsave = f(ij)$
 $nn = nn - (ij - 1)$

do $6 i = 1, nn$
 $f(i) = fsave + float(i - 1)*deltaf$
 $tma(i) = sptma(f(i))$

write $(20, *) f(i)$

write $(30, *) tma(i)$

6: continue

 $a = am/float(nebin)$

Integration for bins. 8.

$$\mathbf{call}\ bins(1, deltaf, nn, nebin, f, tma, a, f\!f)$$

9. Construction of directional spectrum array.

$$nn = 81$$

do 3 $i = 1, nn$
 $thet(i) = -(pi/2.) + float(i-1)*deltat$

3: continue

do 4 $i = 1, nn$
 $d(i) = spd(thet(i))$

4: continue

- write(*,*) 'Directional_spectrum_computed'
- 10. Computation of area under directional spectrum.
 call integ(1, 2, d, 1, nn, deltat, ad)
- 11. Determination of cutoff points.

$$\begin{split} ij &= 1 \\ \textbf{call } integ(2,2,d,ij,nn,deltat,ad) \\ tsave &= thet(ij) \\ nn &= nn - (ij-1) \\ \textbf{do } 8 \ i &= 1,nn \\ thet(i) &= tsave + \text{float}(i-1)*deltat \\ d(i) &= spd(thet(i)) \\ \textbf{write}(40,*) \ thet(i) \\ \textbf{write}(50,*) \ d(i) \end{split}$$

8: continue

$$aa = ad/float(ndbin)$$

12. Determination of amplitudes of components.

$$amp = \operatorname{sqrt}(2.*a*aa)/100.$$

13. Integration for bins.

call bins(2, deltat, nn, ndbin, thet, d, aa, tt)

14. Assure symmetry of directional spectrum.

```
ndd = ndbin/2 + 1
\mathbf{do} \ 15 \ ii = 1, ndd
dif = \mathsf{abs}(\mathsf{abs}(tt(ii)) - \mathsf{abs}(tt(ndbin - ii + 2)))
\mathbf{if} \ (dif > 1. \cdot 10^{-08}) \ \mathbf{then}
tt(ndbin - ii + 2) = (\mathsf{abs}(tt(ii)) + \mathsf{abs}(tt(ndbin - ii + 2)))/2.
tt(ii) = -tt(ndbin - ii + 2)
\mathbf{endif}
15: continue
tt(ndd) = 0.
\langle \ \mathsf{Output} \ \mathsf{Statements} \ \mathsf{17} \ \rangle
\mathbf{end}
```

15. Common and Statements.

This code is used in sections 3 and 18.

16. Open Statements.

This code is used in section 3.

17. Output Statements.

This code is used in section 14.

18. Routine to Read Input.

This routine reads input from file specgen.dat.

```
subroutine readin
  (Common Declarations 15)
    // Read or specify parameters for 'indat.dat'
  read(10,*)(iun(i), i = 1,3)
  read(10,*) mr, nr
  read(10,*) iu, ntype, icur, ibc
  read(10,*) dxr, dyr, dt
  read (10,*) ispace, nd
  read(10,*) (iff(i), i = 1,3)
  read(10,*) isp
  read(10,*) tide
  read(10,*) iopt
  if (iopt \equiv 1) then
    read(10,*) if comp, idcomp
  else
    iopt = 0
  endif
  iinput = 1
  ioutput = 1
  iwave = 1
    // Read parameters for setup and discretezation of spectra
  read(10,*) depth
  read(10,*) fpeak
  read(10,*) fmax
  read(10,*) hs
  read(10,*) gamma
  read(10,*) nebin
  read(10,*) thetm
  read(10,*) sigm
  read(10,*) ndbin
  return
end
```

19. Function to Generate TMA Spectrum.

This routine, originally written by Vincent and Briggs, calculates a TMA spectrum using input parameters provided by the user in the second part of the file *specgen.dat*.

The TMA spectrum is given by the energy density E(f) for frequency f,

$$E(f) = \frac{\alpha g^2}{(2\pi)^4 f^5} \exp\left\{-1.25 \left(\frac{f_m}{f}\right)^4 + (\ln \gamma) \exp\left[\frac{-(f - f_m)^2}{2\sigma^2 f_m^2}\right]\right\} \phi(f, h)$$
(47)

The parameters α , f_m and γ were defined while describing the input. Furthermore σ = Shape parameter defined in the program by,

$$\sigma = \begin{cases} \sigma_a = 0.07 & \text{if } f < f_m \\ \sigma_b = 0.09 & \text{if } f \ge f_m \end{cases}$$

$$\tag{48}$$

The factor $\phi(f, h)$ incorporates the effect of the depth h and is computed following Hughes (1984) by,

$$\phi = \begin{cases} 0.5 (\omega_h)^2 & \text{if } \omega_h < 1\\ 1 - 0.5 (2 - \omega_h)^2 & \text{if } 1 \le \omega_h \le 2\\ 1 & \text{if } \omega_h > 2 \end{cases}$$
(49)

where,

endif

$$\omega_h = 2\pi f \sqrt{\frac{h}{g}}$$

Broad or narrow frequency spectra can be obtained assigning the values 2 and 20 to the parameter γ , respectively. The code follows.

```
function sptma(ff)

common /blk1/fpeak, alpha, gamma, depth, g, fmin, fmax

sa = 0.07

sb = 0.09

twopi = 8.*atan(1.)

if (ff < fmin) then

sptma = 0.

return

endif

if (ff > fmax) then

sptma = 0.

return
```

20. Calculation of four factors:

- High frequency tail equilibrium range factor (p1)
- Low frequency tail factor (Pierson-Moskowitz) (p2)
- Jonswap factor (p3)
- TMA spectrum factor (phi)

$$p1 = (alpha*g*g)/(1558.55*ff^5)$$
 $z2 = -1.25*(ff/fpeak)^{-4}$
 $if(z2 < -28.)$ $z2 = -28$
 $p2 = exp(z2)$
 $ss = sa$
 $if(ff > fpeak)$ $ss = sb$
 $bb = (ff - fpeak)/(ss*fpeak)$
 $bb = .5*bb^2$
 $if(bb > 28)$ $bb = 28$
 $p3 = gamma exp(-bb)$
 $omh = twopi*ff*sqrt(depth/g)$
 $if(omh \le 1.)$ $phi = omh^2/2$.
 $if(omh > 1.)$ $phi = 1. - 0.5*(2. - omh)^2$
 $if(omh \ge 2.)$ $phi = 1$.

21. TMA Spectrum.

$$sptma \,=\, p1*p2*p3*phi$$

return

end

22. Function to Calculate Directional Spectrum.

A directional spectrum $D(\theta)$ was used in conjunction with the frequency spectrum to create a two-dimensional spectrum. The angular spectrum is defined by,

$$D(\theta) = \frac{1}{2\pi} + \frac{1}{\pi} \sum_{i=1}^{J} \exp\left[-\frac{(j\sigma_m)^2}{2}\right] \cos j \left(\theta - \theta_m\right)$$
 (50)

where,

end

 θ_m = mean wave direction = 0^0 .

J = number of terms in the series chosen as 20 in the numerical calculations.

 σ_m is chosen to be 10^0 or 30^0 giving a narrow or broad directional spreading, respectively. All three parameters are input. The code follows.

```
function spd(thet)

common /blk2 / sigm, thetm, l

pi = 4.*atan(1.)

sum = 0.

do 1 i = 1, l

sum = sum + exp(-0.5*(i*sigm)^2)*cos(i*(thet - thetm))

1: continue

spd = (1./(2.*pi)) + (1./pi)*sum

return
```

23. Routine to Calculate Areas under Spectra.

This routine calculates the area under a function fun(i) using Simpson's Rule with a step size of df. The routine performs two different functions.

24. Computation of Total Area.

If option 1 is chosen for the variable opt in the call statement the total area under the curve of the function fun(i) is computed and returned to the calling routine.

```
subroutine integ(opt, ii, fun, isave, n, df, area)

dimension fun(n)

integer opt

if (opt \equiv 1) then

s1 = fun(2)
s2 = fun(3)

do 1 i = 1, n/2 - 2
s1 = s1 + fun(2 + 2*i)
s2 = s2 + fun(3 + 2*i)

1: continue

area = (df/3.)*(fun(1) + 4.*s1 + 2.*s2 + fun(n))

return
```

25. Determination of Cut-off Points.

If option 2 is chosen the left tail of the function fun(i) containing 0.25% of the total area is cut off. If fun(i) is representing the frequency spectrum the right tail containing 1% of the total area is cut off. For the directional spectrum this value is 0.25%. The new values of the starting and ending locations of the array fun(i) are returned to the calling routine.

```
elseif (opt \equiv 2) then
    asave = area
    area = 0.
    icount = 1
   do 2 i = isave, n - 2, 2
      area = area + (df/3.)*(fun(i) + 4.*fun(i+1) + fun(i+2))
      if (icount \equiv 1) then
        if (area > (0.0025*asave)) then
           icount = icount + 1
           isave = i
        endif
      endif
      if (ii \equiv 1 \land area > (0.99*asave)) go to 10
      if (ii \equiv 2 \land area > (0.9975*asave)) go to 10
      if (i \equiv (n-2)) write (*,*) 'Doing Last Larray member'
 2: continue
10: n = i + 2
   return
 endif
 end
```

26. Routine to Divide Spectra into Equal-Energy Bins.

This routine integrates a given function f(x) until the desired energy value, i.e. area value is reached. If the resolution of the given disscrete function is not high enough, twice the resolution is introduced and the spectrum for the new resolution computed. Integration is started again at the bin where convergence was not achieved. Integration is performed by Simpson's Rule.

The program structure is as follows.

```
subroutine bins(ni,dx,n,nbin,x,f,area,xx)

common /blk1/fpeak, alpha, gamma, depth, g, fmin, fmax

common /blk2/sigm, thetm, l

dimension x(200000), f(200000), xx(200000)

real in, ins, inb

eps = 1. \cdot 10^{-03}

in = 0.

isave = 1

dxsave = dx

ii = 1

xx(1) = x(1)
```

27. Main Integration Loop.

This loop integrates the function. If convergence is achieved the integration for that bin is completed. Integration is started again for the next bin.

If the integral value is too small at one step and too big at the other, the resolution is increased by simply dividing Δx by 2 and setting up the spectrum again with the new, finer resolution. Integration is started again at the last point where the integral alue was smaller than the desired value.

```
80: do 2 i = isave, n - 2, 2
      in = in + (dx/3.)*(f(i) + 4.*f(i+1) + f(i+2))
        // Check for convergence
      if (abs(in - float(ii)*area) < eps) then
        write (*, *) 'Bin_\', ii, '\completed.'
        xx(ii+1) = x(i+2)
        ii = ii + 1
        isave = i + 2
        go to 80
      endif
      if (in < (float(ii)*area)) then
        ins = in
      elseif (in > (float(ii)*area)) then
        inb = in
        dx = dx/2.
        n = 2*n - 1
        isave = 2*i - 1
        xr = x(i)
        in = ins
        if (ni \equiv 1) then
           do 4 j = isave, n
             x(j) = dx *float(j-1) + dxsave
             f(j) = sptma(x(j))
        4: continue
        elseif (ni \equiv 2) then
           do 5 j = isave, n
             x(j) = xr + \mathsf{float}(j - isave) * dx
             f(j) = spd(x(j))
```

```
5: continue
endif
go to 80
endif
2: continue
xx(nbin + 1) = x(n)
return
end
```

8 Appendix D: Answers to Frequently Asked Questions

The following section contains answers to frequently asked questions.

1. What is the maximum current this model can handle?

The program can handle currents that are of the same order as the wave speed. However, opposing currents which are strong enough to stop the waves are not handled properly. This extension will be added in the future.

- 2. In which situations should the user specify totally reflecting lateral boundaries (ibc = 0), and when should partially transmitting lateral boundaries (ibc = 1) be used? It is common practice to use the closed boundary condition as the amount of reflection can be identified and the domain width can be chosen so that there are no reflections in the region of interest. The final run may be made using partially transmitting boundaries, but it should be noted that there will still be a certain amount of reflected waves present in the domain.
- 3. What is the difference between a user-specified subdivisions and user-specified subgrid?

User-specified subdivisions and user-specified subgrids are two completely independent options. The user has to specify the desired number of subdivisions in the y-direction nd. However, the user has an option between specifying the number of subdivisions in the x-direction md(ir) or having the program specify those. If the program specifies md(ir), the switch ispace has to be set to zero and the program performs interpolations of the depth and velocity grids. If the user choses to specify the md(ir) values himself, the switch ispace has to be set to unity. In this case, the user again has a choice: The user can specify isp = 0, then the program will perform interpolations. The user can set isp = 1, in which case he/she has to specify a subgrid in the file subdat.dat.

4. What are the guidelines for the specification of the number of subdivisions md(ir) in the x-direction?

It is recommended that the number of subdivisions md(ir) be determined by the program. However, if the user wants to specify md(ir)'s, the user should make sure that the subdivided grid is at least as fine as the program would have determined. In order to achieve this, the user can first run the program by letting it pick its own subdivisions. Then the user can chose his/her subdivision using the programs subdivisions as a guideline.

5. What are the guidelines for the choice of the number of subdivisions nd in the y-direction?

nd should be chosen such that the final subdivided reference grid cells have increments in the x- and y-directions that are fairly close in size.

6. Where is the origin of the grid?

The origin of the domain is always chosen to be the lower right-hand corner of the domain with the x-axis pointing in the propagation direction and the y-axis pointing towards the left.